



LUXE

Proposal for a new experiment using a
Laser and XFEL to test quantum physics in the strong-field regime

Beate Heinemann (DESY and University of Freiburg)
on behalf of LUXE Collaborators

Birmingham, February 19th 2020





OUTLINE

LUXE = “Laser Und XFEL Experiment”

- Scientific Motivation
- Accelerator and Laser
- Particle Detection and Simulation Results
- Conclusions



Letter of Intent for the LUXE Experiment

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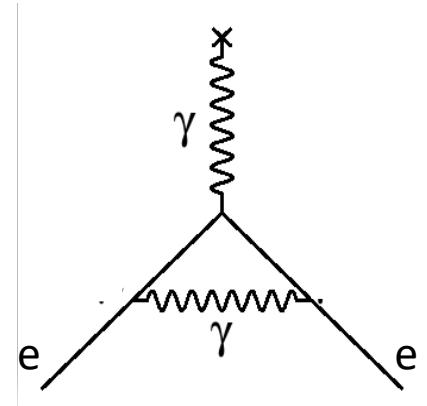
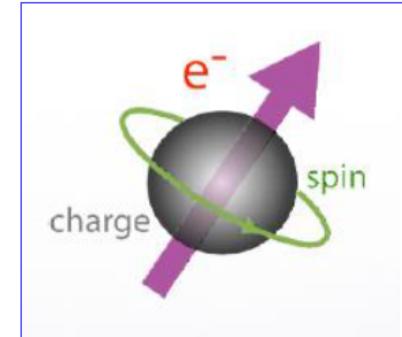
arXiv:1909.00860



SCIENTIFIC MOTIVATION

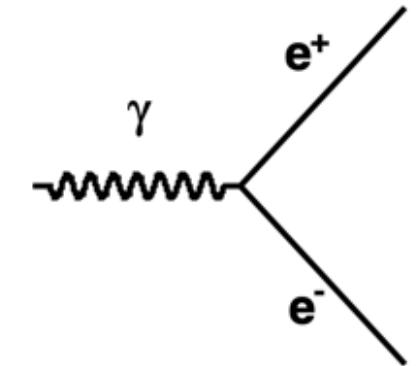
REMINDER: QUANTUM ELECTRODYNAMICS

- **Relativistic field theory of electrodynamics**
 - Perturbation theory in terms of coupling constant α
- **World's most precisely tested theory**
 - Anomalous magnetic dipole moment ($g-2$) of electron:
 - Zero at leading order => first corrections calculated by Schwinger (1947)
 - Based on precise measured and calculated (includes terms of 5th order: α^5) values, extract $1/\alpha = 137.035\ 999\ 070\ (98)$
 - Precision better than 10^{-9} , consistent with other measurements
 - Anomalous magnetic dipole moment of muon shows interesting tension
 - New experiment at FNAL ("Muon g-2") will improve precision by factor 4



QED: WHAT DO WE NOT KNOW?

- What happens if electrons or photons propagate in a very strong field?
 - QED expects that vacuum becomes unstable e.g. for nucleus with $Z > 137$. Spontaneous creation of e^+e^- pairs (“boiling of vacuum”)
- Historical developments:
 - 1930s: Initial discussions of EM in strong field in literature (Sauter, Euler, Heisenberg) => introduction of “critical field”
 - 1951: First non-perturbative calculations by Julian Schwinger
 - 1990s: E144 experiment at SLAC



$$\varepsilon_{crit} = \frac{m_e^2 c^3}{\hbar e} \simeq 1.3 \cdot 10^{18} \text{ V/m}$$



HEISENBERG AND EULER: THE CRITICAL FIELD



Folgerungen aus der Diracschen Theorie des Positrons.

Von W. Heisenberg und H. Euler in Leipzig.

Mit 2 Abbildungen. (Eingegangen am 22. Dezember 1935.)

Aus der Diracschen Theorie des Positrons folgt, da jedes elektromagnetische Feld zur Paarerzeugung neigt, eine Abänderung der Maxwellschen Gleichungen des Vakuums. Diese Abänderungen werden für den speziellen Fall berechnet, in dem keine wirklichen Elektronen und Positronen vorhanden sind, und in dem sich das Feld auf Strecken der Compton-Wellenlänge nur wenig ändert. Es ergibt sich für das Feld eine Lagrange-Funktion:

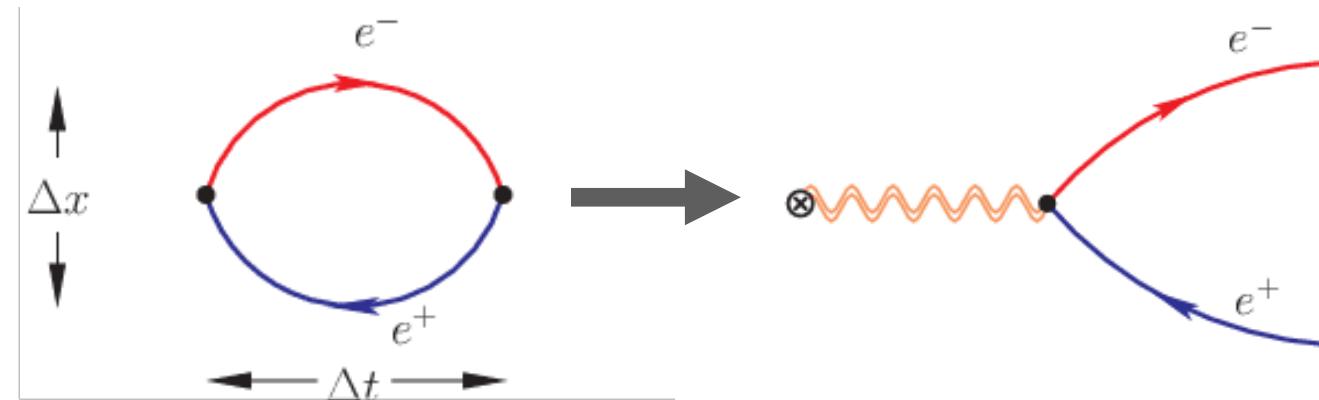
$$\mathfrak{L} = \frac{1}{2} (\mathfrak{E}^2 - \mathfrak{B}^2) + \frac{e^2}{\hbar c} \int_0^\infty e^{-\eta} \frac{d\eta}{\eta^3} \left\{ i\eta^2 (\mathfrak{E}\mathfrak{B}) \cdot \frac{\cos\left(\frac{\eta}{|\mathfrak{E}_k|} \sqrt{\mathfrak{E}^2 - \mathfrak{B}^2 + 2i(\mathfrak{E}\mathfrak{B})}\right) + \text{konj}}{\cos\left(\frac{\eta}{|\mathfrak{E}_k|} \sqrt{\mathfrak{E}^2 - \mathfrak{B}^2 + 2i(\mathfrak{E}\mathfrak{B})}\right) - \text{konj}} + |\mathfrak{E}_k|^2 + \frac{\eta^2}{3} (\mathfrak{B}^2 - \mathfrak{E}^2) \right\}.$$

$\mathfrak{E}, \mathfrak{B}$ Kraft auf das Elektron.

$$\left(|\mathfrak{E}_k| = \frac{m^2 c^3}{e \hbar} = \frac{1}{137^a} \frac{e}{(e^2/m c^2)^2} = \text{"Kritische Feldstärke".} \right)$$

THE SCHWINGER PROCESS

J. Schwinger: *On Gauge Invariance and Vacuum Polarization,*
Phys. Rev. 82 (1951) 664



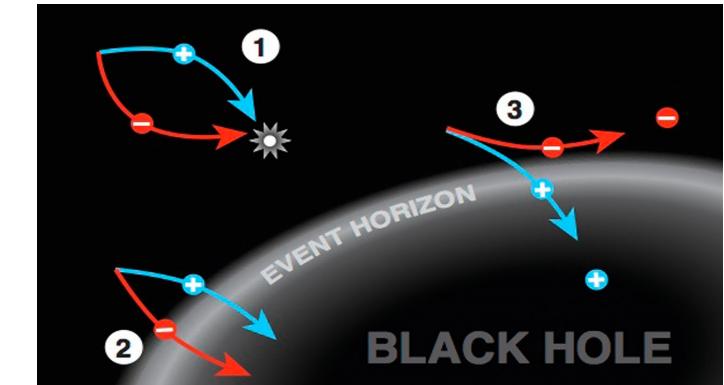
Photon in electric field: simplified

- The EM force is $F = e\varepsilon$
- Energy needed to separate e^+e^- pair: $E = Fd_{min}$
- Heisenberg: $\Delta t \geq \frac{\hbar}{2\Delta E} \Rightarrow \Delta t_{min} = \frac{\hbar}{4mc^2} \Rightarrow$ minimum distance: $d_{min} = 2c\Delta t_{min} = \frac{\hbar}{2mc} = \lambda_c/2$
- Virtual pair becomes real if $E = Fd_{min} = \frac{\hbar e \varepsilon}{2mc} > 2mc^2 \Rightarrow$ possible if $\varepsilon > \frac{4m^2 c^3}{\hbar e} = 4\varepsilon_{crit}$

$$P \propto \exp\left(-\frac{d}{\lambda_c}\right) \propto \exp\left(-\pi \frac{\varepsilon_{crit}}{\varepsilon}\right)$$

ANALOGY TO HAWKING RADIATION

- Energy needed to create on-shell e^+e^- pair: $\Delta E = 2mc^2$
- Grav. Field near the event horizon: $F = \frac{G_N M m}{r_s^2}$
- Schwarzschild radius $r_s = \frac{2G_N M}{c^2}$. $\Rightarrow F = \frac{mc^4}{4G_N M}$
- Energy to separate pair: $E = F d_{min} = \frac{mc^4}{4G_N M} \times \frac{\hbar}{mc} = \frac{\hbar c^3}{4G_N M}$



H. Murayama

Hawking radiation possible if virtual pair becomes real, i.e. $\frac{\hbar c^3}{4G_N M} > 2mc^2$



WHY EXPLORE STRONG-FIELD QED?

- Relevant to numerous phenomena in our Universe

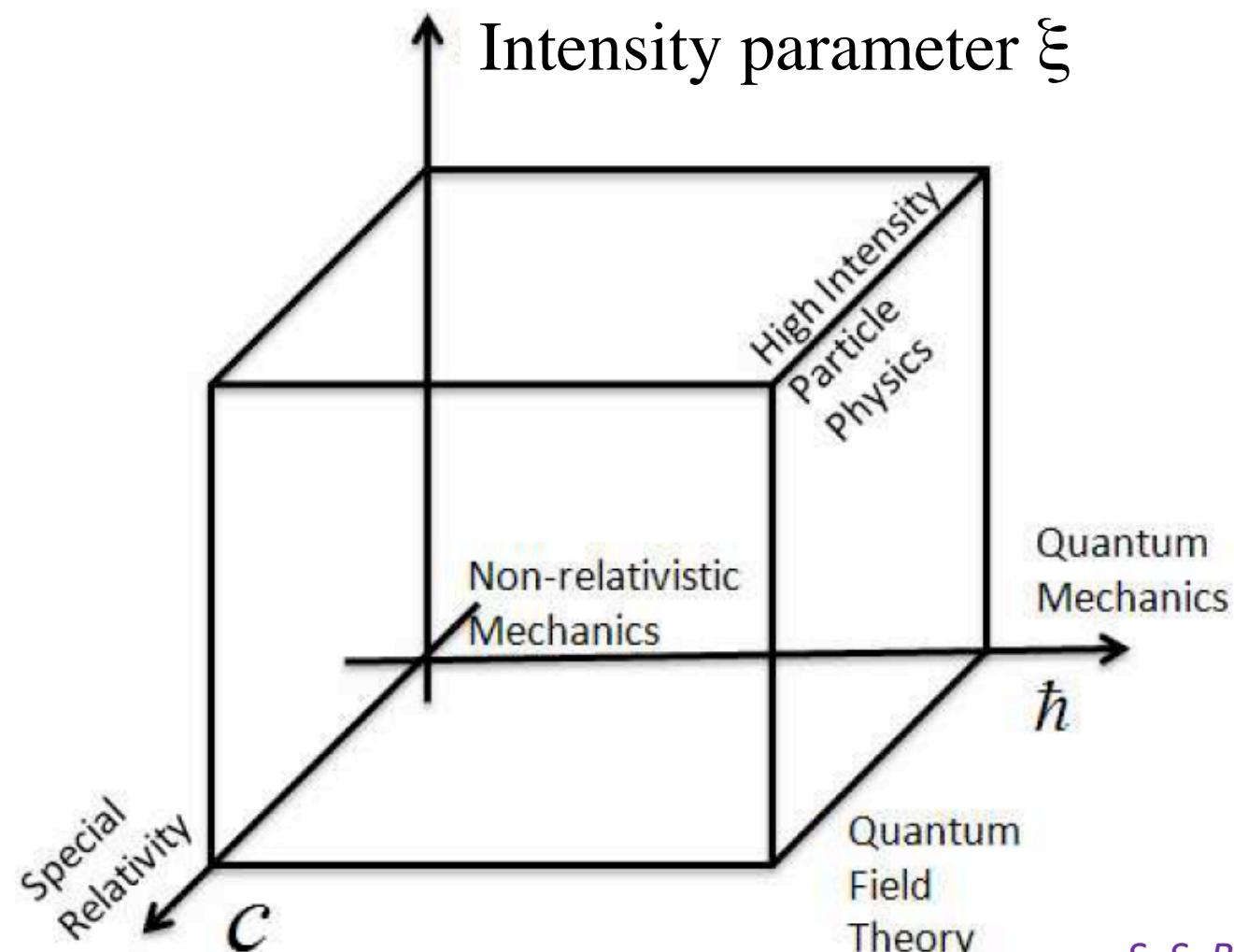
- Astrophysics:

- Hawking radiation, surface of neutron stars (magnetars), early Universe
- Condensed matter and atomic physics (nuclei with $Z>137$)
- Accelerator physics: high energy e^+e^- colliders

- Main goals:

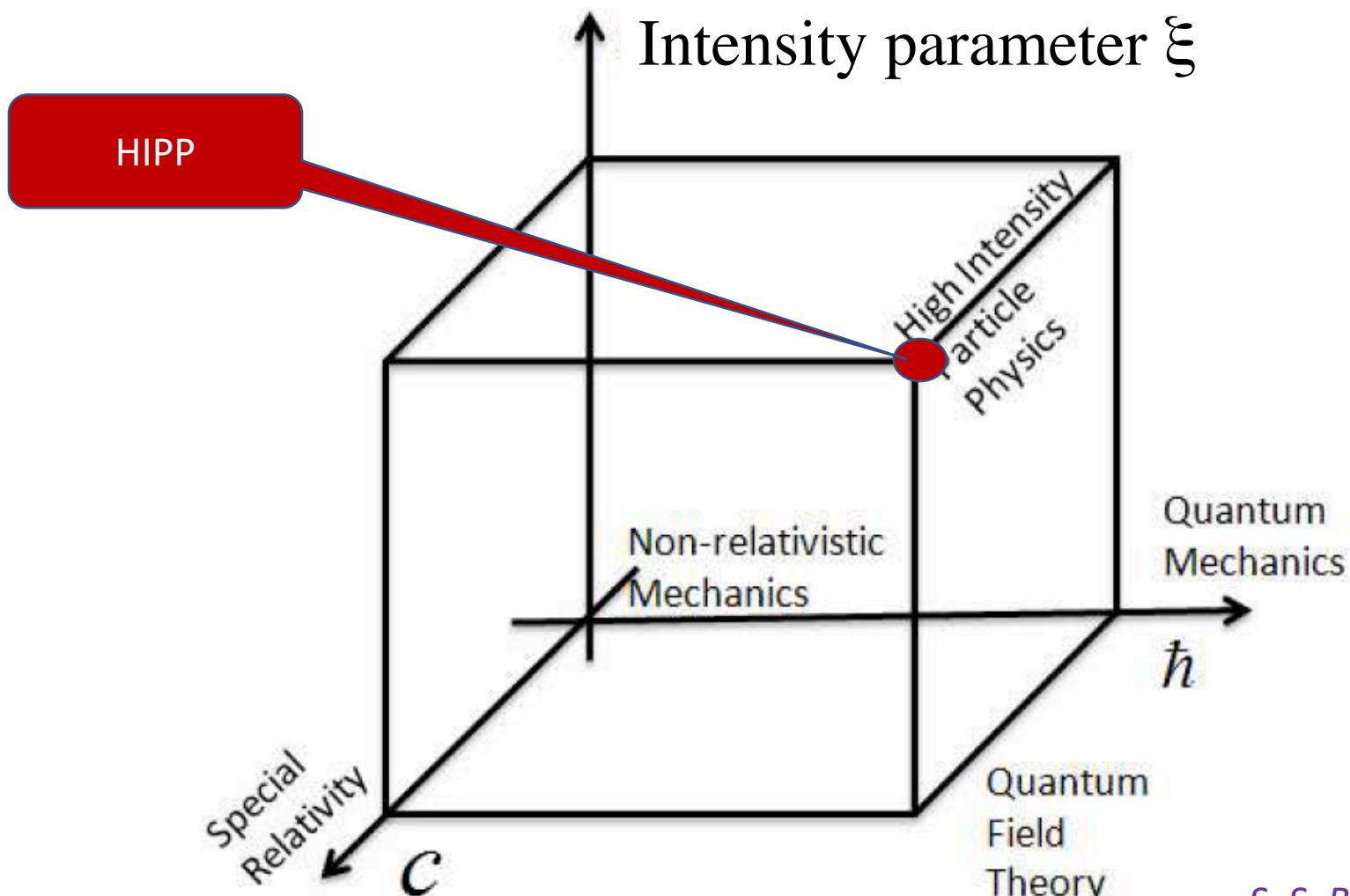
- Testing theoretical predictions in novel regime
 - gain deeper understanding of quantum physics
- Measure transition from perturbative to non-perturbative regime
 - could teach us about other non-perturbative regimes, e.g. understanding confinement [Gribov, hep-ph/9902279]
- **Schwinger field has never been reached experimentally in clean environment**
 - Exciting to be the first to explore this ... we might be surprised what we find!

THEORIES ON A CUBE



S. S. Bulanov, W. Leemans et al.

THEORIES ON A CUBE



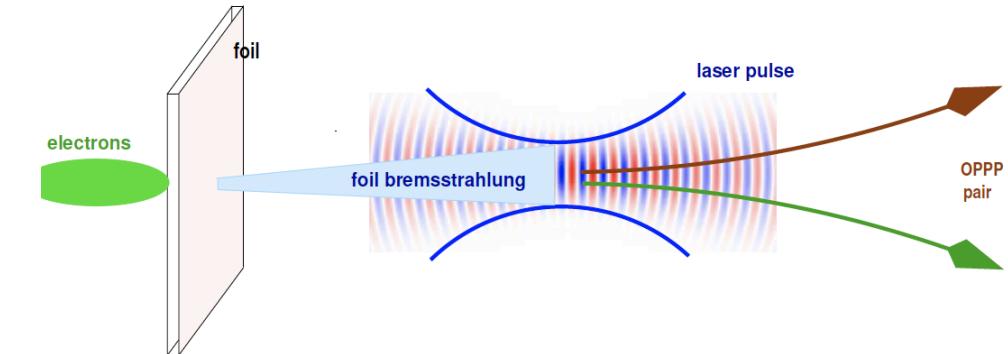
S. S. Bulanov, W. Leemans et al.

LASER AND PHOTON BEAM

- Use laser to generate electric field
- Use high energy electron beam

$$\xi = \frac{e\varepsilon_L}{m_e \omega_L c}$$

$$\chi \approx \gamma \frac{\varepsilon_L}{\varepsilon_{crit}}$$



- Laser power required to reach **Schwinger field** ($\chi_\gamma \sim 1$):

- Non-relativistic photons: $I=2 \times 10^{29} \text{ W/cm}^2$
- EU.XFEL, $E_\gamma \approx 10 \text{ GeV}$: $\approx 10^{20} \text{ W/cm}^2$
- ELI-NP, $E_\gamma \approx 1 \text{ GeV}$: $\approx 10^{22} \text{ W/cm}^2$

=> Much beyond currently achievable values

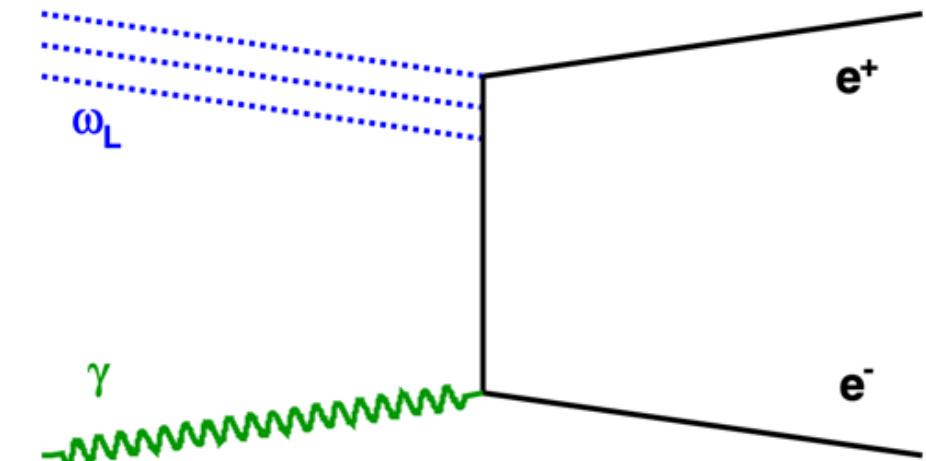
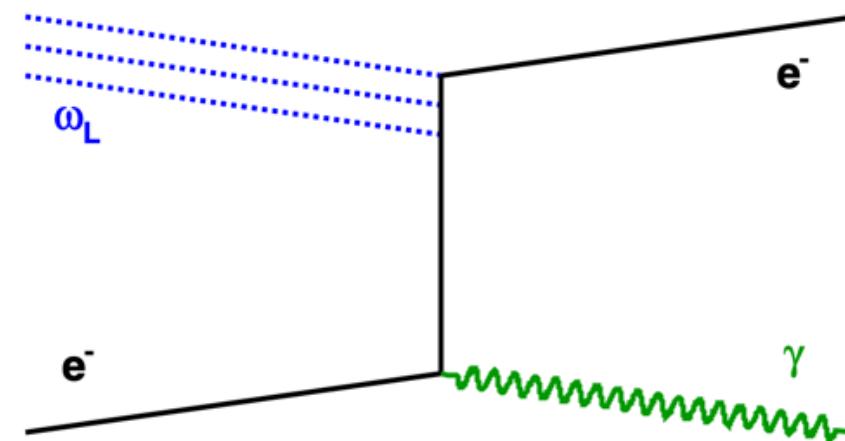
=> Can use well-tested laser technology

=> State-of-the-art laser needed

MAIN PROCESSES OF INTEREST



Low-energy photons
from laser

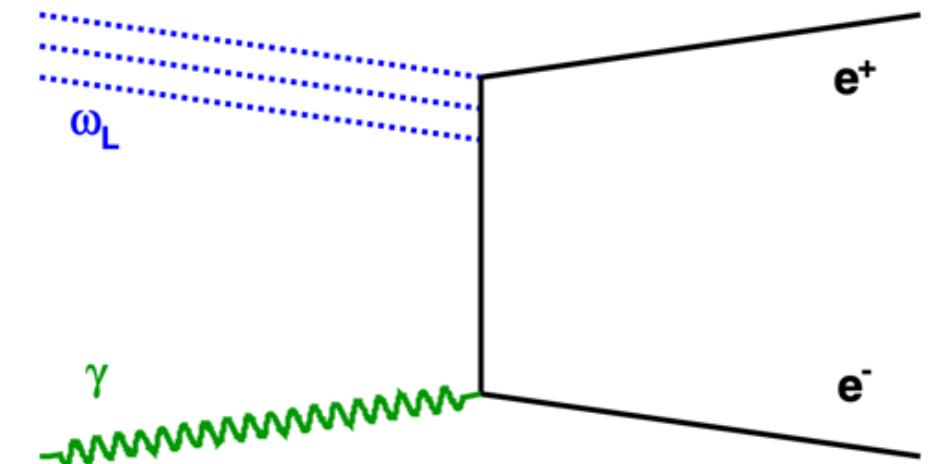


High energy electron or photon interacts with laser

- Also higher order process $e^- + n\omega_L \rightarrow e^- e^+ e^-$
- Via two steps ($e^- + n\omega_L \rightarrow e^- + \gamma$ and then $\gamma + n\omega_L \rightarrow e^+ e^-$) or one step

CROSS SECTION OF QED PROCESSES

- Perturbative QED valid
 - For n photons $\sigma \propto \alpha^n$
 - With $\alpha \propto e^2 \propto \xi^2$ it follows: $\sigma \propto \xi^{2n}$
- If $\xi \gtrsim 1$ all orders can contribute ~equally => cannot truncate series any more
 - All-order calculation needs to be performed (which is hard)
- Example for asymptotic result for $\xi \gg 1$ and $\chi < 1$: $\sigma \propto \chi e^{-8/(3\chi)}$
 - Since $\chi \propto \sqrt{\alpha}$ cannot expand perturbatively
 - Result not proportional to powers of α



**Observation of deviation from power-law is
the experimental signature of strong QED**

PAIR PRODUCTION PROCESS

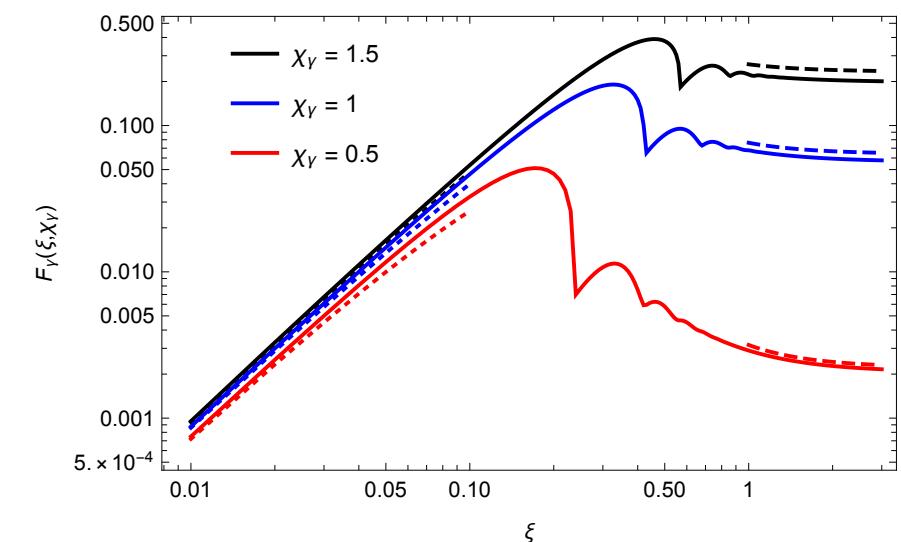
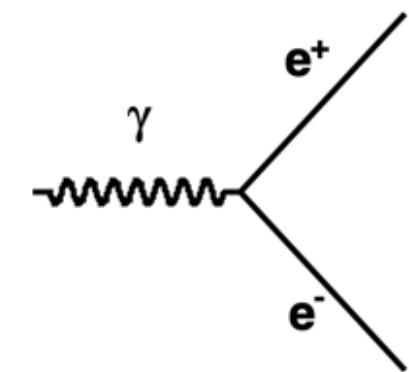
- Process not possible in vacuum in classical electrodynamics
- Pair production in a constant static field (Schwinger process)

$$\frac{\Gamma_{\text{SPP}}}{V} = \frac{m_e^4}{(2\pi)^3} \left(\frac{|\mathbf{E}|}{E_c} \right)^2 \sum_{n=1}^{\infty} \frac{1}{n^2} \exp \left(-n\pi \frac{E_c}{|\mathbf{E}|} \right) \propto \exp \left(-\pi \frac{E_c}{|\mathbf{E}|} \right)$$

- Pair production in plane wave laser: asymptotic result

$$\Gamma_{\text{OPPP}} \rightarrow \frac{3}{16} \sqrt{\frac{3}{2}} \alpha m_e (1 + \cos \theta) \frac{|\mathbf{E}|}{E_c} \exp \left[-\frac{8}{3} \frac{1}{1 + \cos \theta} \frac{m_e}{\omega_i} \frac{E_c}{|\mathbf{E}|} \right]$$

- Good agreement between full calculation and asymptotic result for $\xi \ll 1$ and $\xi > 1$



PAIR PRODUCTION PROCESS

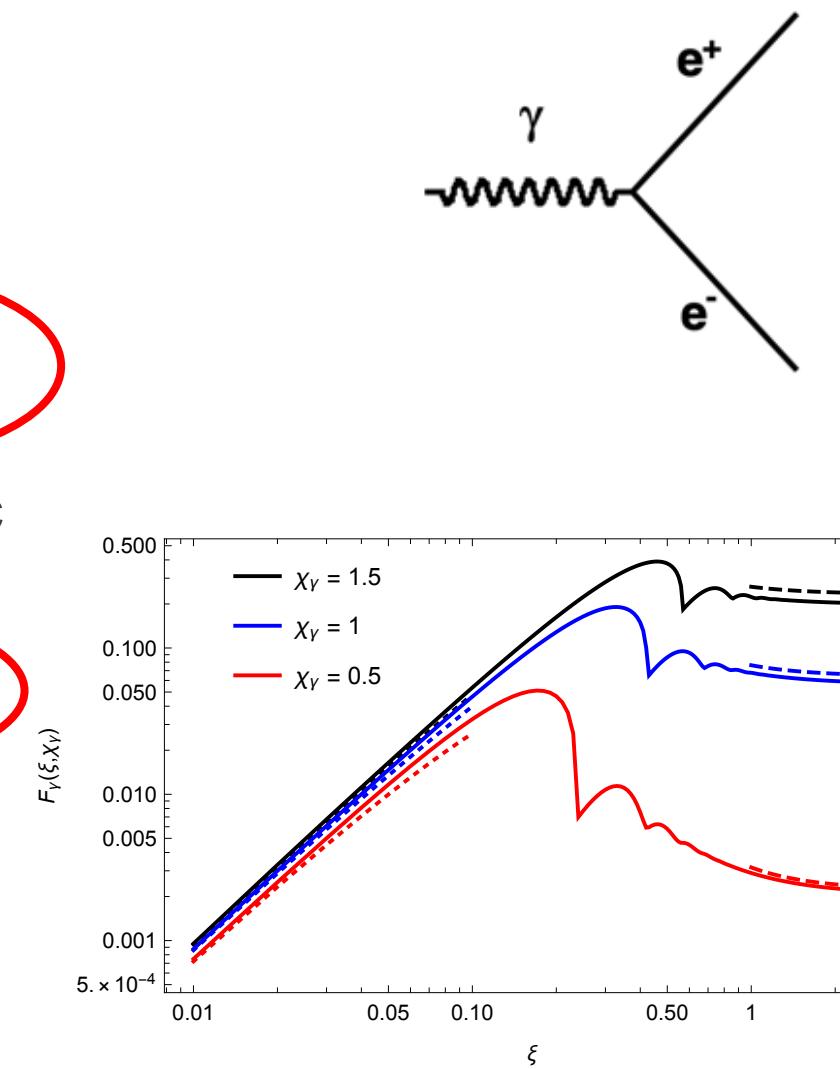
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- Pair production in plane wave laser: asymptotic result

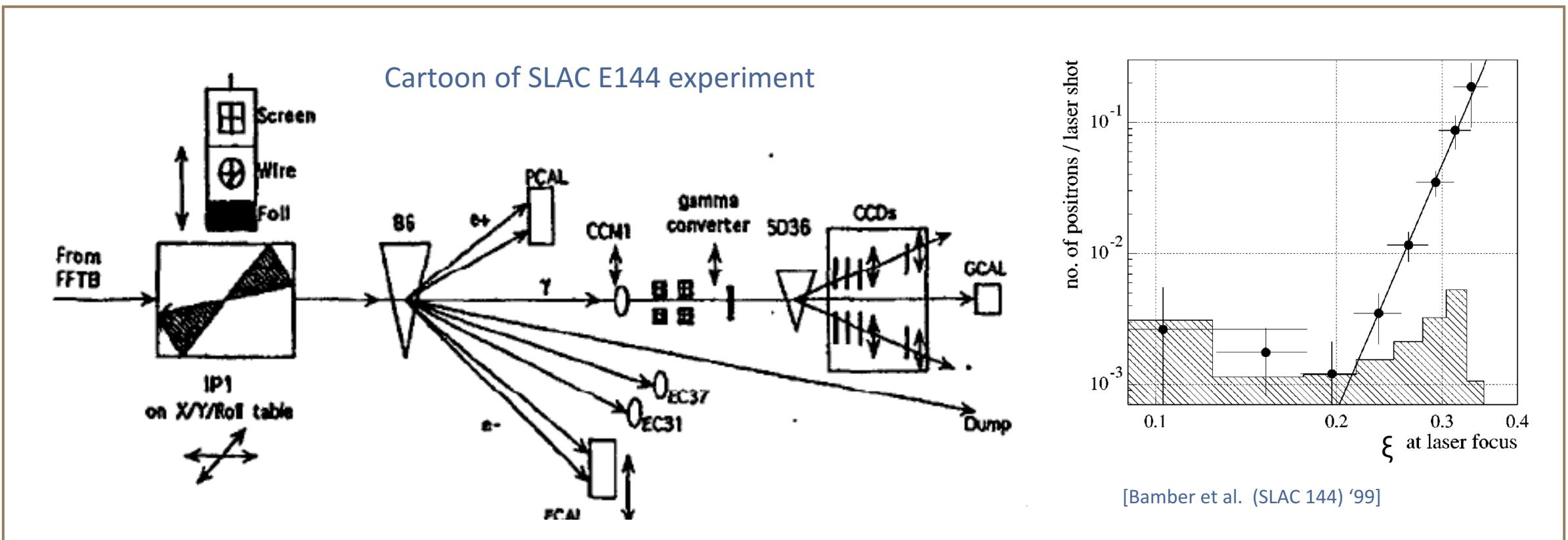
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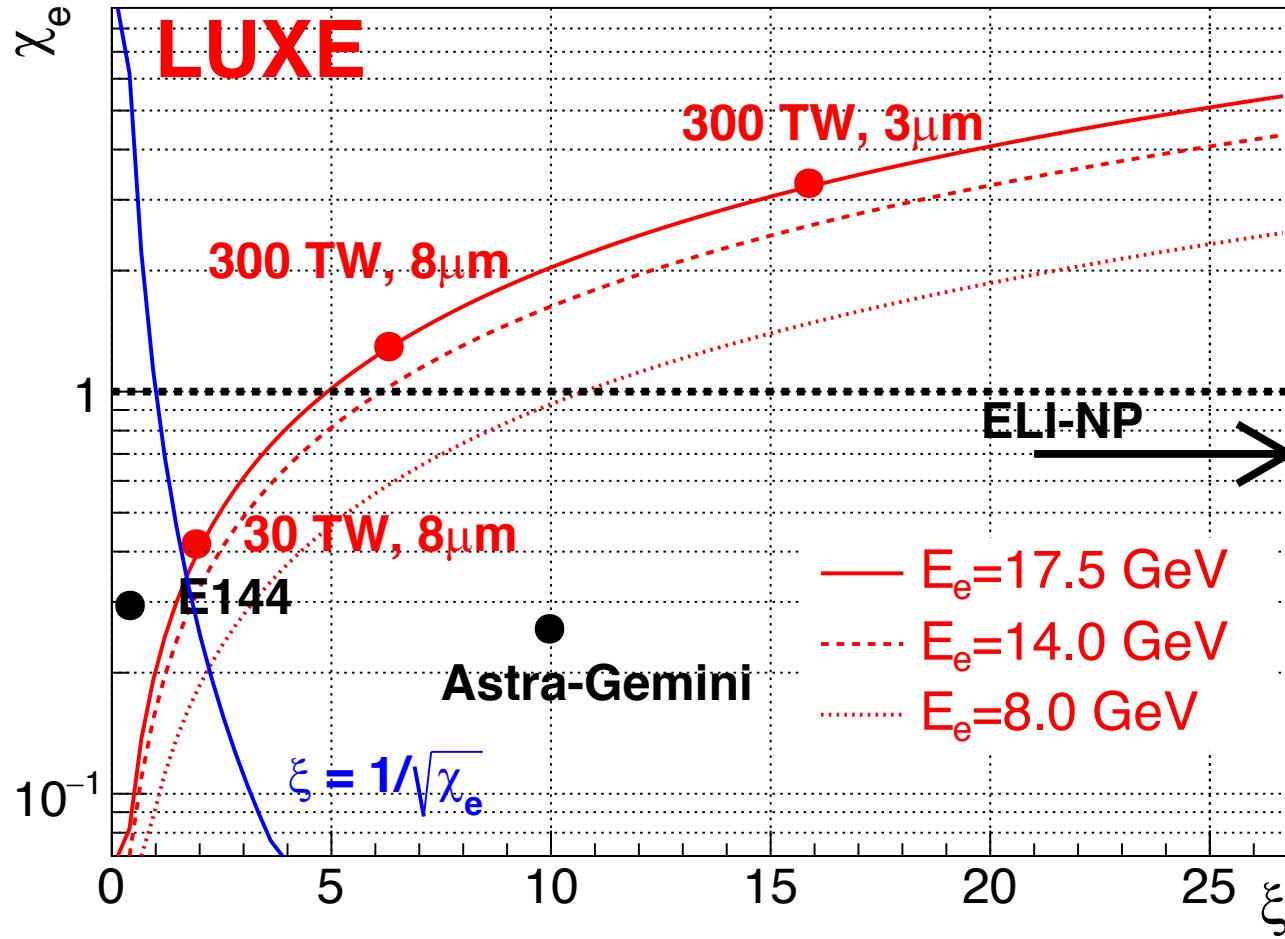


EXPERIMENT E144 AT SLAC

- Experiment at SLAC in 1990s with $E_{\text{beam}} = 46.6 \text{ GeV}$ achieved $\chi \leq 0.25$
 - Did observe two-step process $e^- + n\omega_L \rightarrow e^- e^+ e^-$
 - Saw the expected strong rise with ξ^{2n} but did not reach the critical field



PARAMETER SPACE



Intensity parameter:

$$\xi = \sqrt{4\pi\alpha} \left(\frac{\mathcal{E}_L}{\omega_L m_e} \right) = \frac{m_e \mathcal{E}_L}{\omega_L \mathcal{E}_{cr}}$$

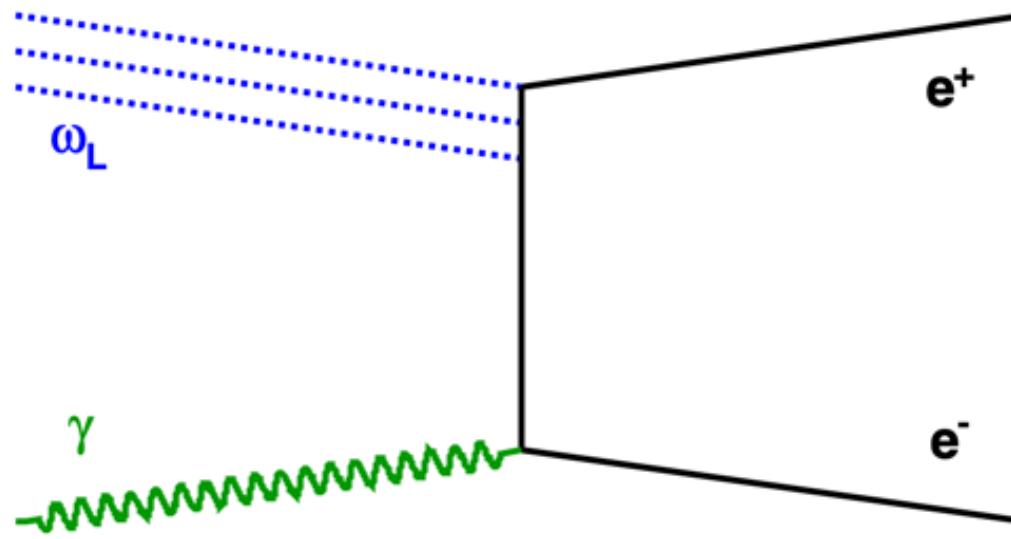
Quantum parameters:

$$\chi_e = (1 + \cos \theta) \frac{E_e}{m_e} \frac{\mathcal{E}_L}{\mathcal{E}_{cr}}$$

$$\chi_\gamma = (1 + \cos \theta) \frac{E_\gamma}{m_e} \frac{\mathcal{E}_L}{\mathcal{E}_{cr}}$$

ABSORBING LIGHT WITH LIGHT

Low-energy photons from laser

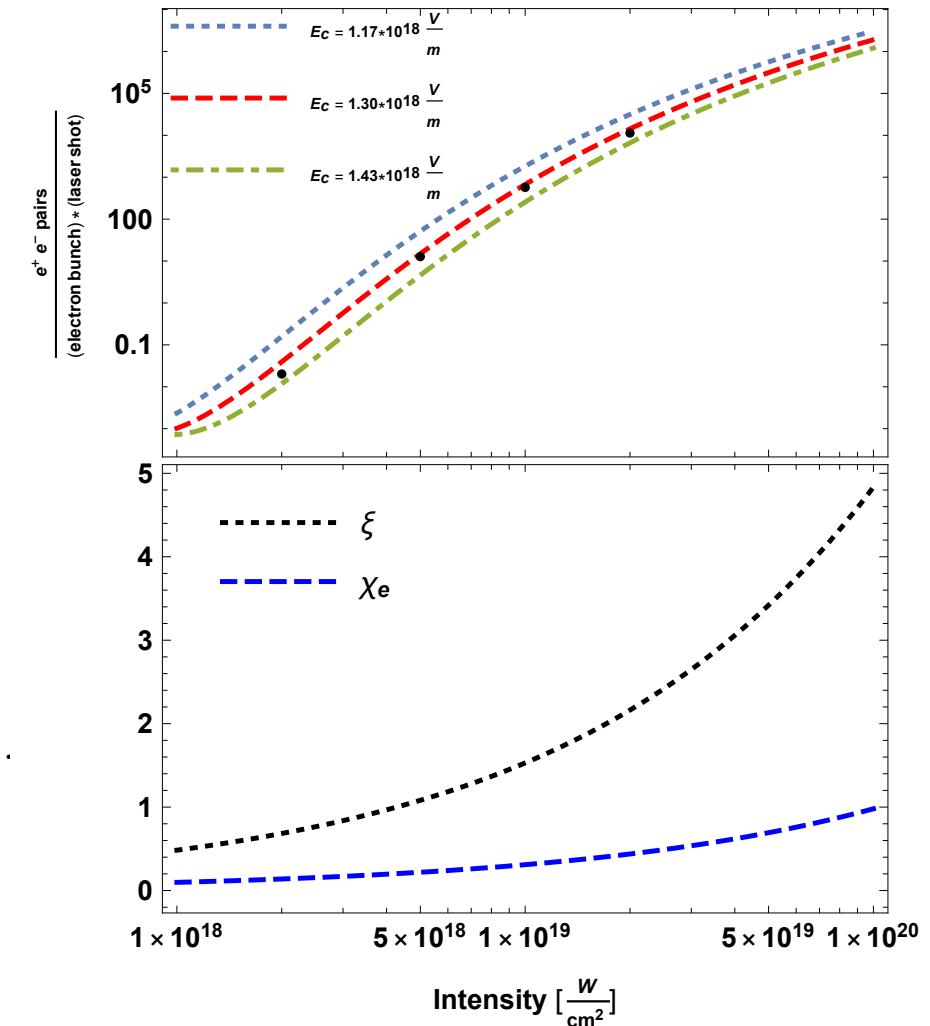


High-energy (relativistic) photon

$$\Gamma_{\text{BPPP}} \rightarrow \frac{9}{128} \sqrt{\frac{3}{2}} \alpha E_e (1 + \cos \theta)^2 \left(\frac{|\mathbf{E}|}{E_c} \right)^2 \exp \left[-\frac{8}{3} \frac{1}{1 + \cos \theta} \frac{m_e}{E_e} \frac{E_c}{|\mathbf{E}|} \right] \frac{X}{X_0} .$$

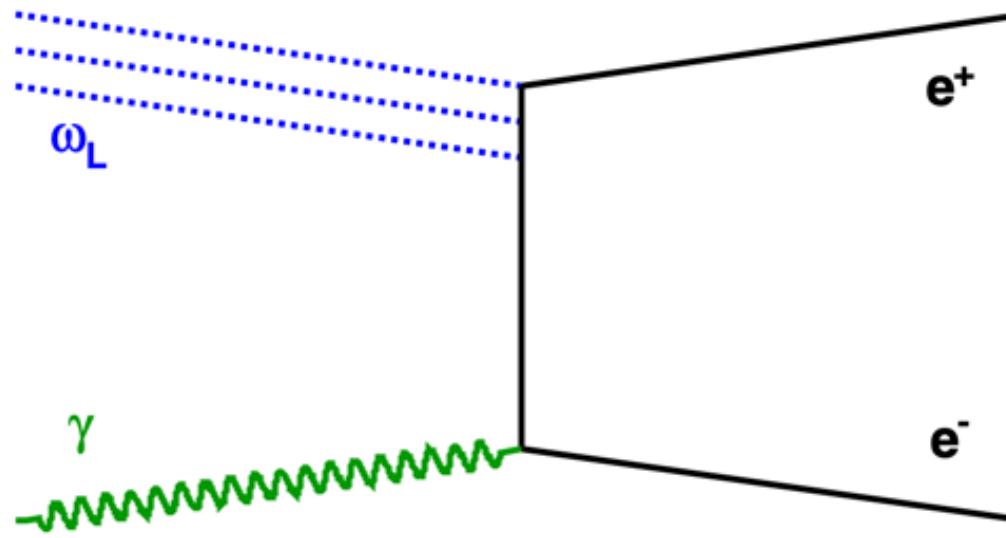
For $N(\omega_L) > 5$: $\sqrt{s} > 2mc^2$

$E_e = 17.5 \text{ GeV}$, $e^- b = 6 \times 10^9$, $\frac{X}{X_0} = 0.01$, L. s. = 35 fs, $\theta = \frac{\pi}{12}$, $w = 1.55 \text{ eV}$



ABSORBING LIGHT WITH LIGHT

Low-energy photons from laser



High-energy (relativistic) photon

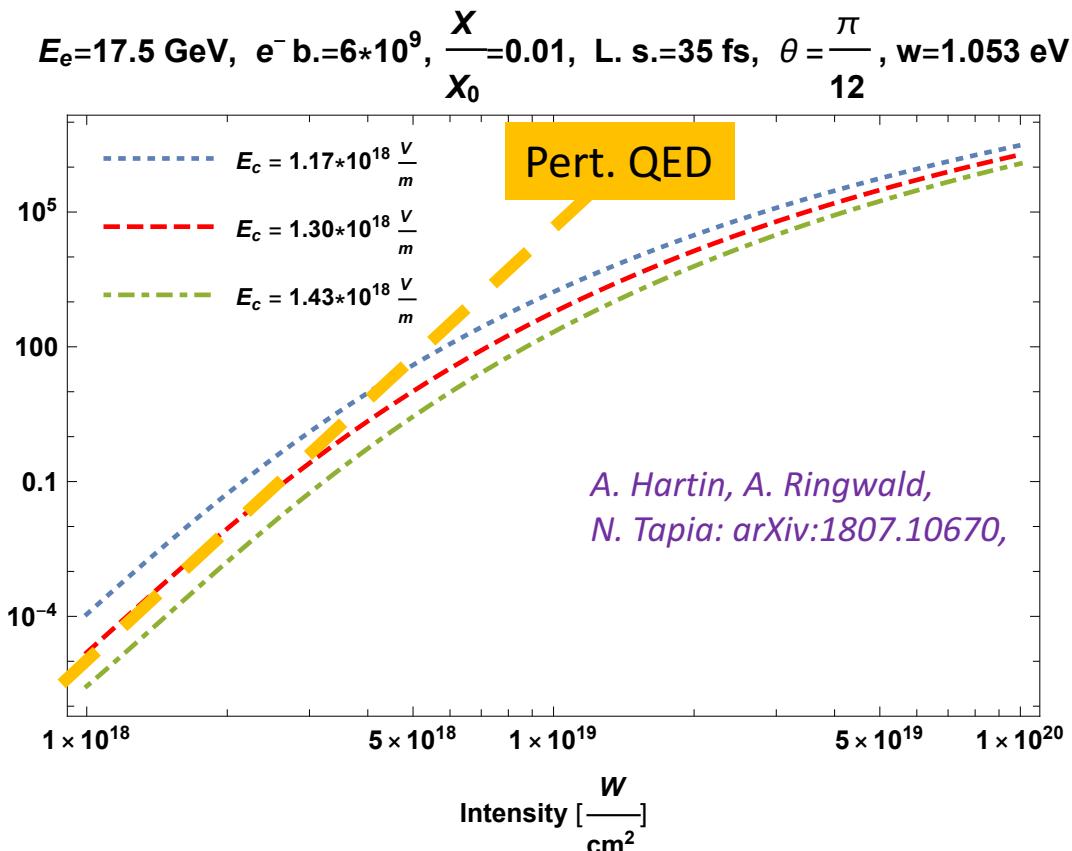
- Prediction for rate of positrons per laser shot

$$\xi \ll 1: R_{e^+} \propto \xi^{2n} \propto I^n$$

👉 Perturbative regime: strong rise, follows power-law

$$\xi \gg 1: R_{e^+} \propto \chi_\gamma \exp\left(-\frac{8}{3\chi_\gamma}\right)$$

👉 Non-perturbative regime: departure from power-law





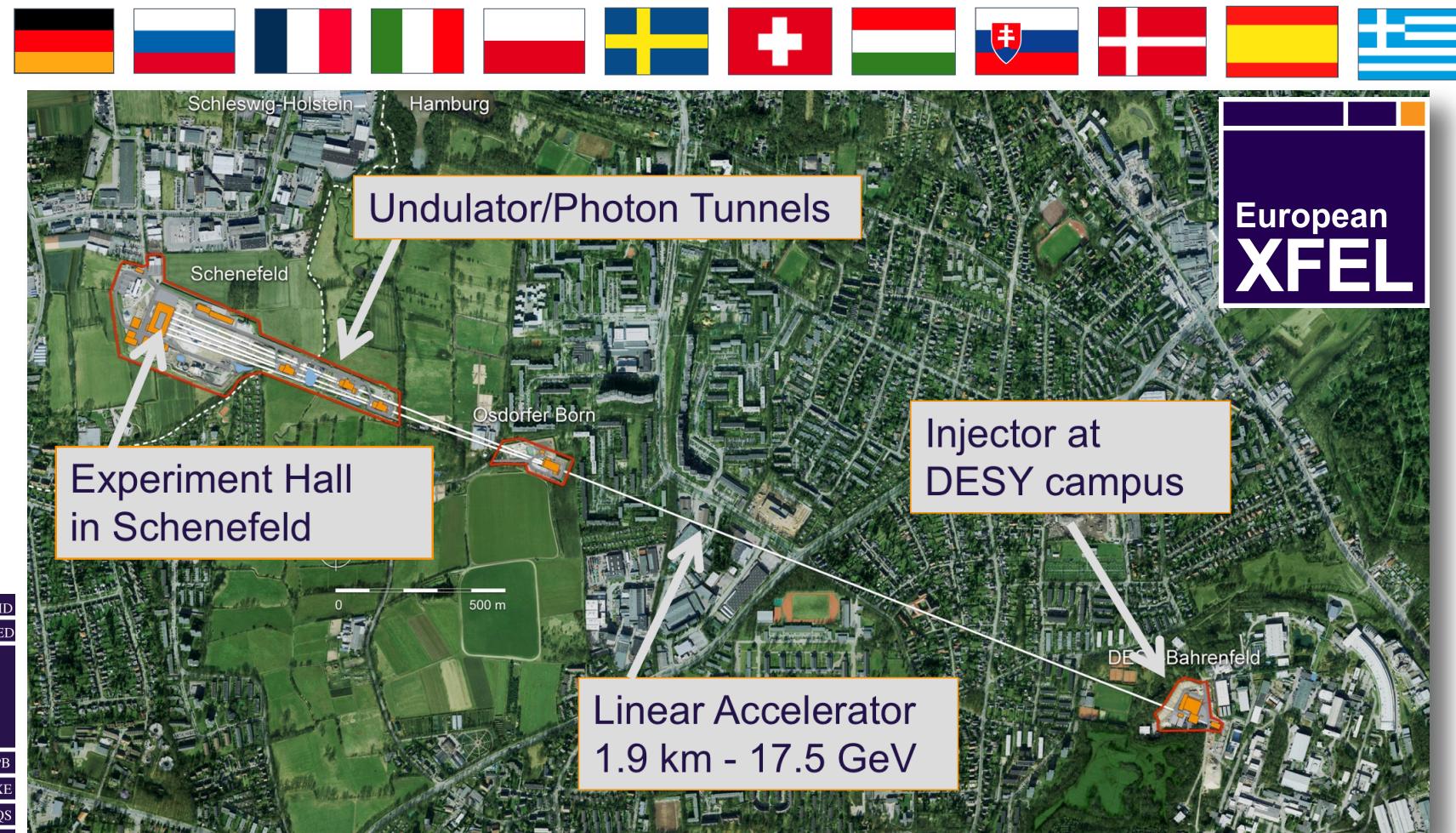
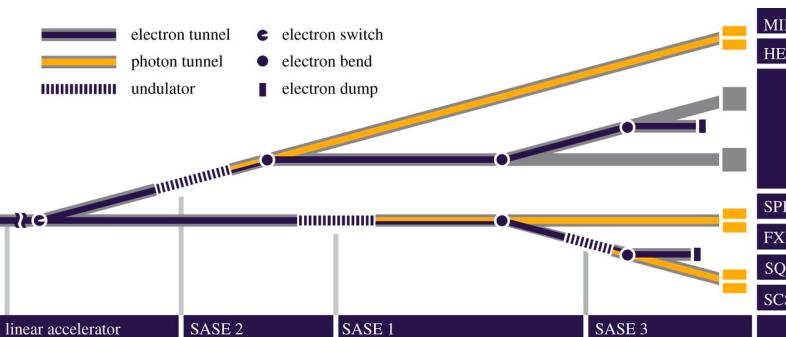
ACCELERATOR AND LASER



THE EUROPEAN XFEL

Electron accelerator:

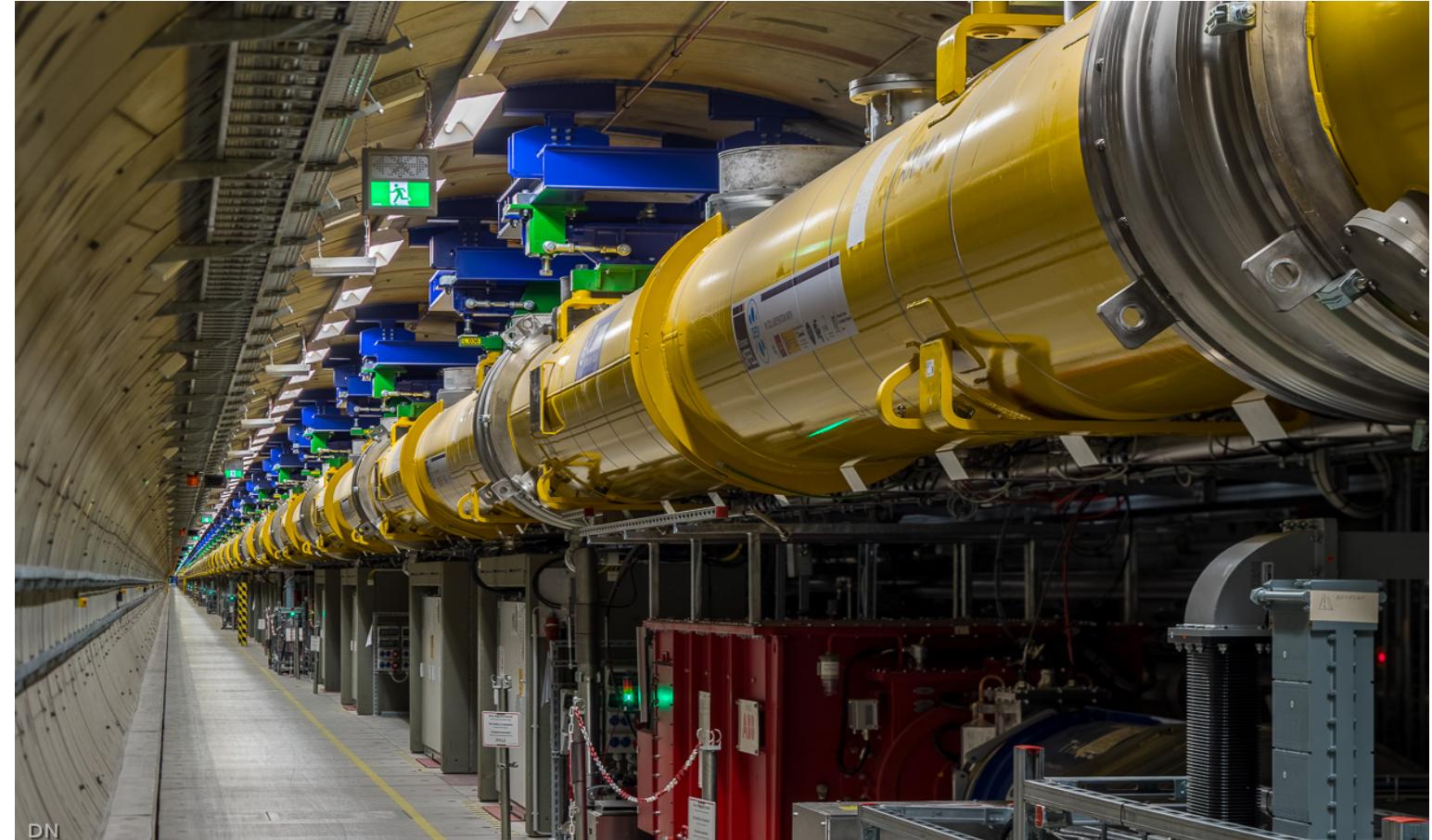
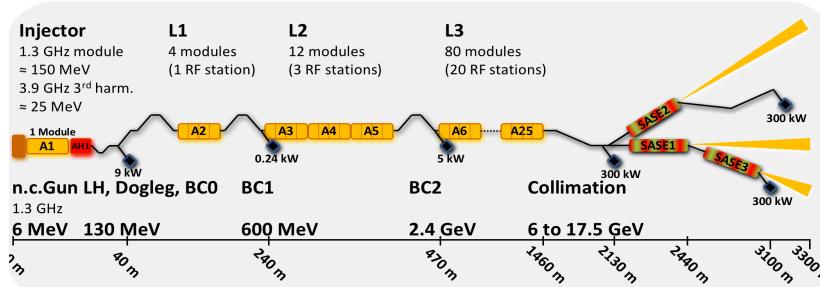
- 2.1 km 17.5 GeV SCRF linear accelerator
- 2700 electron bunches at rate of 10 Hz
- X-ray photons produced in undulators
- Experiments for physics, material science, chemistry, biology, ...





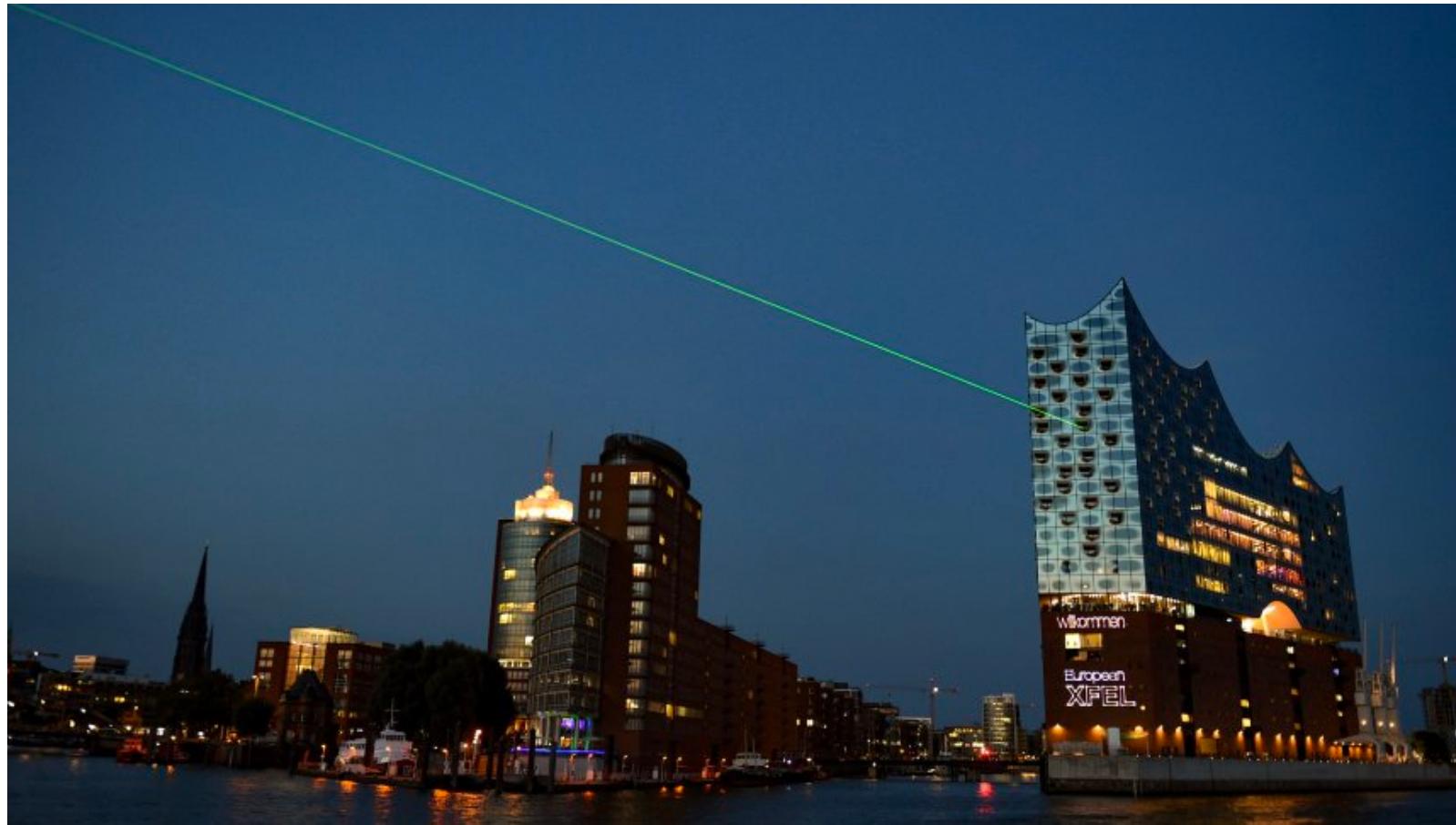
THE EUROPEAN XFEL

View along L3 accelerator section and undulator





EUROPEAN XFEL INAUGURATION



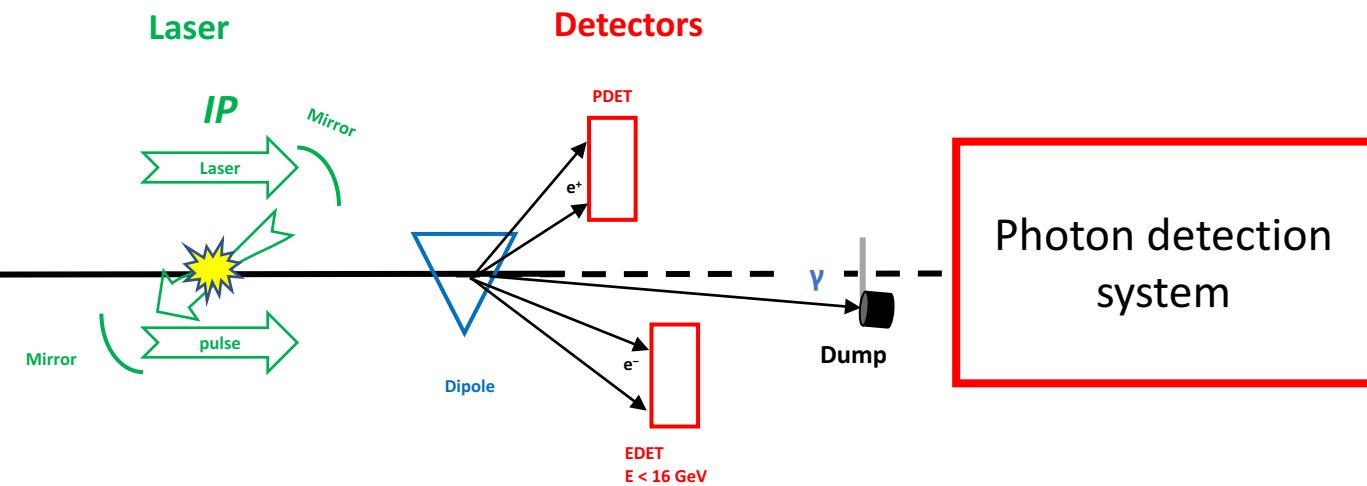
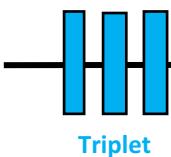
Operating since September 2016



ELECTRON LASER COLLISIONS

Compton and trident processes: $e^- + n\omega \rightarrow e^- + \gamma$ and $e^- + n\omega \rightarrow e^-e^+e^-$

one bunch of electron beam from XFEL-EU



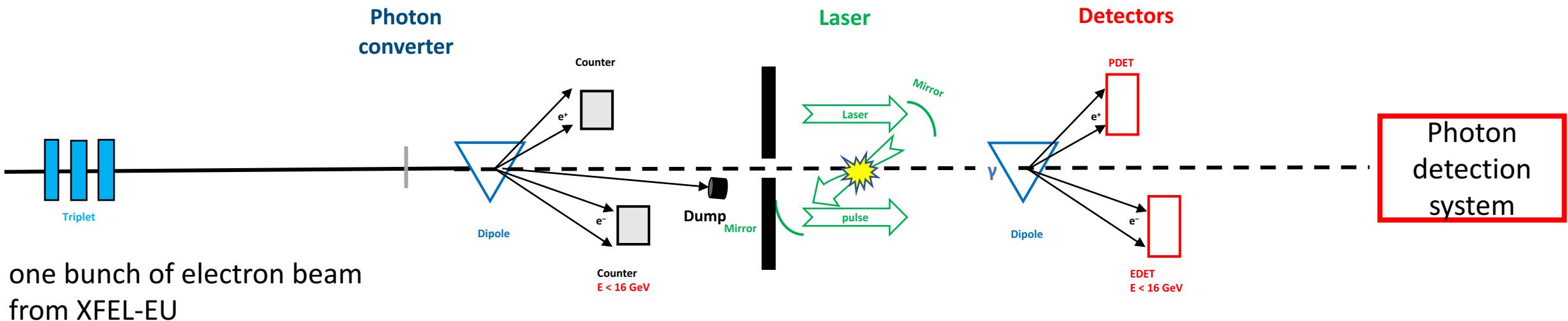
Kicker and triplet to select single bunch and focus it

Electron- Laser interaction area

Dipole and detectors to observe e^+e^- pairs

PHOTON LASER COLLISIONS

Pair production (Breit-Wheeler) process: $\gamma + n\omega \rightarrow e^- + e^+$



one bunch of electron beam
from XFEL-EU

Kicker and triplet to
select single bunch
and focus it

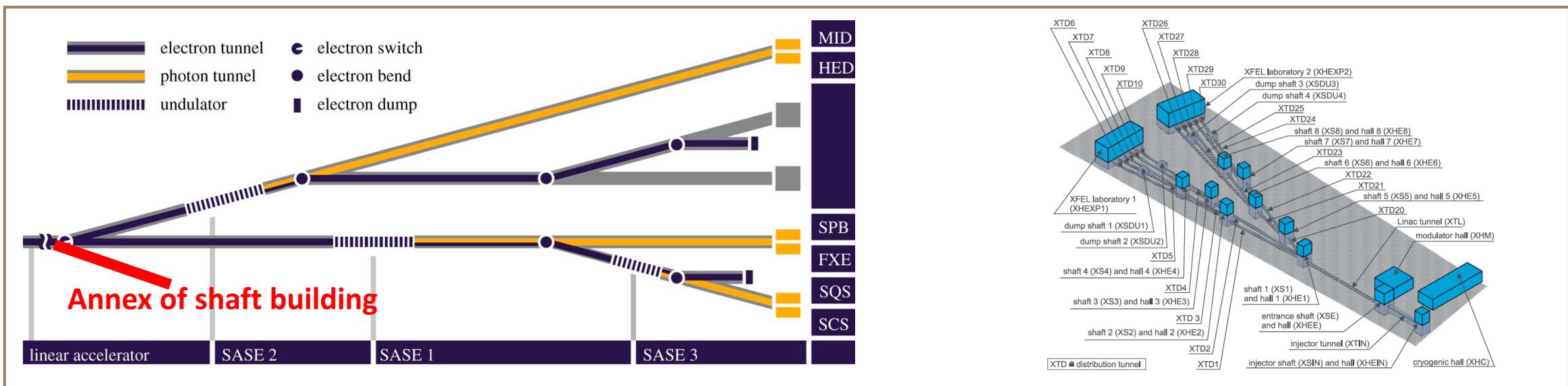
Dipole and detectors to
remove e^+e^- pairs and
monitor photon flux

Photon- Laser
interaction area

Dipole and
detectors to
observe e^+e^- pairs

LOCATIONS IN EU.XFEL TUNNEL

- Location at EU.XFEL:
 - Annex of shaft building XS1: at end of electron accelerator
 - Was build for 2nd EU.XFEL fan foreseen for later (late 2020s)
- Design aims to have no impact on photon science programme
 - Use only 1 of the 2700 bunches in bunch train (kicked out by fast kicker magnet)

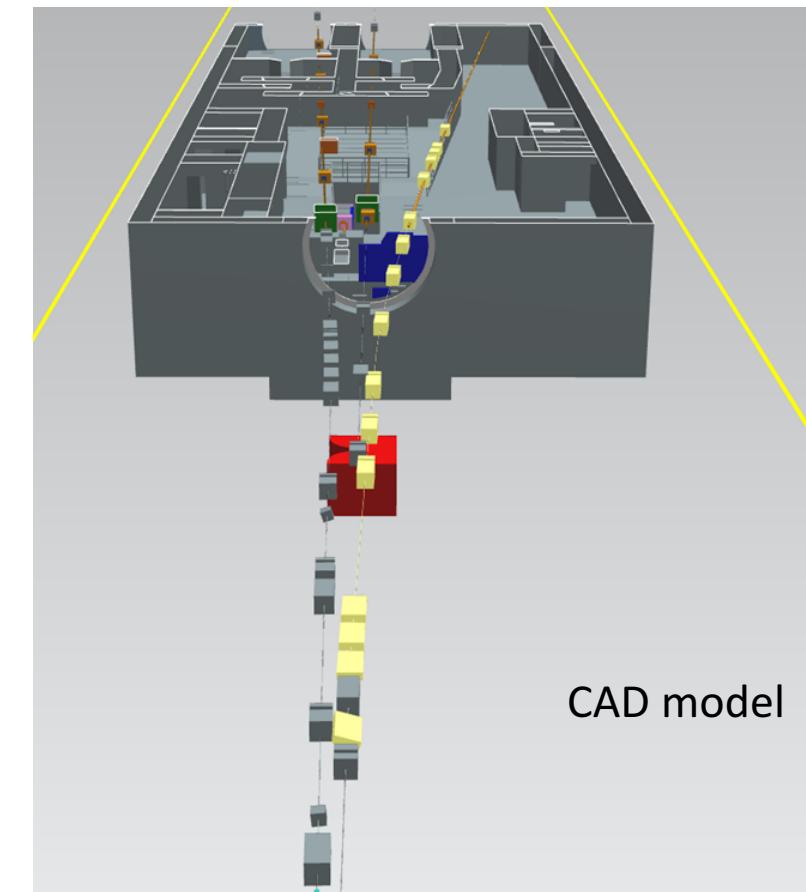
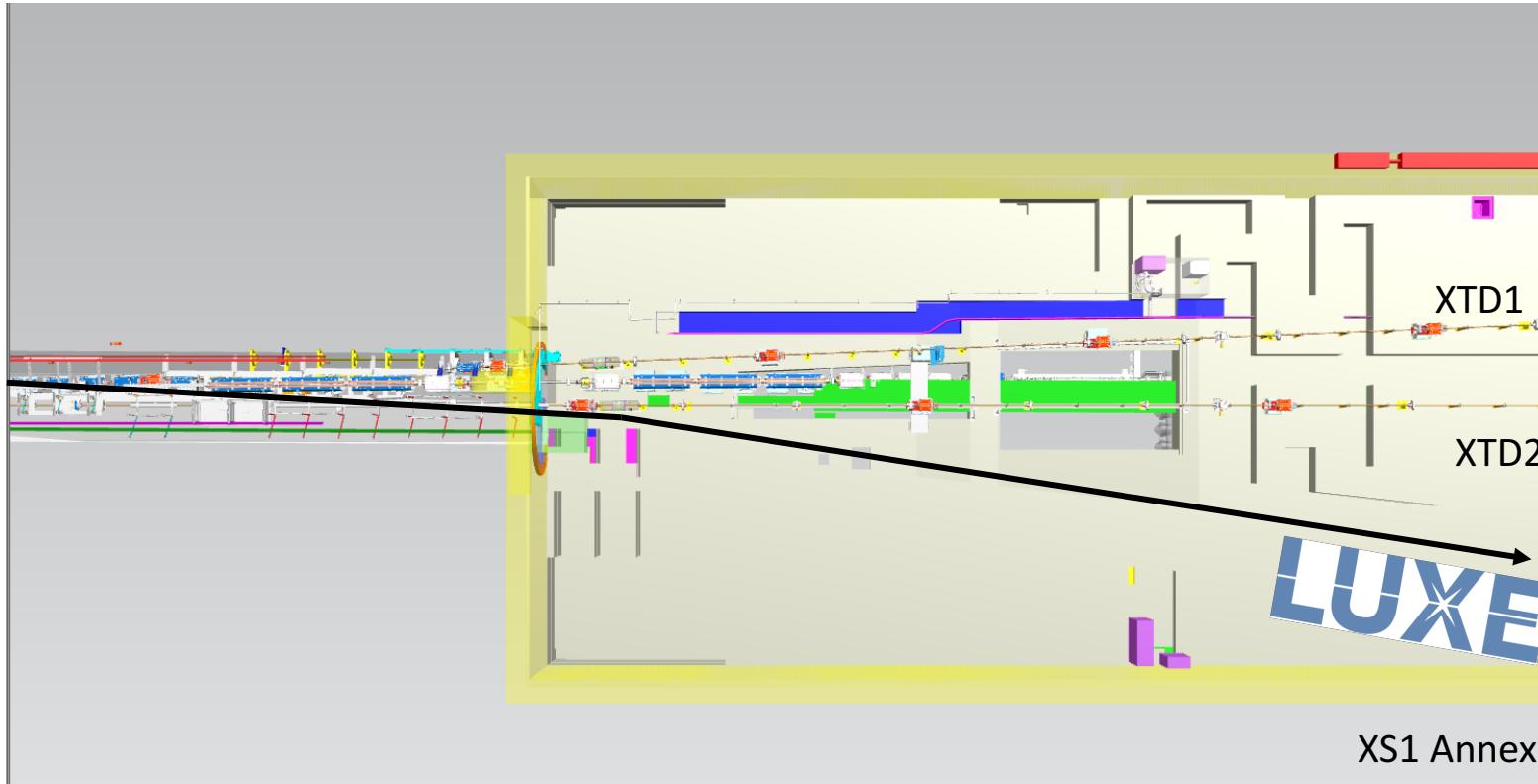


LOCATION





SCHEMATIC VIEW: BEAM EXTRACTION AND TRANSFER



M. Huening, M. Scheer, F. Burkart, W. Decking

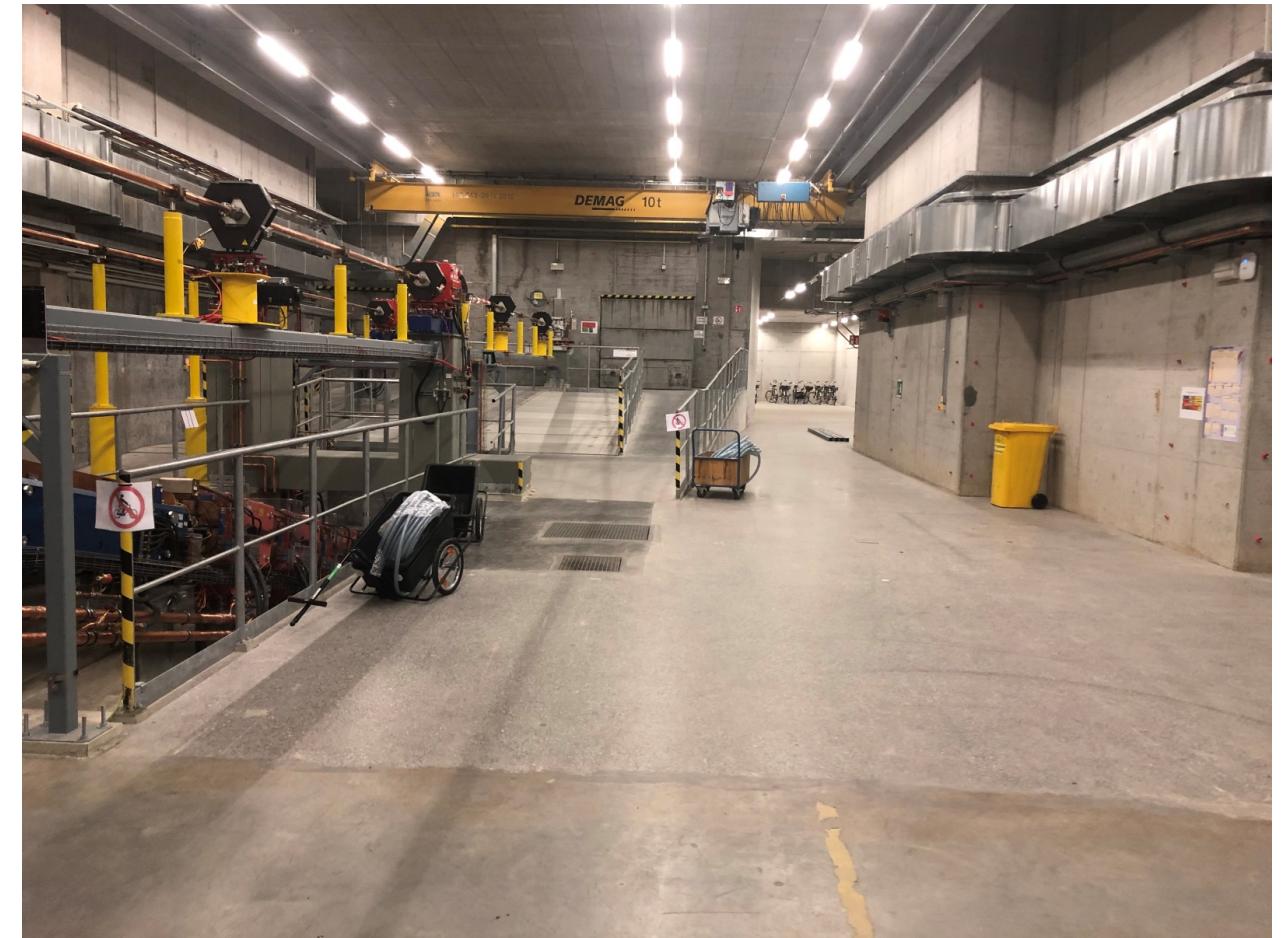


PICTURE OF TUNNEL AT XS1 ANNEX

Shaft located at end of linear accelerator of European XFEL

Dimensions of annex

- 60m long, 5.4m wide, 5m high



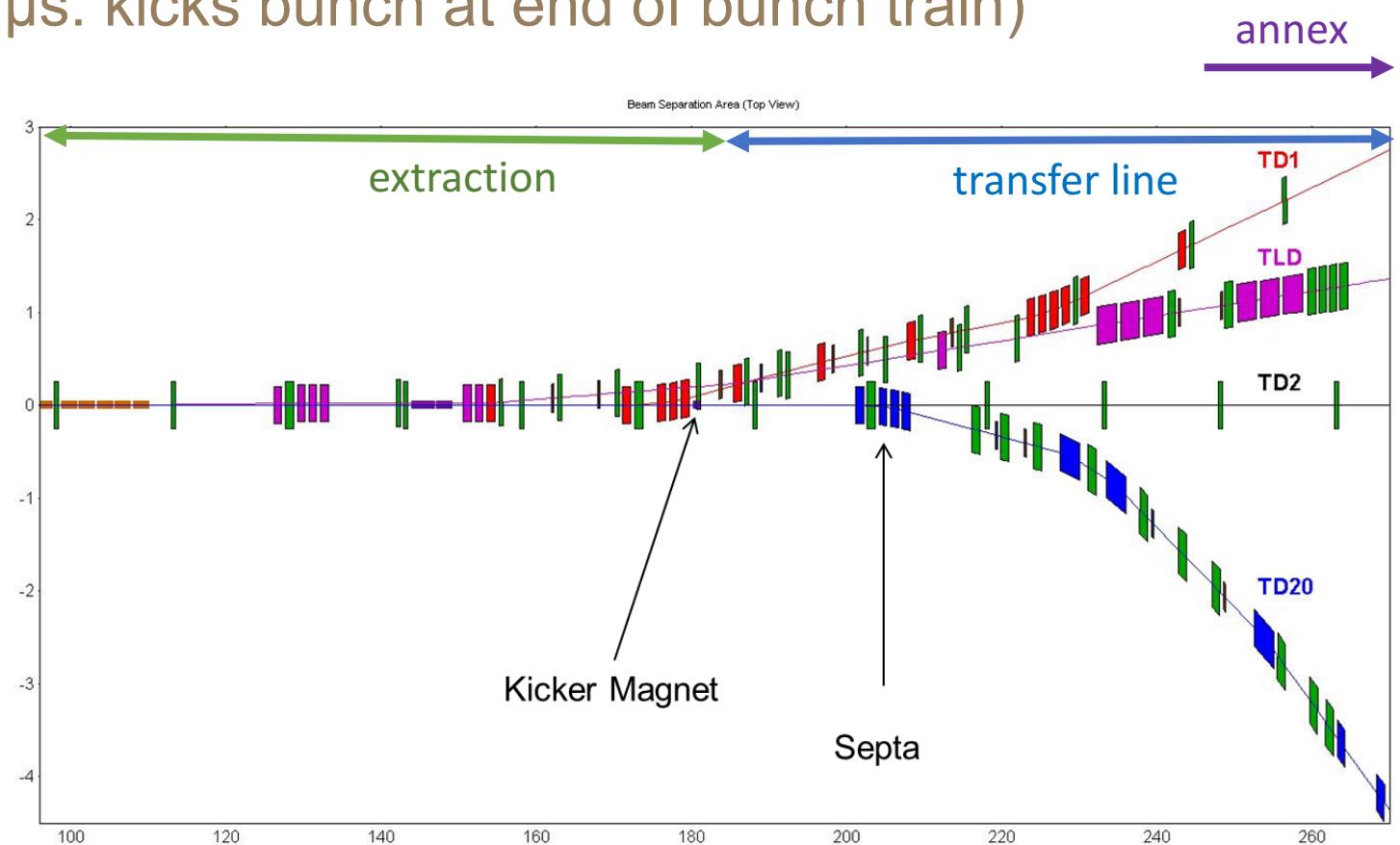
BEAMLINE LAYOUT

Design of magnets for beam extraction and then beam transfer to LUXE

- Most magnets use design already operating today in XFEL.EU
- New fast kicker magnets (2 μ s: kicks bunch at end of bunch train)

Installation requires

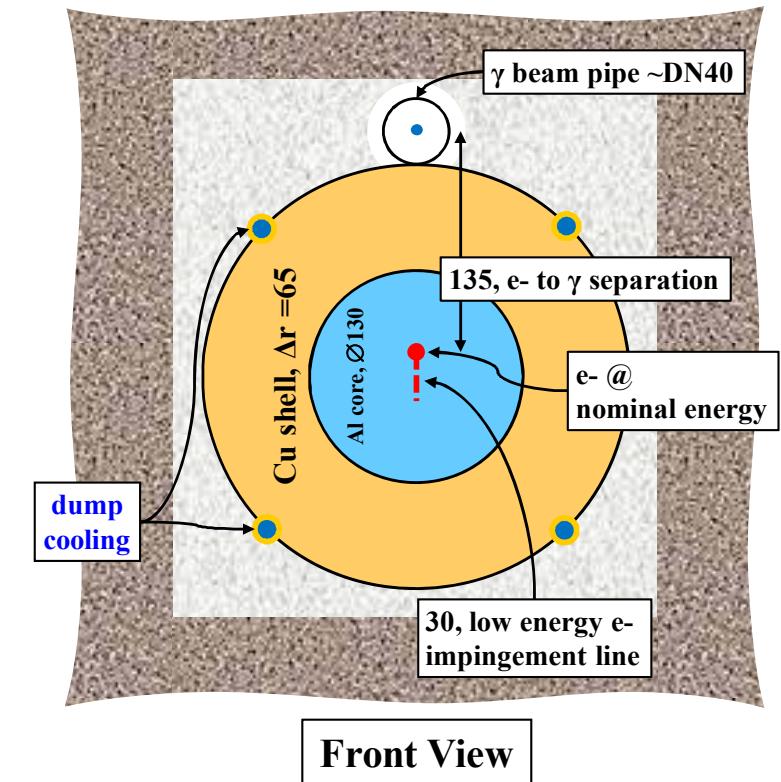
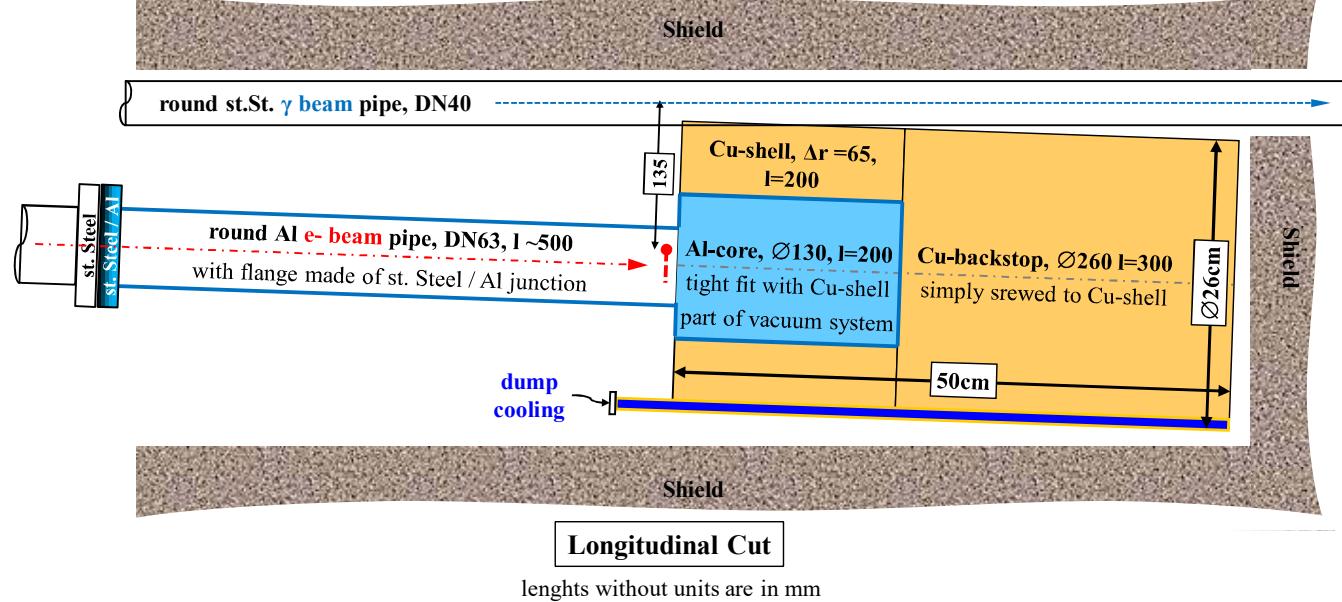
- 5 weeks for extraction
- 7 weeks for transfer line



F. Burkart, W. Decking

BEAM DUMP

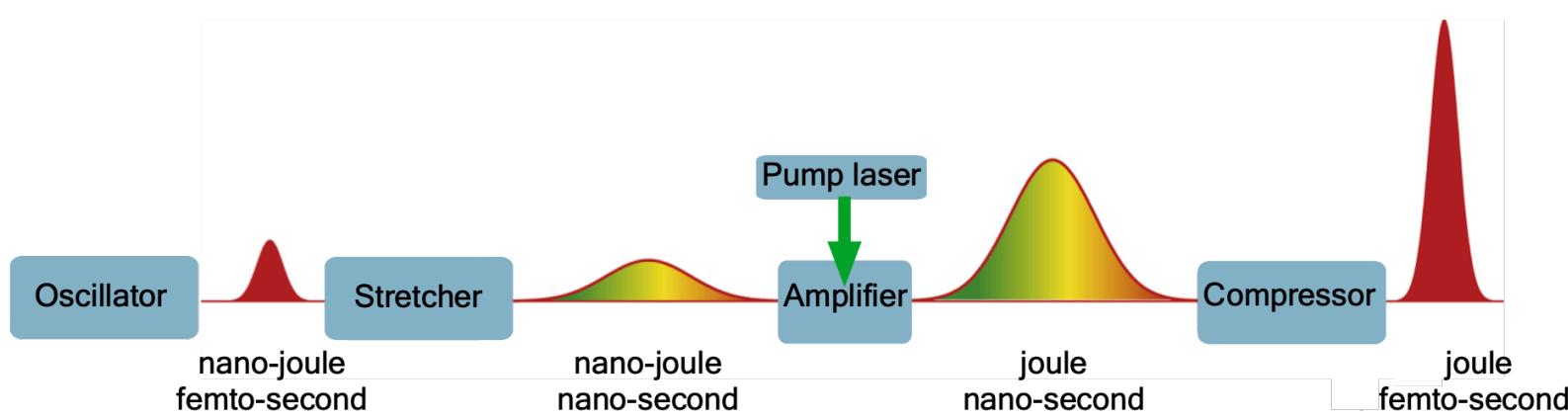
Beam needs to be safely dumped, design (with radioprotection group) well advanced



F. Burkart, M. Schmitz (DESY)



LASER TECHNOLOGY



© Nobel Media AB. Photo: A.
Mahmoud
Gérard Mourou
Prize share: 1/4



© Nobel Media AB. Photo: A.
Mahmoud
Donna Strickland
Prize share: 1/4

- Use Chirped Pulse Amplification (CPA) technique
 - Half of the NP 2018 shared by Gerard Mourou and Donna Strickland "for their method of generating high-intensity, ultra-short optical pulses."
 - Ti:Sa laser with 800 nm wavelength
 - Energy focussed strongly in both time and space => high intensity

LASER PARAMETERS

Parameter	Initial stage	Stage 1	Stage 2
Laser energy after compression [J]	0.9	9	
Percentage of laser in focus [%]	40	40	
Laser energy on focus [J]	0.36	3.6	
Laser pulse duration [fs]	30	30	
Laser repetition rate [Hz]	1	1	
Laser-beam crossing angle [degrees]	17	17	
Laser focal spot FWHM [μm]	8	8	3
Peak intensity [10¹⁹ W/cm²]	1.6	16	110
Peak intensity parameter ξ	2	6.2	16
Peak quantum parameter χ: Ebeam=17.5 GeV	0.41	1.3	3.3
Ebeam=14.0 GeV	0.32	1.0	2.6

Laser intensity:

$$I = \frac{E_L}{\Delta t \pi d^2}$$

with

E_L : energy (J)

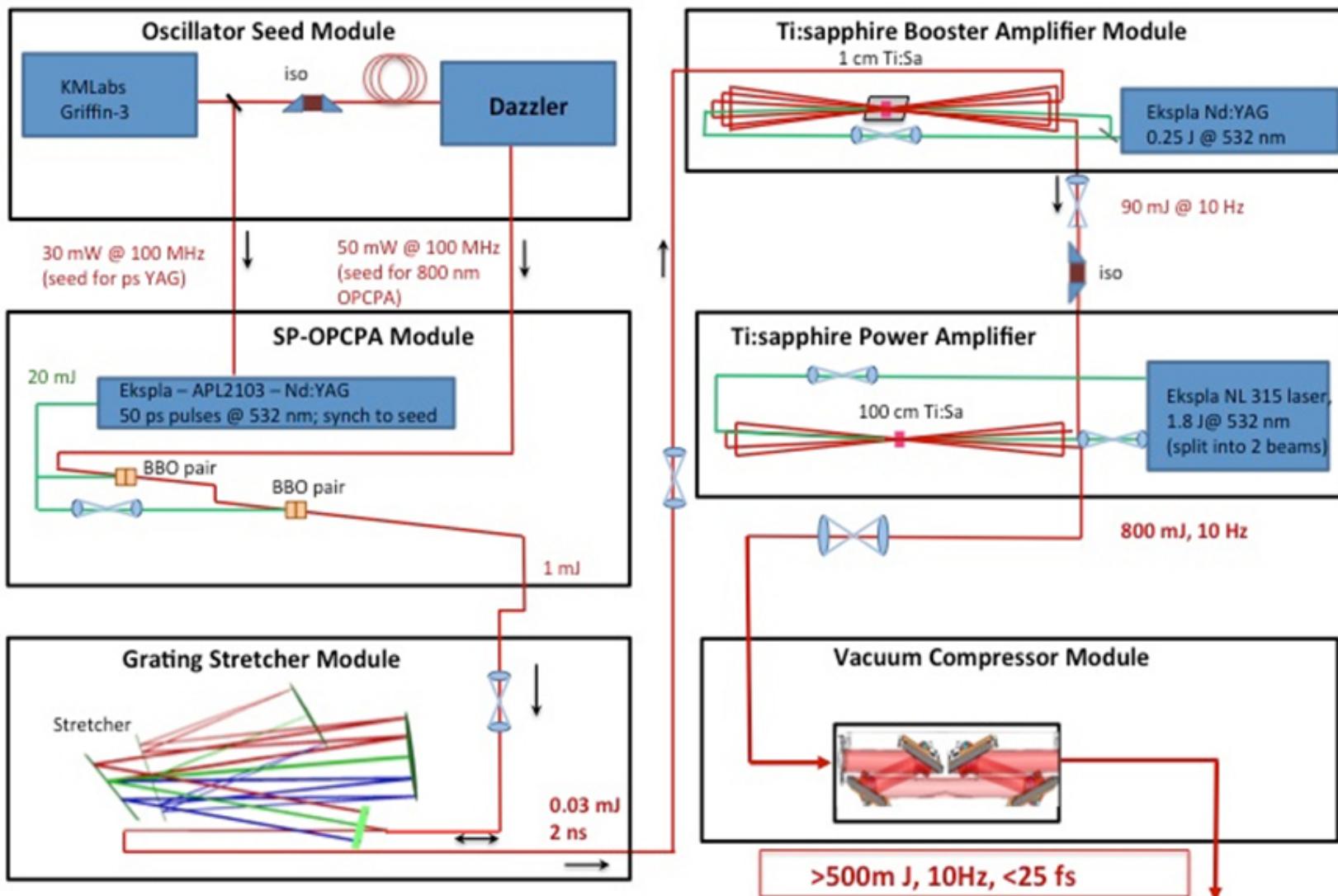
Δt : pulse length (s)

πd^2 : focus area (m^2)

Lower intensities achieved by de-focussing laser or stretching pulse



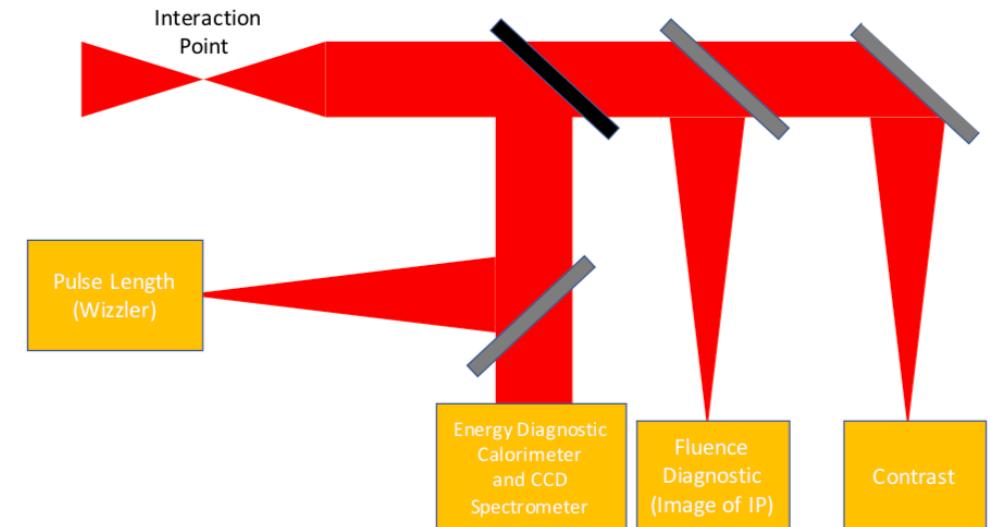
LASER DESIGN



I. Pomerantz (Tel Aviv), G. Sarre
(Belfast), M. Zepf (HZI Jena, Jena,
Belfast) and others

LASER DIAGNOSTICS

- Aim to control intensity at level of 5-10%
 - Cannot measure it directly
- Several diagnostics measurements planned to measure parameters
 - Energy
 - Fluence (Energy/area)
 - Pulse length
- Laser shots can vary by ~15% for stable laser at this power
 - System can be used to tag intensity of individual shots





PARTICLE DETECTION AND SIMULATION RESULTS

RATES OF PARTICLES

e+laser

Location	particle type	rate for $\xi = 2.6$	rate for $\xi = 0.26$
e^- detector	$e^-, E_e < 16 \text{ GeV}$	1.5×10^9	6×10^6
e^+ detector	e^+	15.3	< 0.01
Photon detector	γ	6×10^{10}	1×10^7
Photon detector (W foil)	e^+ and e^-	6×10^6	1×10^4
Photon detector (W wire)	e^+ and e^-	1.5×10^5	1×10^2

γ +laser

Location	particle type	rate for $\xi = 2.6$	rate for $\xi = 1.2$
e^- detector behind converter	$e^-, E_e < 13 \text{ GeV}$	2×10^7	
e^+ detector behind converter	e^+	9×10^4	
photons after converter	γ	1.3×10^8	
e^\pm detector behind IP	e^-/e^+	5	1×10^{-2}
Photon detector	γ	1.3×10^8	
Photon detector	e^+ and e^-	160	

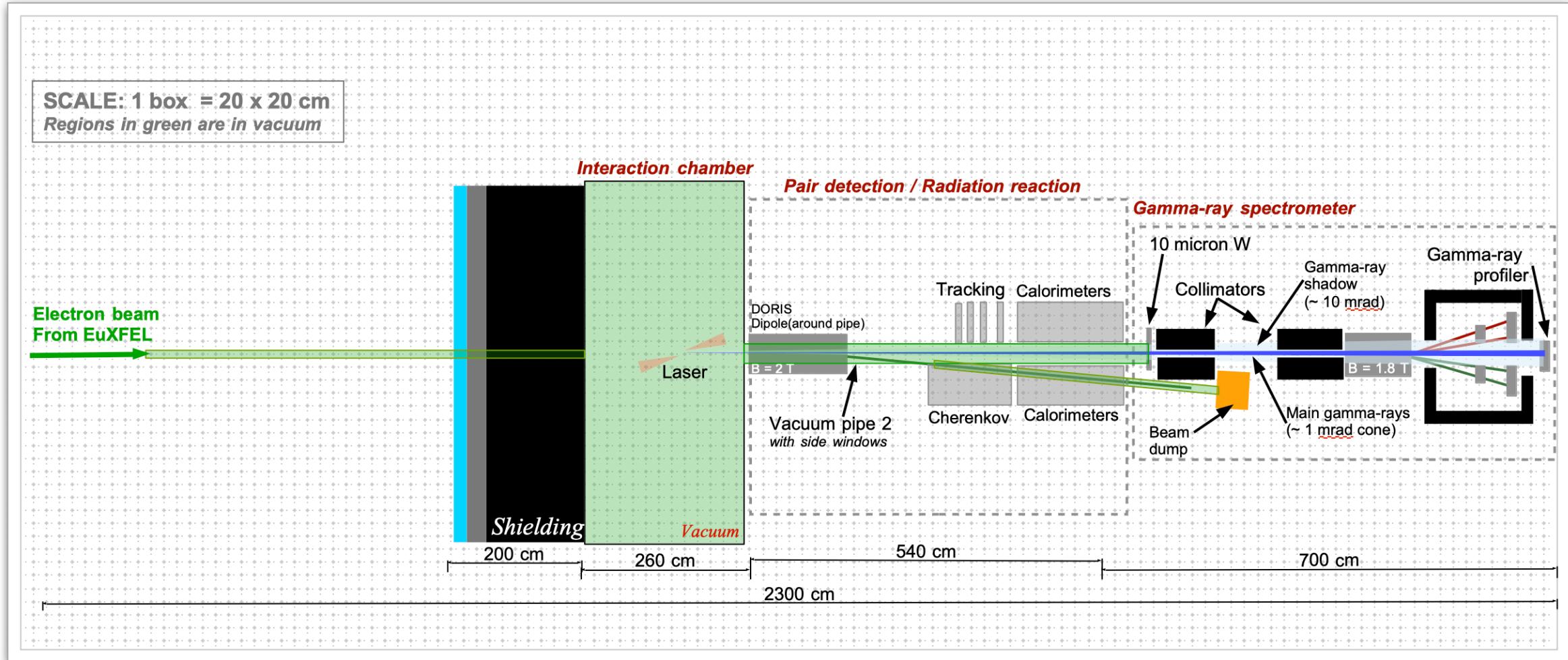
~1000/ μm

~0.01/m

=> Very different rates of particles => need different technologies

ELECTRON LASER COLLISIONS

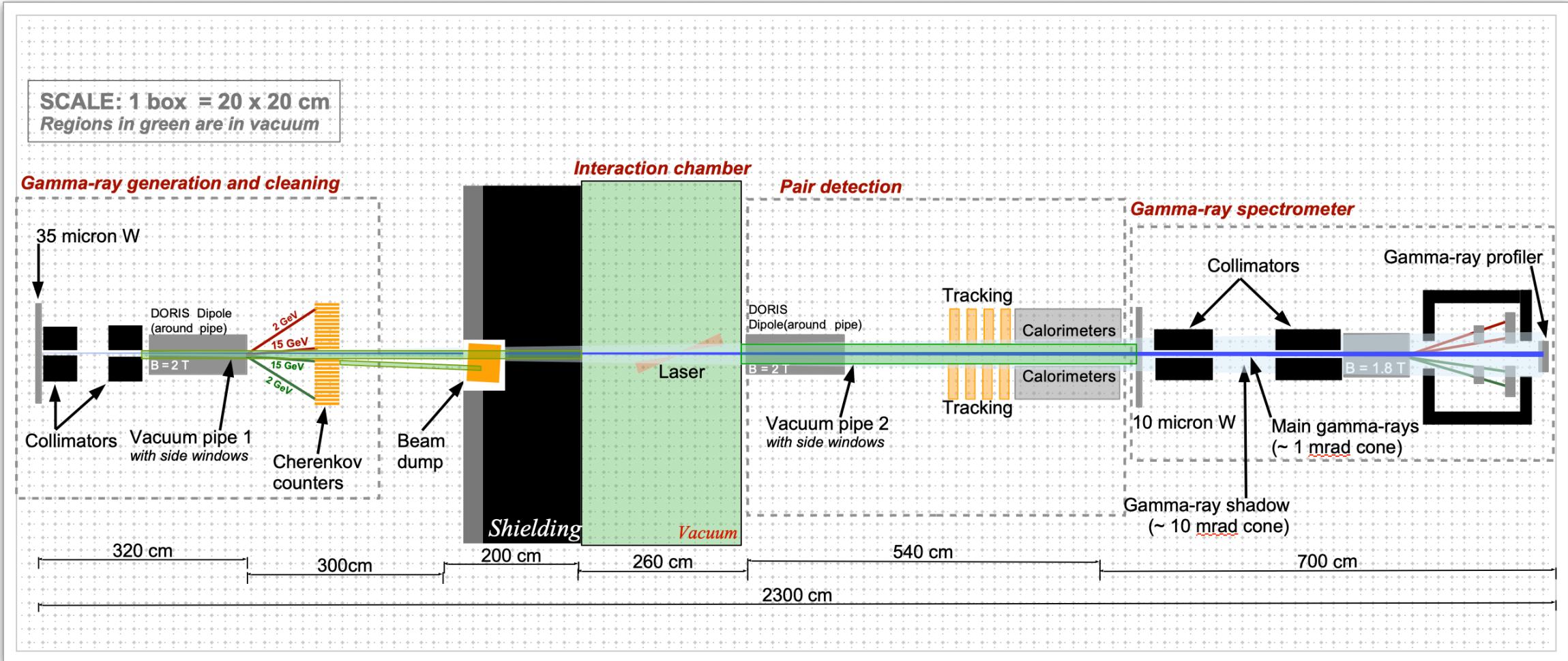
Compton and trident processes: $e^- + n\omega \rightarrow e^- + \gamma$ and $e^- + n\omega \rightarrow e^-e^+e^-$



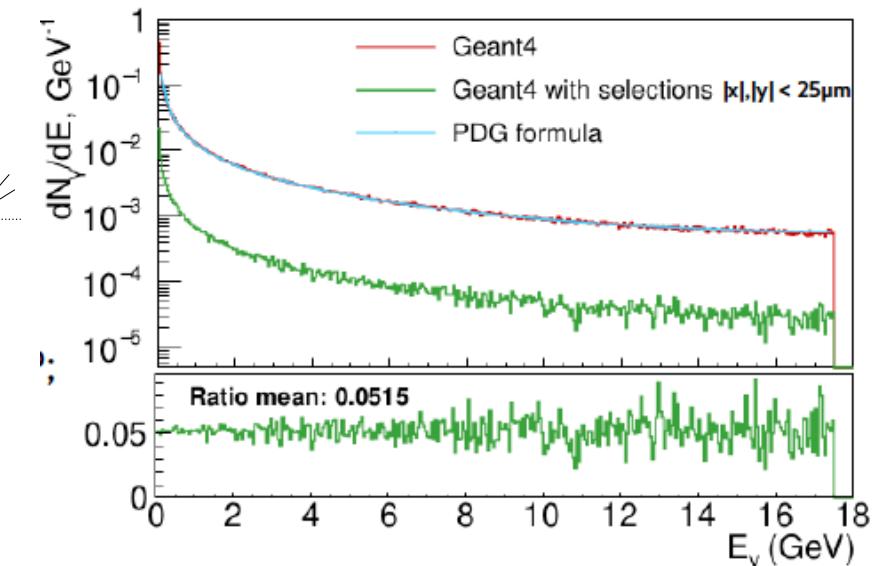
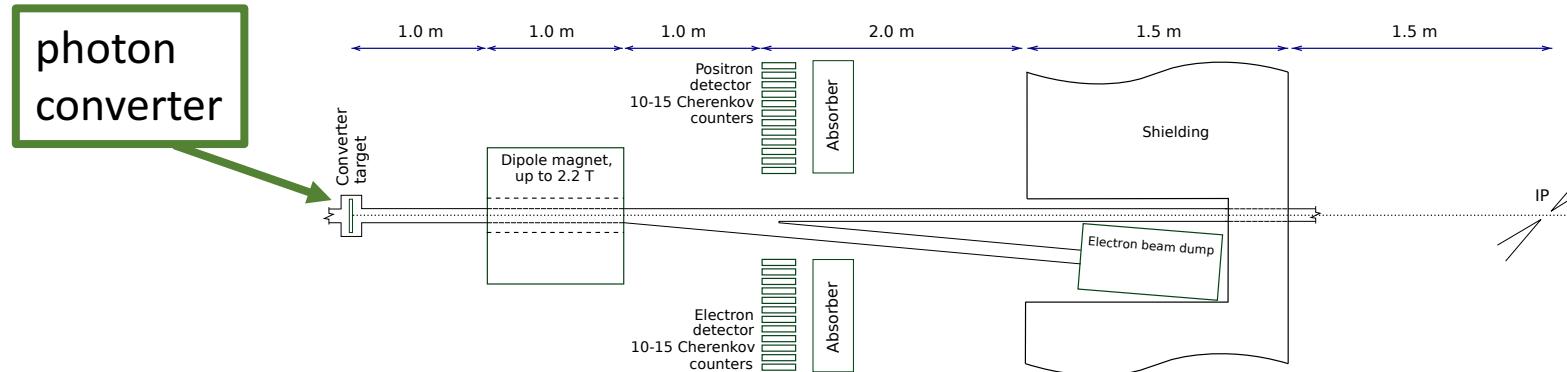


PHOTON LASER COLLISIONS

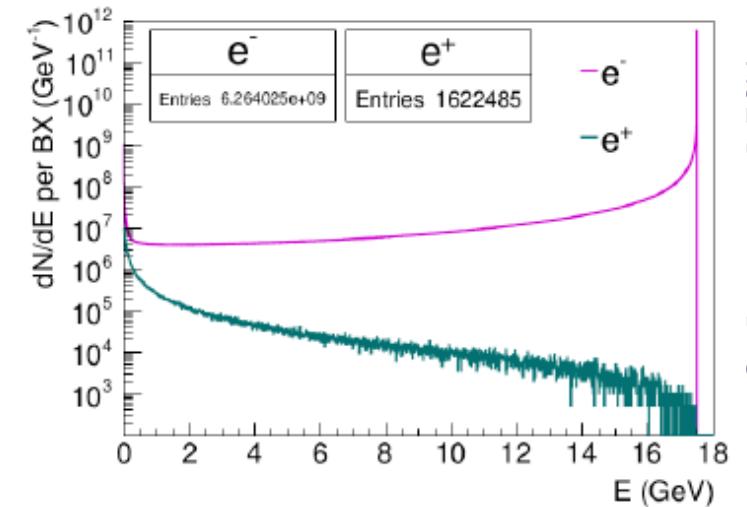
Pair production (Breit-Wheeler) process: $\gamma + n\omega \rightarrow e^- + e^+$



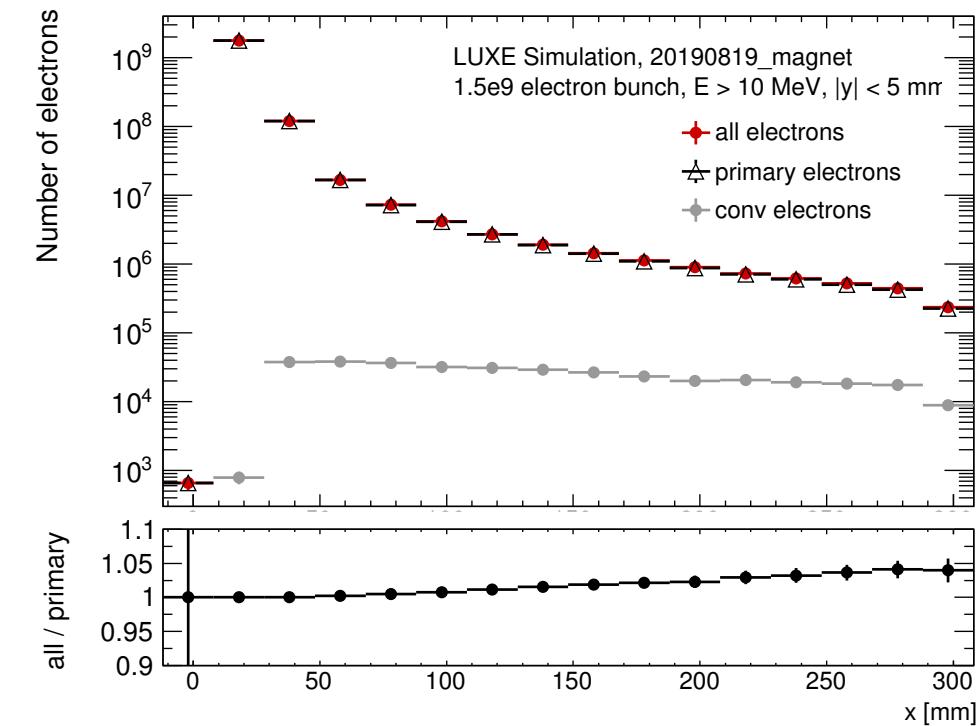
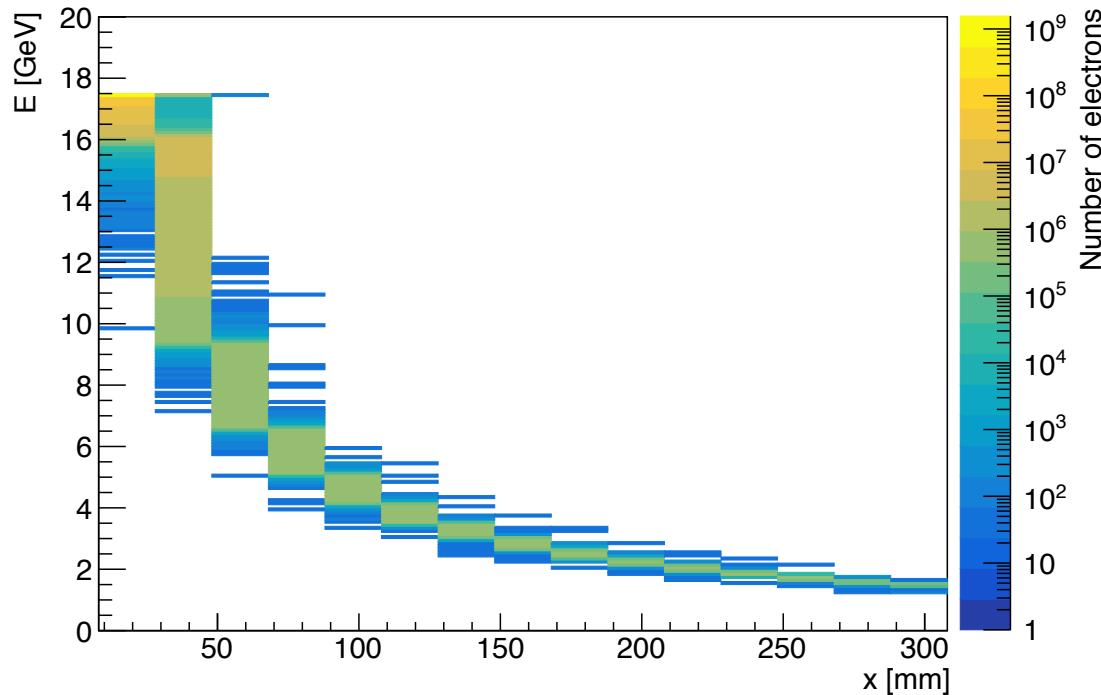
HIGH-ENERGY PHOTON FLUX



- **Simulation of converter using Geant4**
 - Tungsten Target with $0.01 X_0$ ($35 \mu\text{m}$) => 1% at IP
- **Spectrum of photon energies important to know**
 - Measure by observing electrons and positrons right after dipole magnet
- **Particle detection**
 - 2T magnet followed by array of Cherenkov detectors measures flux vs impact position => energy spectrum



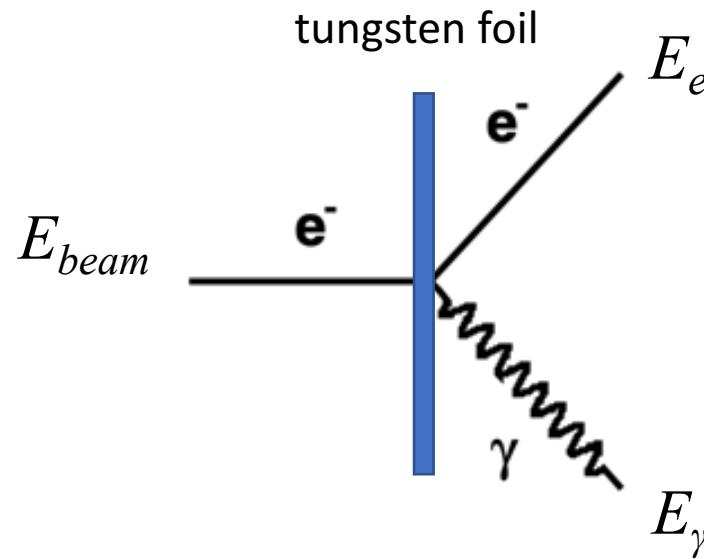
PHOTON FLUX MEASUREMENT: ELECTRONS



Electron energy measured based on position behind dipole magnet

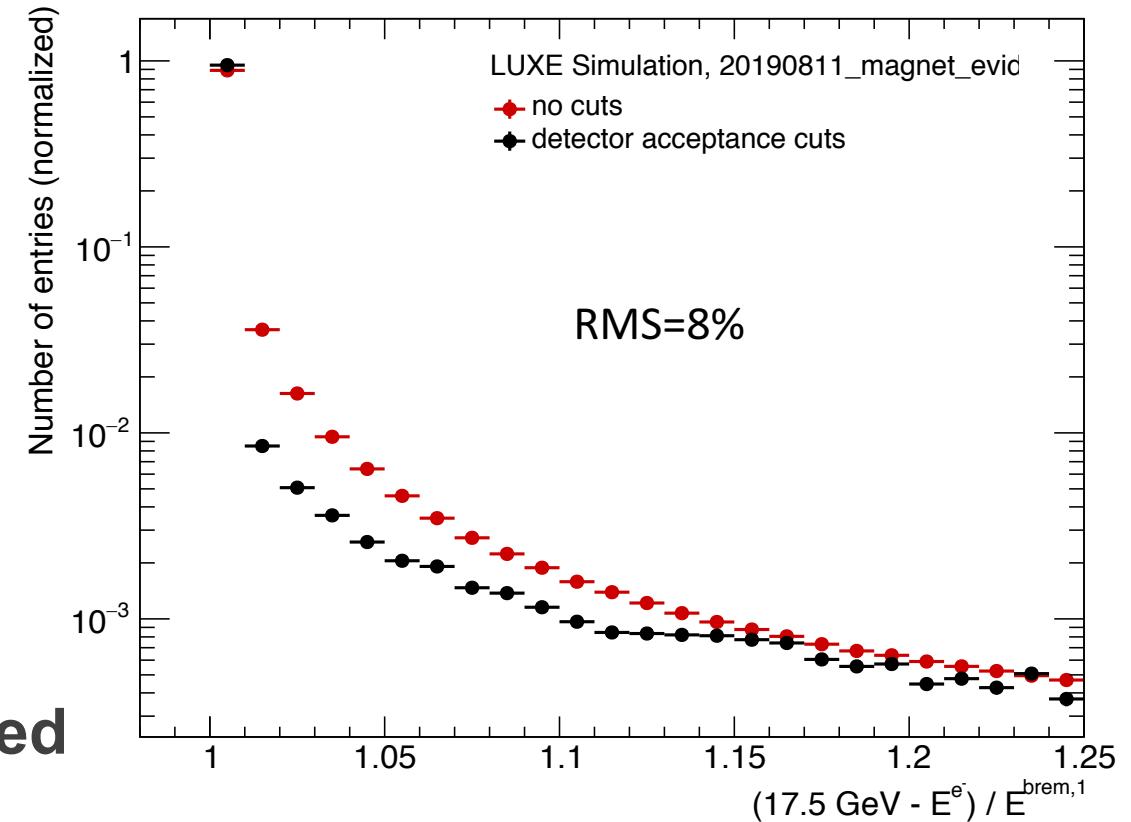
- Dominated by primary electrons
- Contamination of converted electrons small (estimated from positron flux)
- Electron rates high: $\sim 10^5\text{-}10^7/\text{mm/event}$

PHOTON ENERGY MEASUREMENT

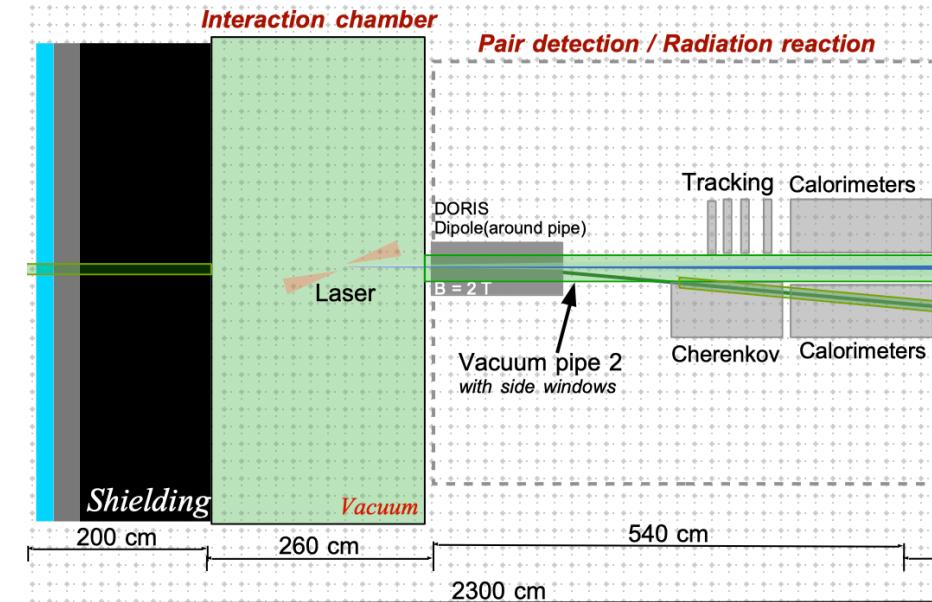
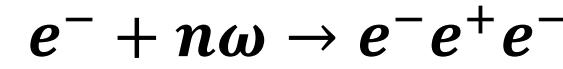
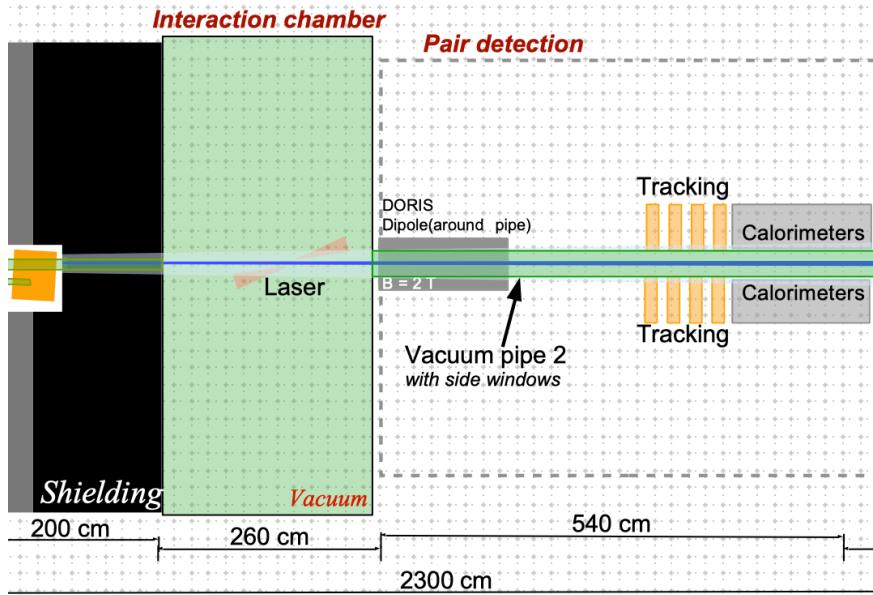
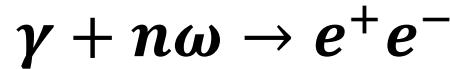


Photon energy determined from measured electron energy to within $\sim 10\%$:

$$E_\gamma = E_{beam} - E_e$$



ELECTRON AND POSITRON DETECTORS

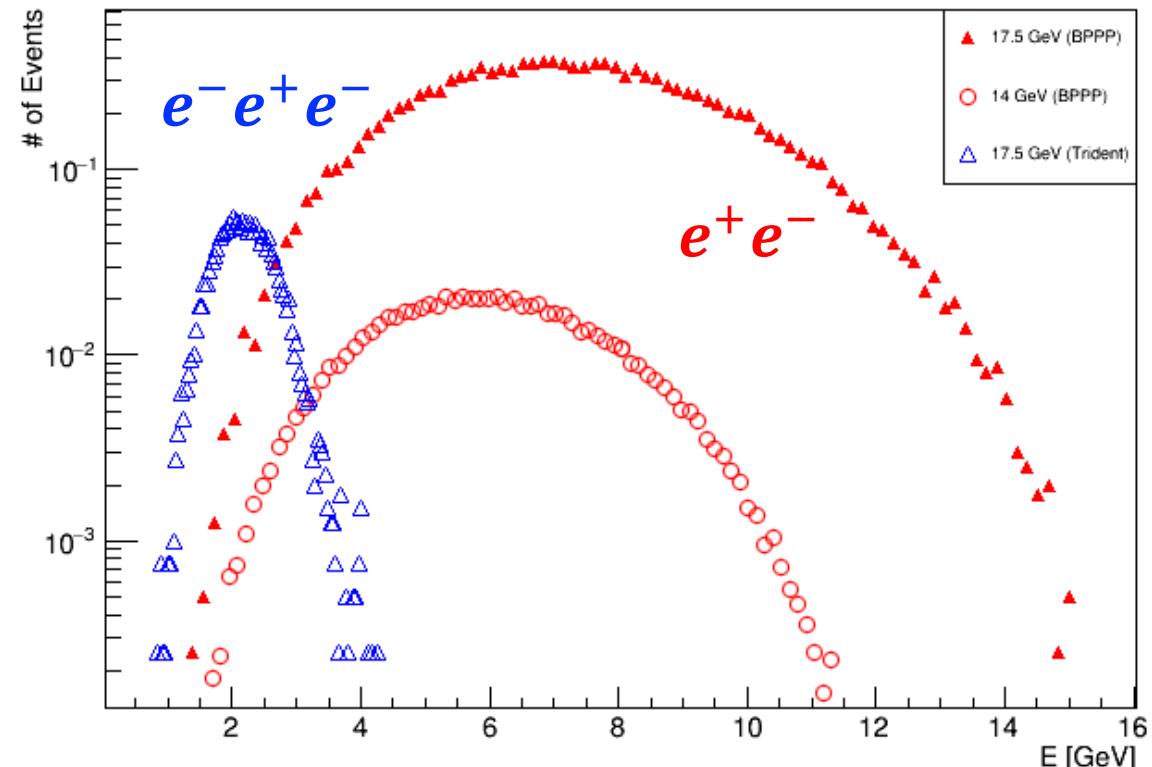


- **Pair production:**
 - e^+ and e^- rate $\sim 0.01\text{-}100$ \Rightarrow silicon pixel detectors and calorimeters
- **Trident:**
 - e^+ rate $\sim 0.01\text{-}100$ \Rightarrow silicon pixel detectors and calorimeters
 - e^- rate $\sim 10^6\text{-}10^9$ \Rightarrow Cerenkov counters and calorimeter/absorber

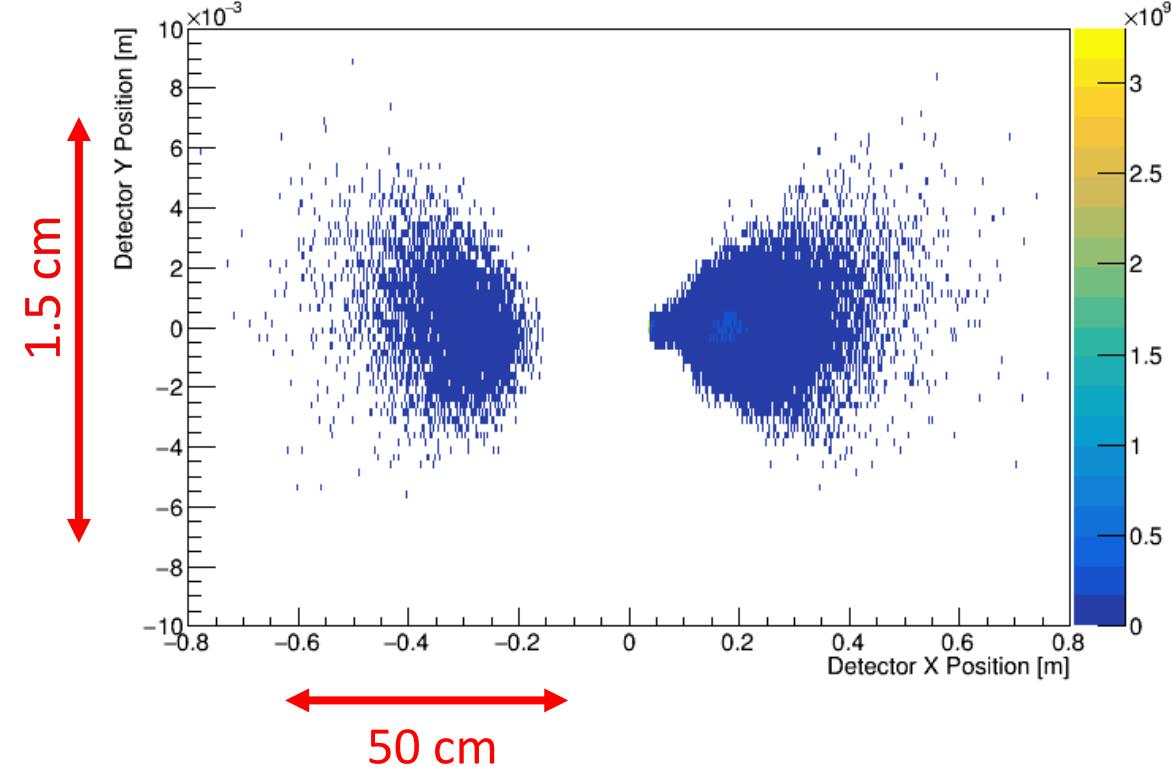
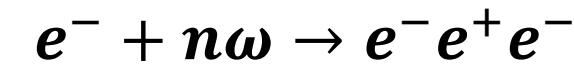
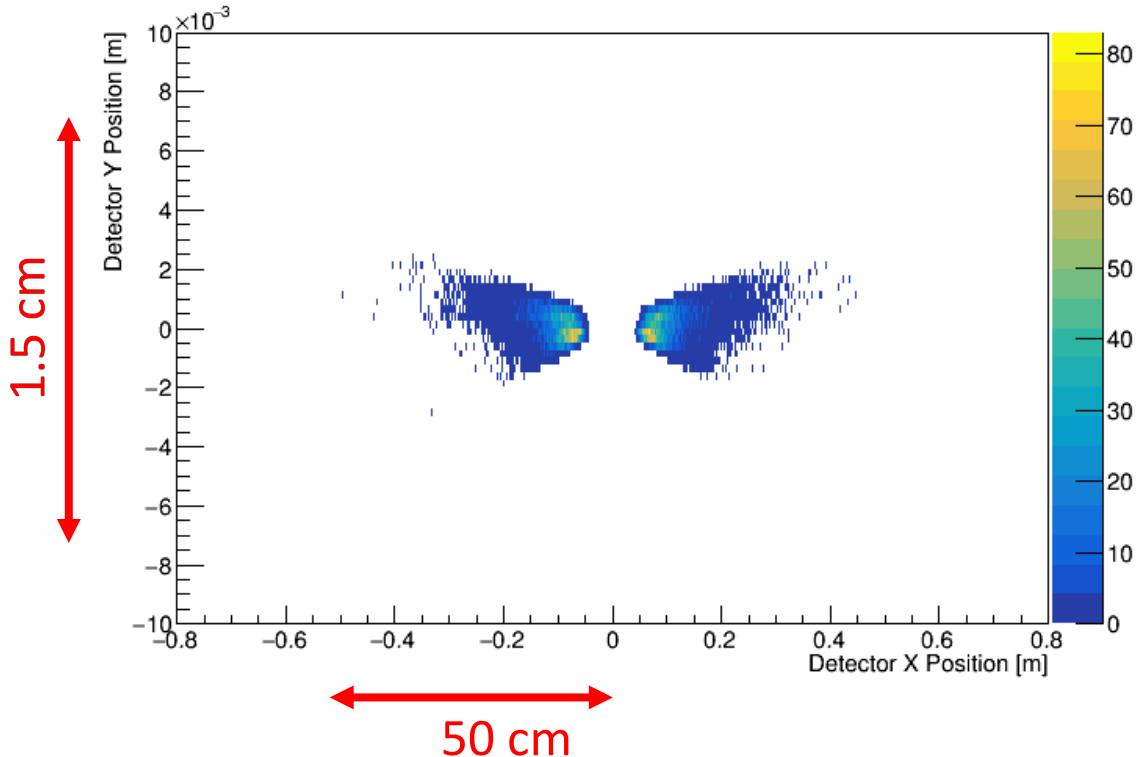
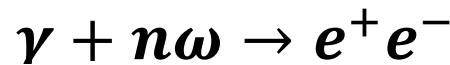
SIMULATION RESULTS

- Monte Carlo simulation of expected signatures used
 - By A. Hartin, UCL
- Energy spectrum spans 1-15 GeV
 - Energies significantly lower for trident process
- For trident process uses “two-step” process only
 - Calculation of one-step trident ongoing

Positron Energy Spectrum

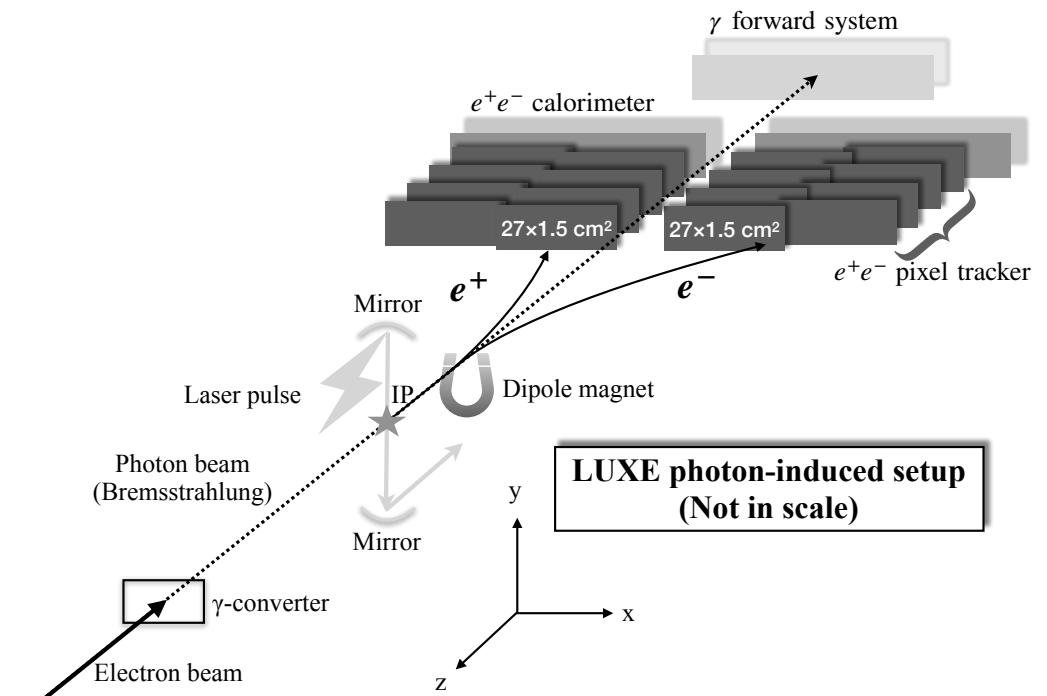
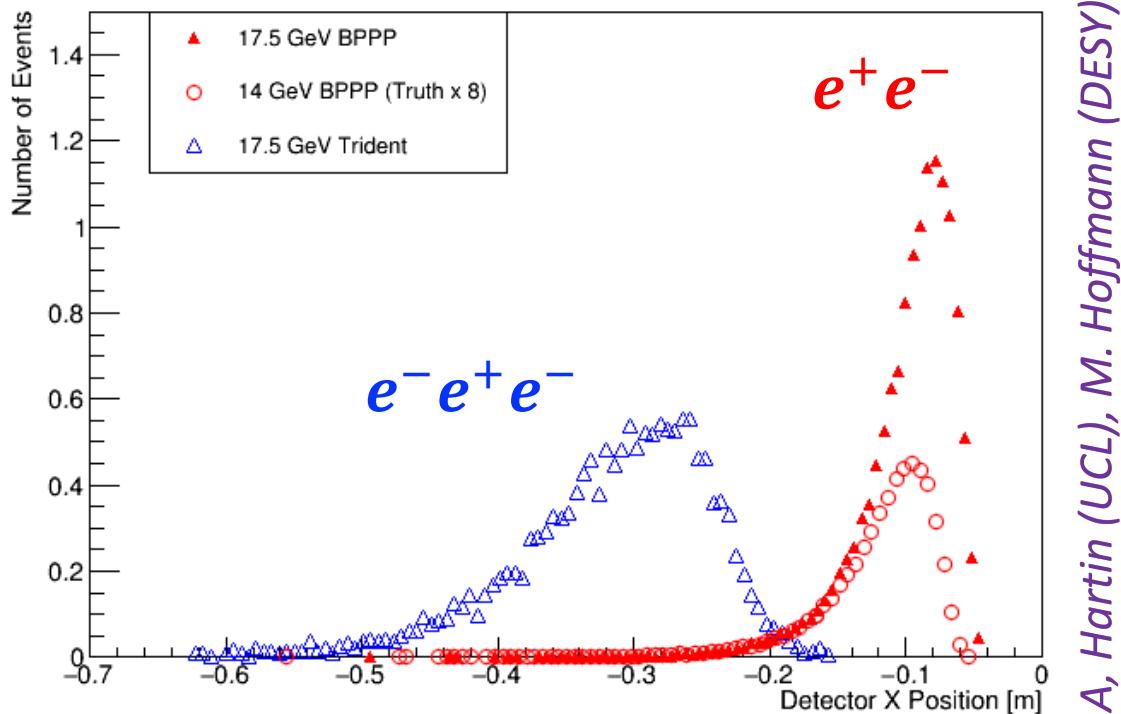


DETECTOR OCCUPANCIES AFTER INTERACTION POINT



- Vertical direction: very small spread for both processes
- Horizontal direction: particles contained within ~ 50 cm

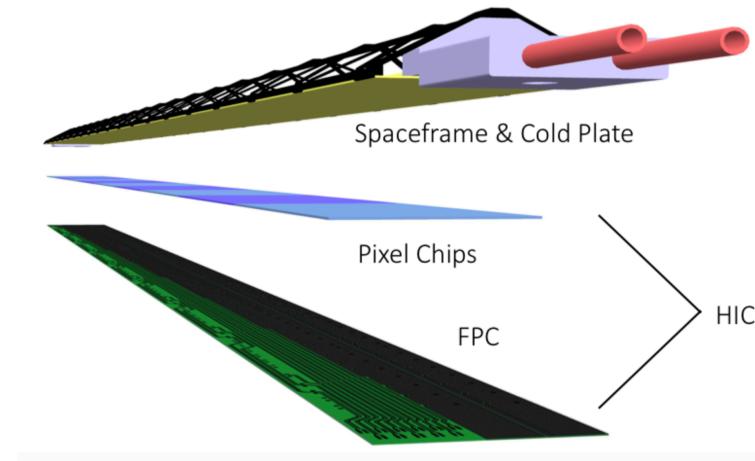
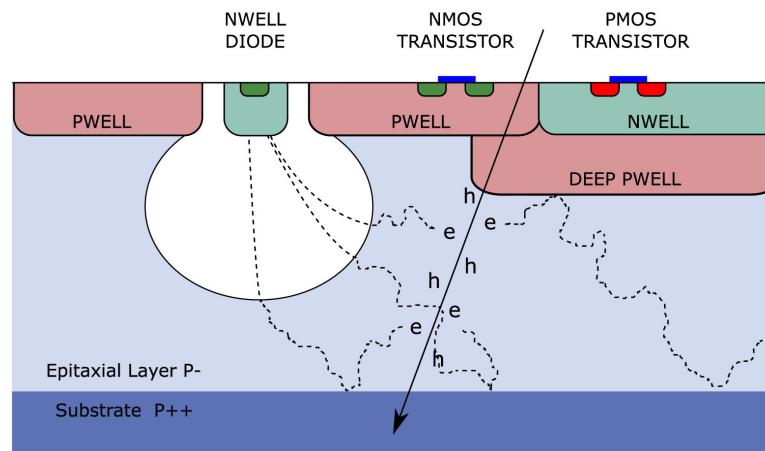
HIT POSITION AT FIRST DETECTOR PLANE



Detectors need to span about ~50 cm to have acceptance >95%:

- $e^- + n\omega \rightarrow e^-e^+e^-$ process: acceptance ~95%
- $\gamma + n\omega \rightarrow e^+e^-$ process: acceptance >99%

SILICON DETECTORS



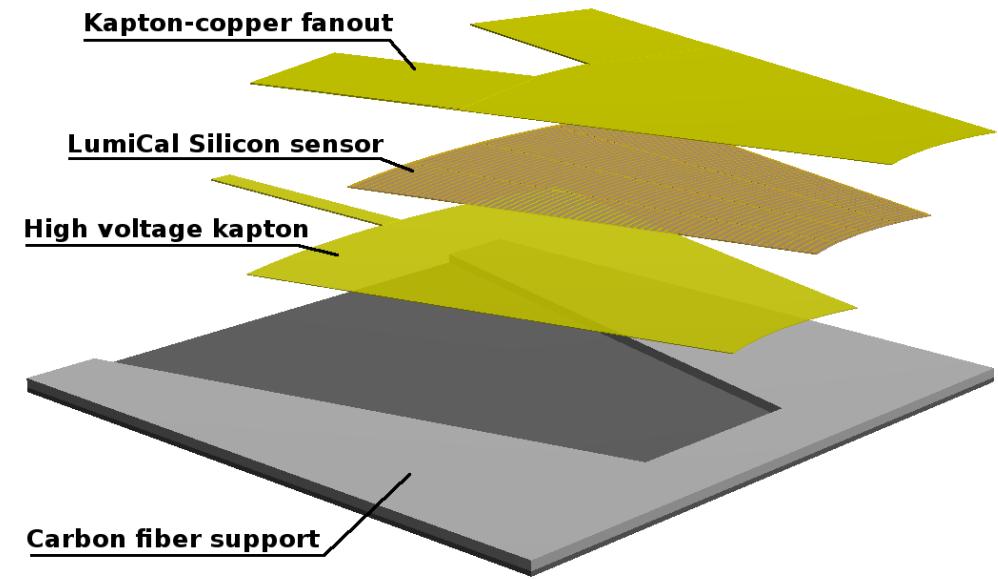
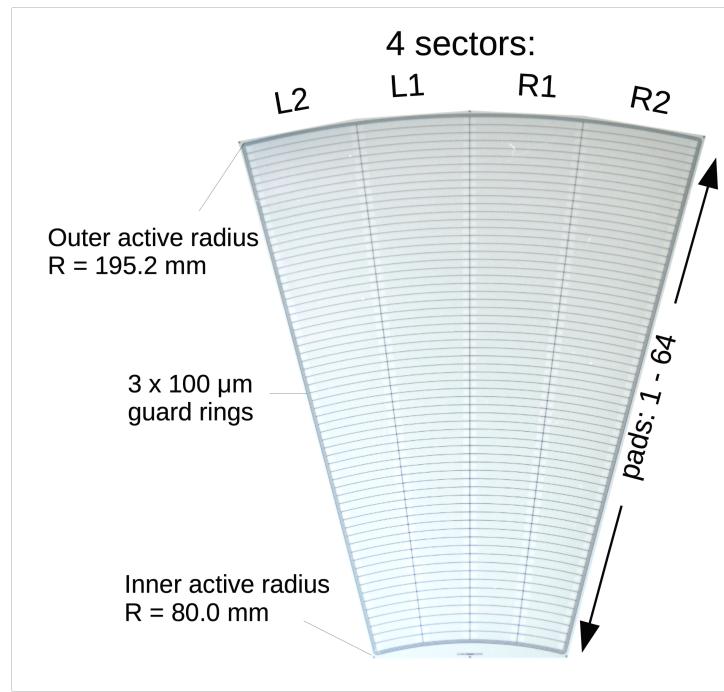
ALPIDE pixel detectors

- Developed by ALICE collaboration
- Staves of 27 cm length; sensor size $1.5 \times 1.5 \text{ cm}^2$
 - Achieve full coverage with two staves placed next to each other
- Pixel size: $27 \times 29 \mu\text{m}^2 \Rightarrow$ Spatial resolution $\sim 5 \mu\text{m}$
- Plan to use four layers staggered behind each other

N. Hod (Weizmann Inst.)

Redundant tracking possible, important for beam background rejection

CALORIMETERS

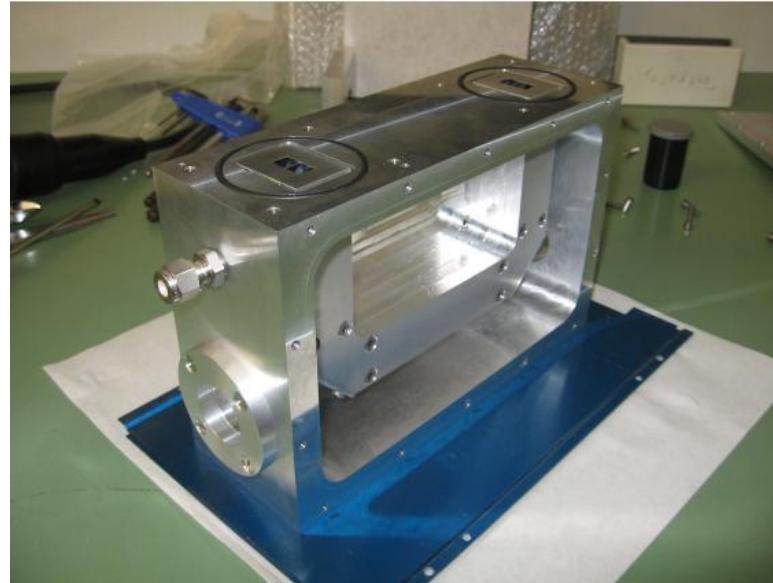
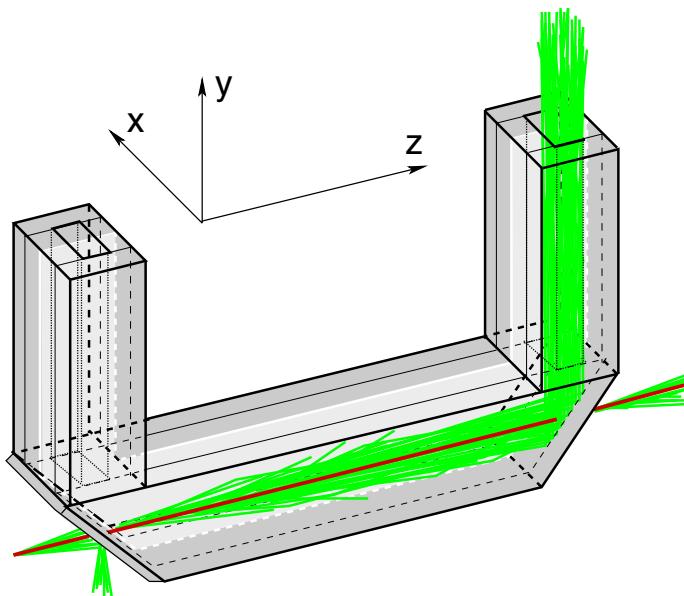


Y. Benhammou, H. Abramowicz, A. Levy (Tel Aviv U)

High granularity silicon Tungsten calorimeter

- Developed for luminosity measurement at linear colliders (LUMICAL)
- 20 tungsten absorber plates (3.5mm), Si layers in gaps ($320\text{ }\mu\text{m}$)
- Geometry adapted to fit needs of LUXE ($\sim 50\text{cm}$ long, vertical spread $<1\text{mm}$)
- Moliere radius 8 mm, Prototyped and test beam measurements available

CHERENKOV COUNTERS



Use Cherenkov detectors in high-flux regions

- Use design developed for ILC polarimeters
- Linearity better than 0.1% over dynamic range spanning 10^3
- Threshold of ~ 10 MeV => robust against background from low energy radiation
- Plan to use array of 15 detectors with cross section of $2 \times 2 \text{ cm}^2$

POSITRON RATE VS LASER INTENSITY

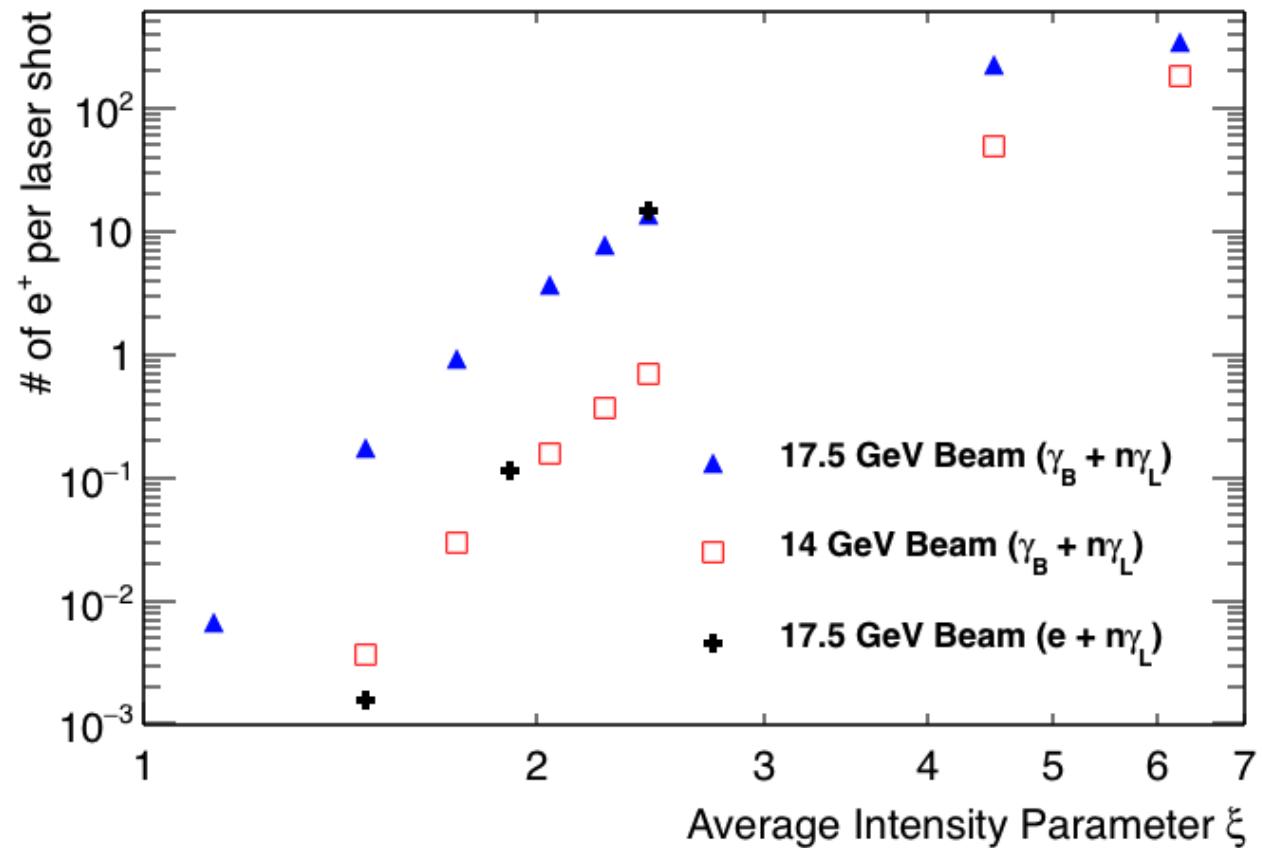
Main expected result of experiment

Low laser intensity

- Encounter power-law behaviour

High intensity

- Should observe deviation from power-law behaviour
- Aim to quantify by extracting coefficient



POSITRON RATE VS LASER INTENSITY

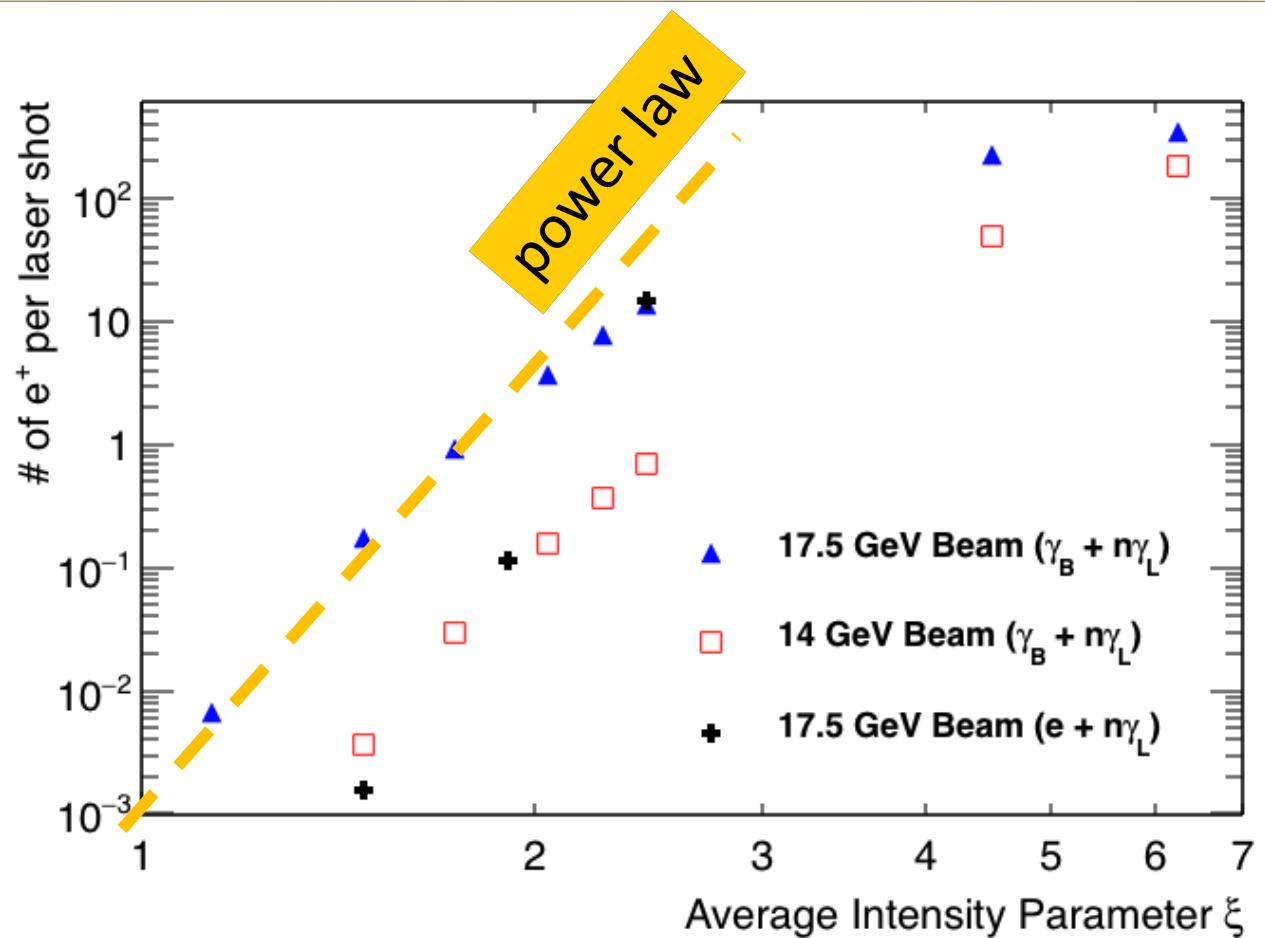
Main expected result of experiment

Low laser intensity

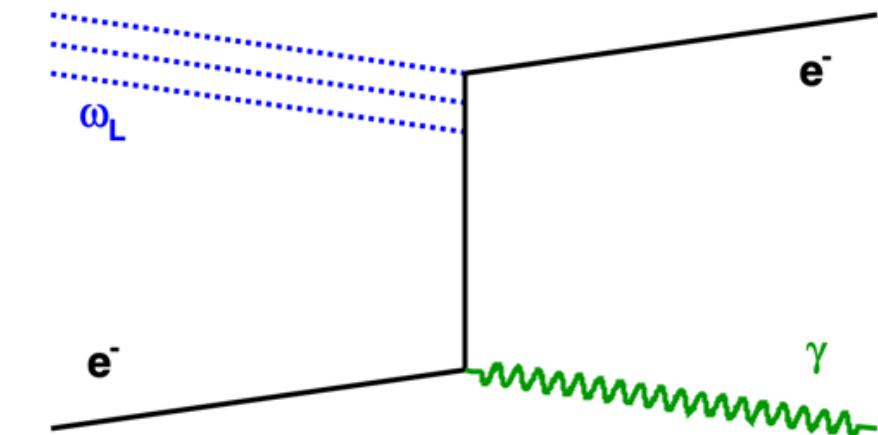
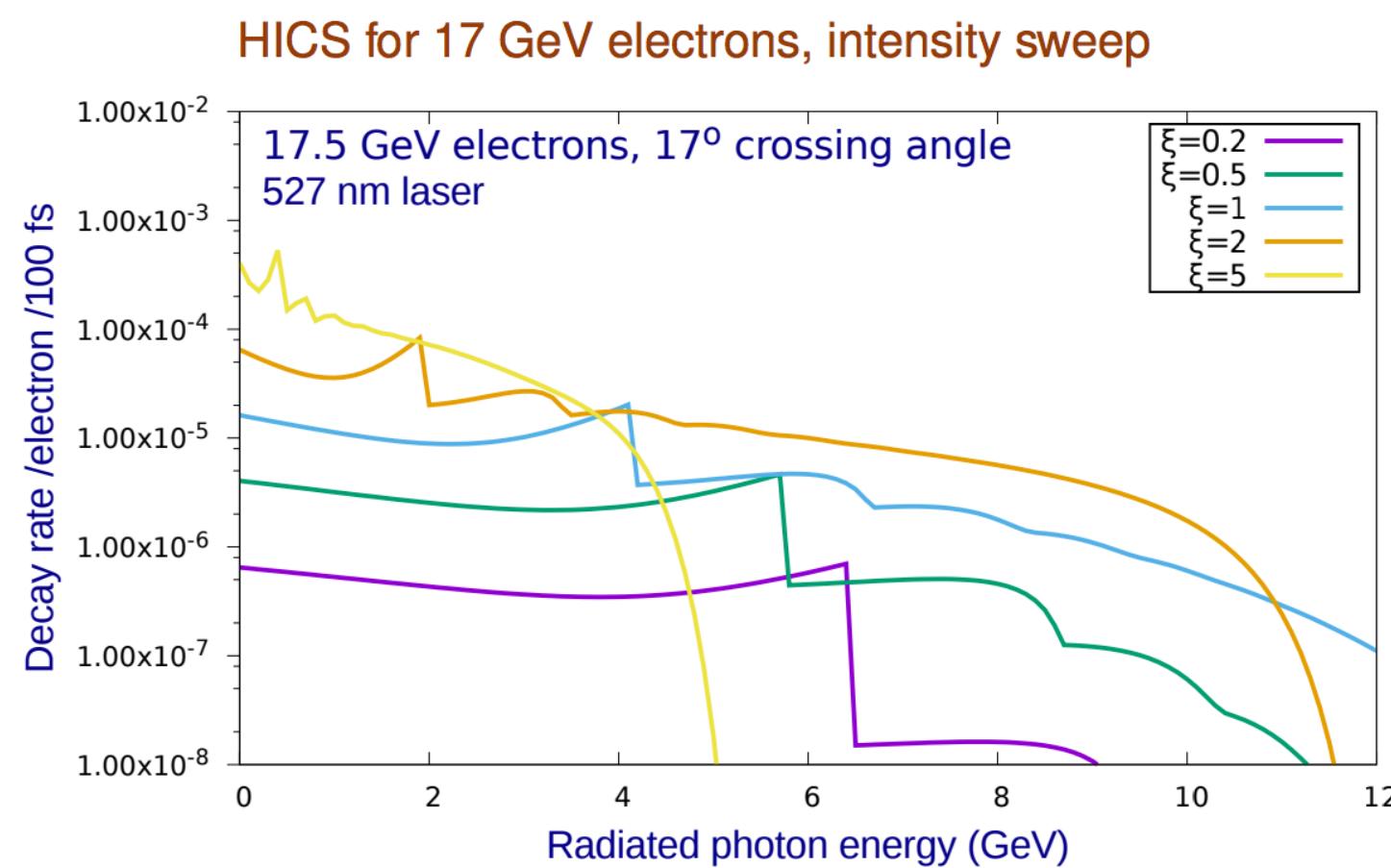
- Encounter power-law behaviour

High intensity

- Should observe deviation from power-law behaviour
- Aim to quantify by extracting coefficient



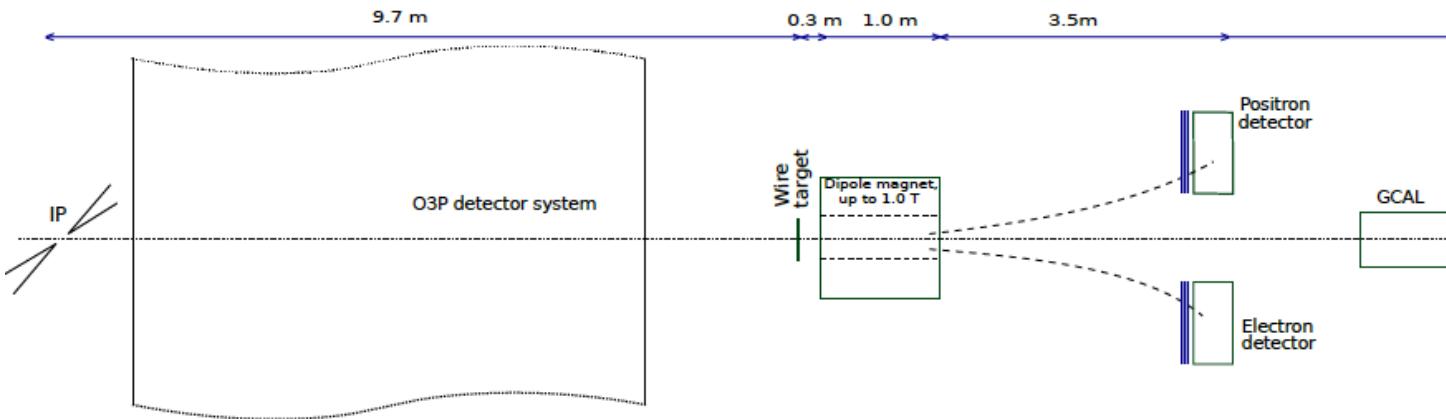
NON-LINEAR COMPTON SCATTERING: $e^- + n\omega_L \rightarrow e^- + \gamma$



$$p_i + k \frac{\xi^3}{2\chi_i} \rightarrow p_i^2 = m^2(1 + \xi^2)$$

A, Hartin (UCL)

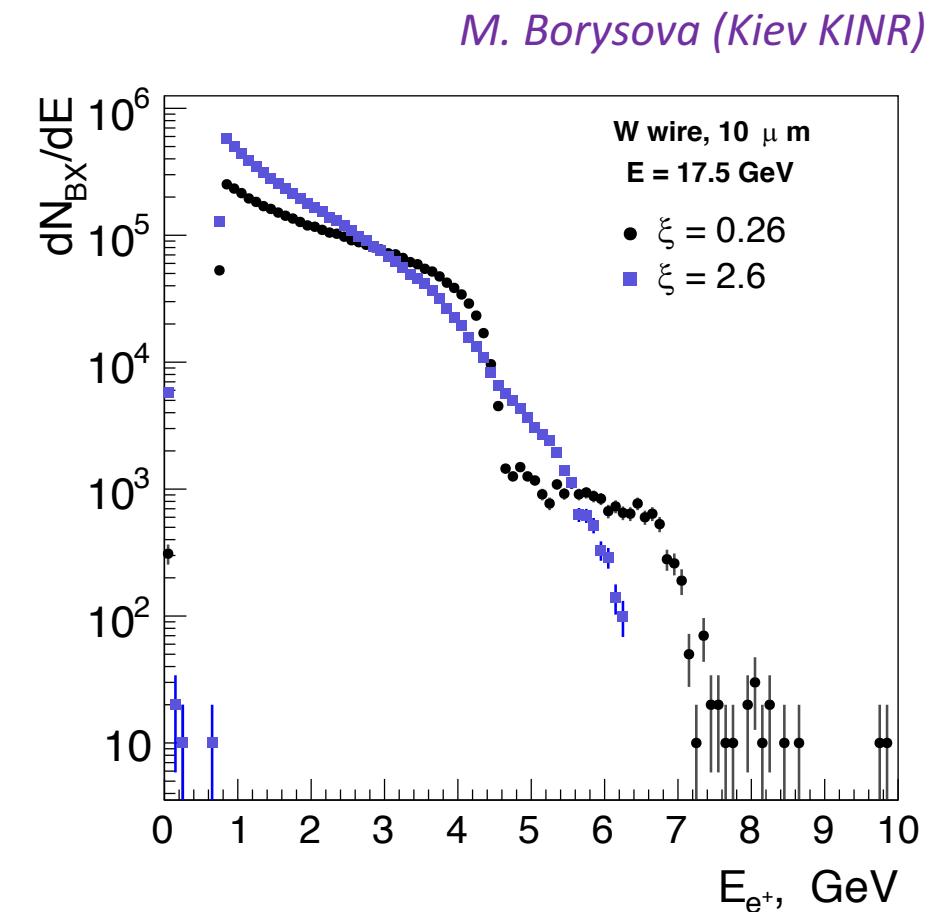
NON-LINEAR COMPTON PROCESS



Measure photon flux and energy in “photon detection system”

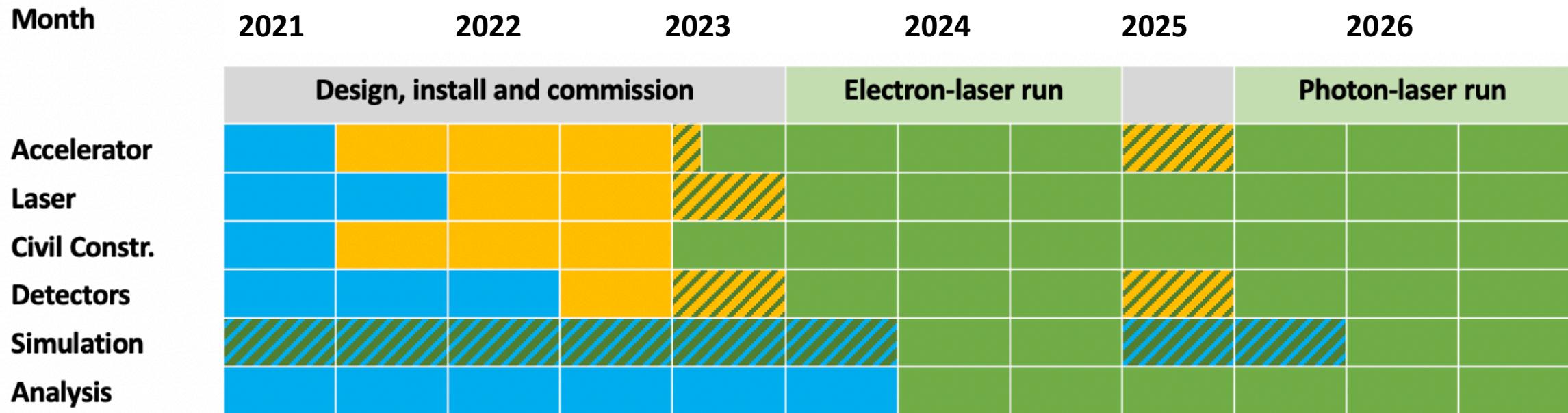
- Photon flux very high ($>10^7$ per laser shot)
- Thin wire to convert photons to e^+e^- pairs

Compton edges observable in e^\pm energy spectra at low ξ





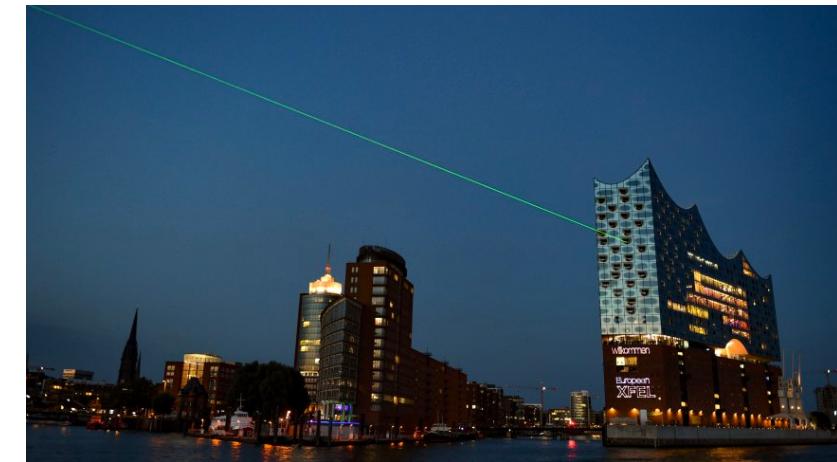
TENTATIVE TIME SCALES: 2020 AND BEYOND



- Fall 2020
 - CDR for LUXE experiment
- Nov/Dec 2020
 - Start preparatory work for installation; main installation following year

CONCLUSIONS

- **LUXE will boil the vacuum using a minute fraction of European XFEL electron beam**
 - Measure several phenomena predicted more than 60 years ago
 - Test quantum field theory in a new regime
- **International collaboration of performed feasibility study**
 - “Letter of Intent” released in September
 - **Only possible in synergy between accelerator, laser and particle physicists**



S. Weinberg: “*My advice is to try crazy ideas and innovative experiments. Something will come up.*”



STEVEN WEINBERG

Steven Weinberg (03/2019, interview at APS):

Do you think the problems faced by particle physicists today are different from those that you faced as a young scientist?

I do. It was a different situation 50 years ago. Back then, we had experimental data coming out of our ears, and a lot of it didn't seem to fit any pattern. The problems seemed formidable, but there were so many ways to go with new theories. It really was a thrilling time to be a physicist.

Nowadays, it's very hard to think of a challenge that we can get our teeth into. The current puzzles don't offer theorists many opportunities to propose solutions that can be tested experimentally.

Do you have any advice to offer the next generation?

*Winston Churchill had a motto at the beginning of World War II: "Keep buggering on." In that spirit, I think it's better to do something than to do nothing. **My advice is to try crazy ideas and innovative experiments. Something will come up.***



Steven Weinberg,
NP 1979

THANKS!!

DESY directorate:

- DESY Strategy Fund funded many of studies presented here

DESY technical groups:

- MVS (Vacuum Modification)
- MIN (Kicker, Beam Dump)
- D3 (radio protection advice)
- MEA (installation and Magnets)
- ZM1 (Construction Input)
- MKK (Power/Water)
- IPP (CAD integration)

DESY divisions

- MXL, MPY, MPY1 (from M), FLC (from FH)





BACKUP SLIDES

NON-LINEAR COMPTON SCATTERING: $e^- + n\omega \rightarrow e^- + \gamma$

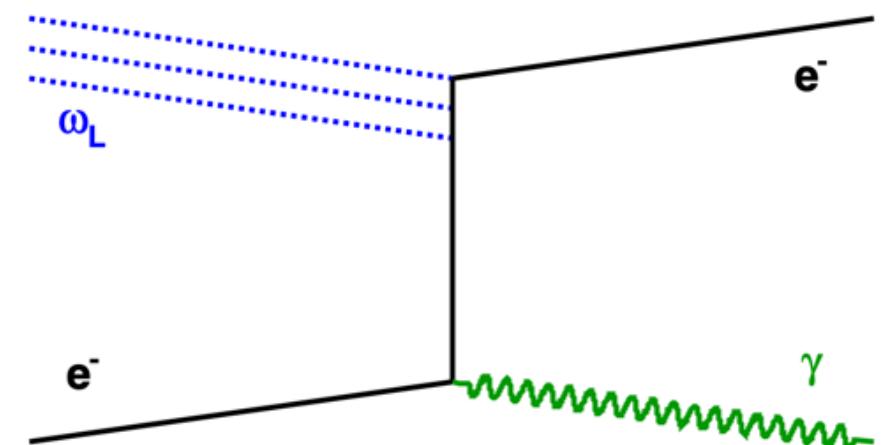
- Rate of high-intensity Compton scattering proportional to

$$\sum_n \delta^{(4)} \left[p_i + k \frac{\xi^3}{2\chi_i} + nk - p_f - k \frac{\xi^3}{2\chi_f} - k_f \right]$$

- Even for small n expect shift of Compton edge due to effective increase of electron rest mass

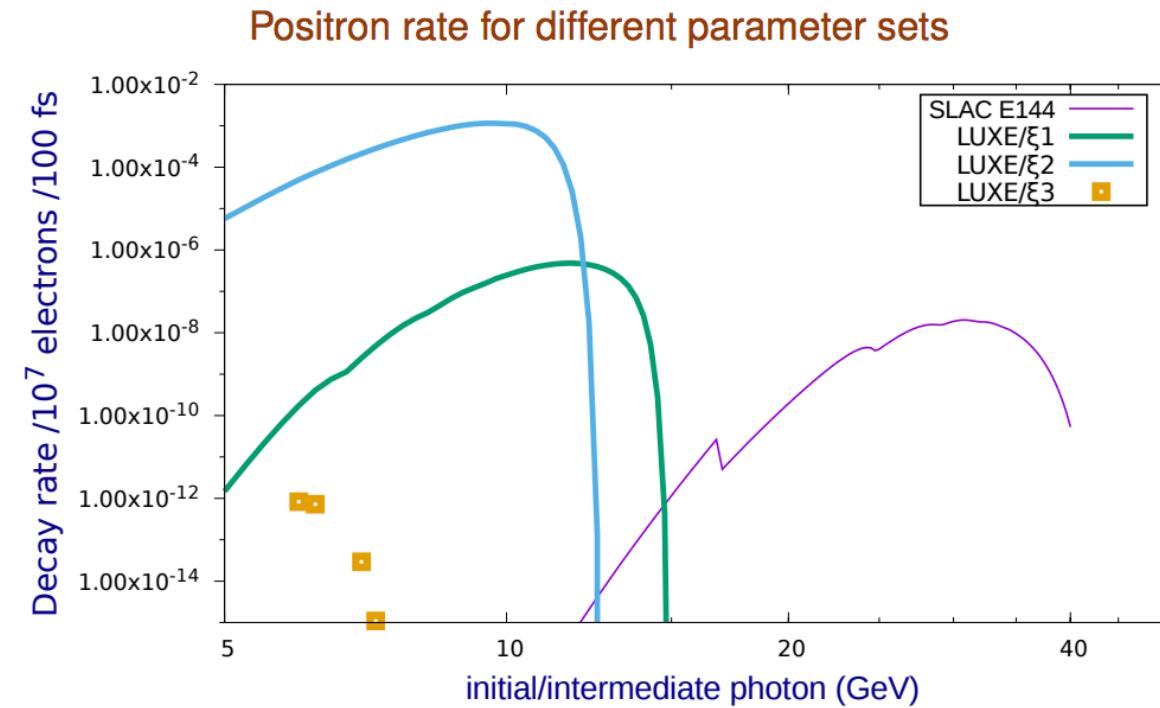
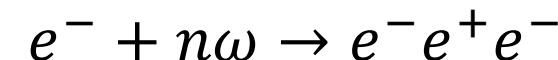
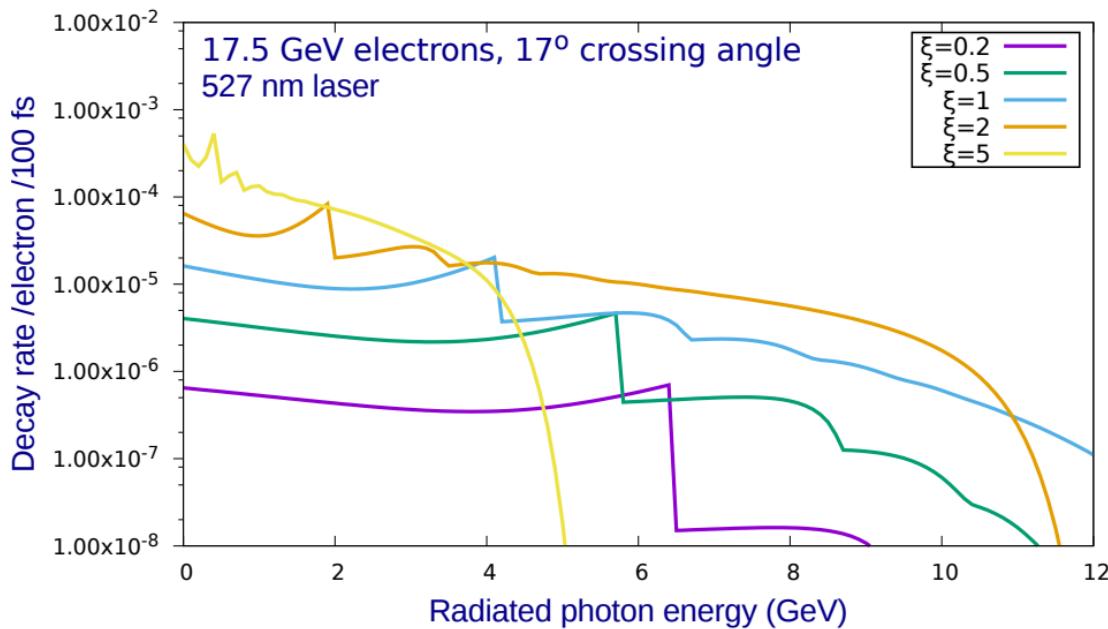
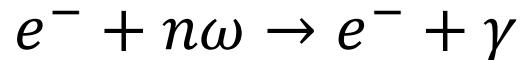
$$p_i + k \frac{\xi^3}{2\chi_i} \rightarrow p_i^2 = m^2(1 + \xi^2)$$

- Has never been observed





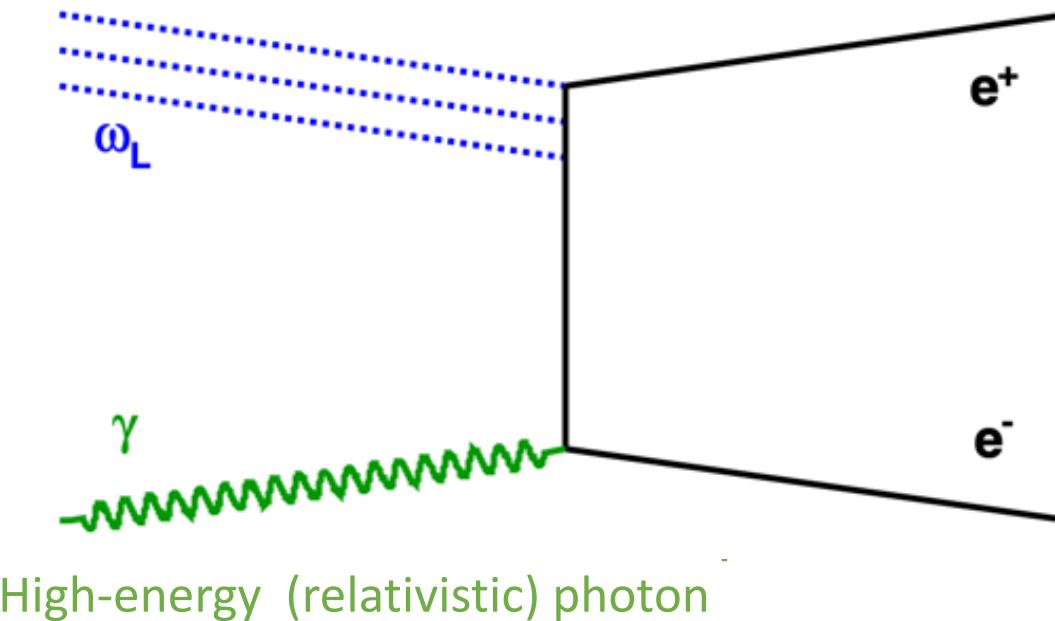
MEASUREMENTS OF MASS SHIFT AND TRIDENTS



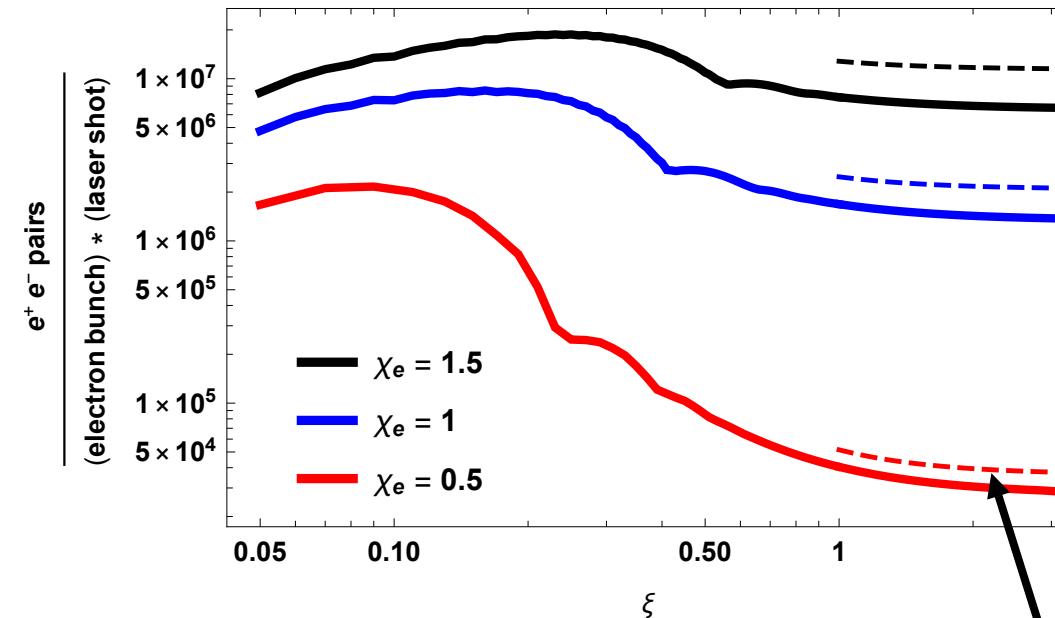
Plots from A. Hartin, IJMPA 33, 1830011 (2018)

ABSORBING LIGHT WITH LIGHT

Low-energy photons from laser

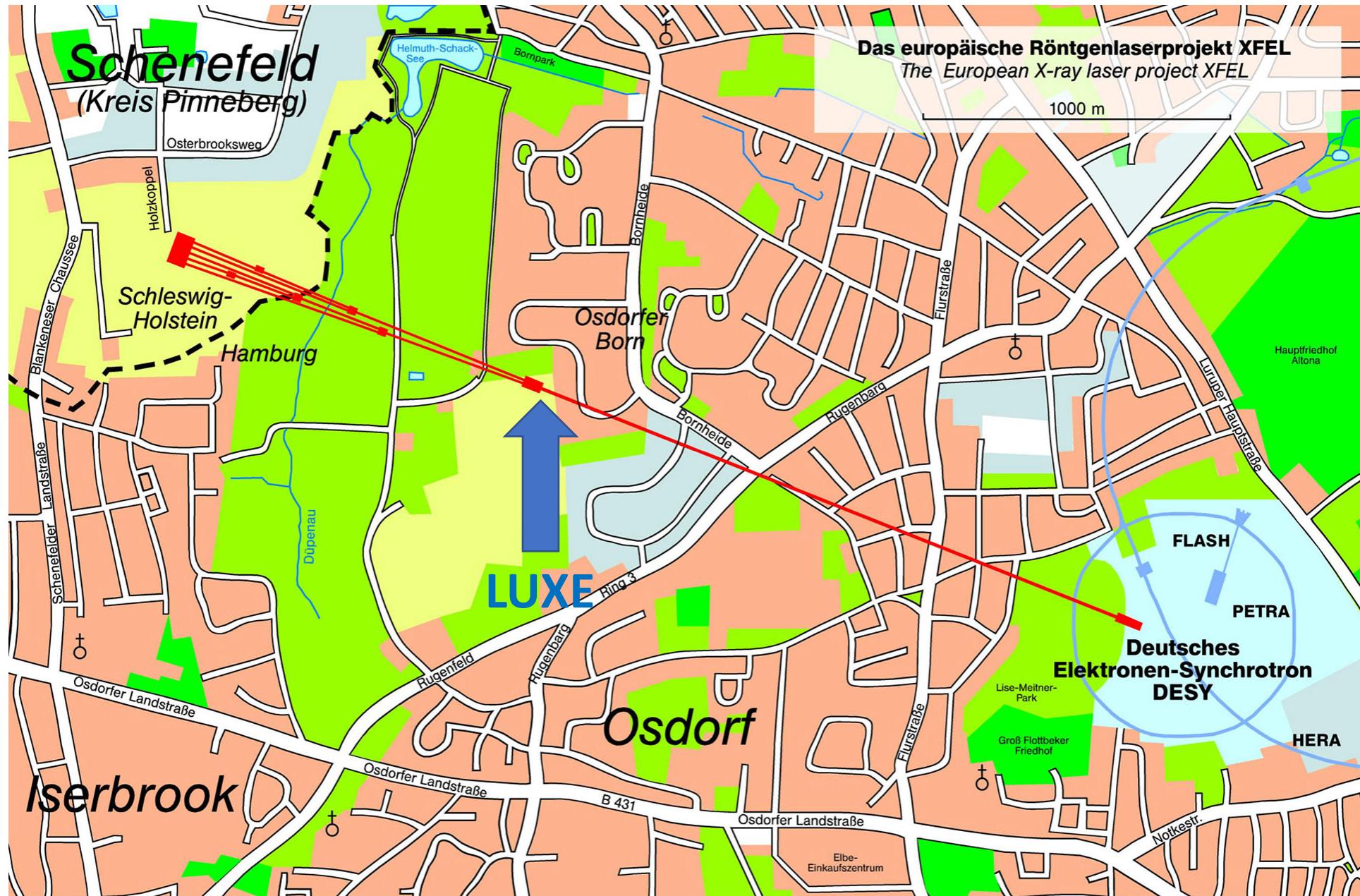


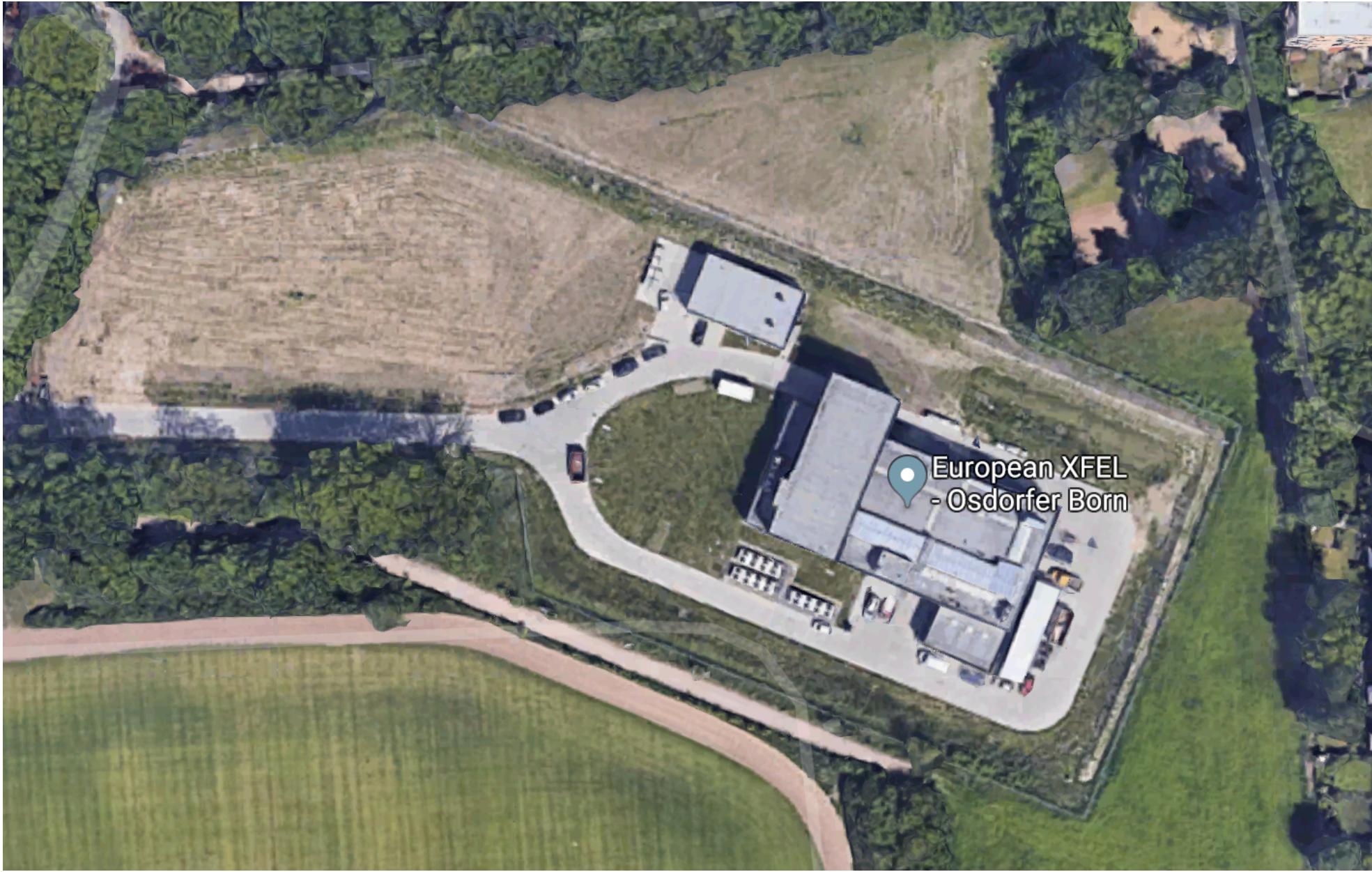
$$E_e = 17.5 \text{ GeV}, e^- \text{ bunch} = 6 \times 10^9, \frac{X}{X_0} = 0.01, \text{ Laser shot} = 35 \text{ fs}$$



Asymptotic limit

- Use spectrum of high energy photons created via Bremsstrahlung
 - Full calculation agrees with asymptotic limit for $\xi > 1$ and $\chi \lesssim 1$





BEAM DUMP

Beam needs to be safely dumped, design well advanced

F. Burkart, M. Schmitz (DESY)

