

*Christoph Englert*

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# Three avenues for Higgs phenomenology

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*Particle Physics Seminar*

*Birmingham*

*16/01/2019*

# Overview

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- ▶ *Improving the expected: SM-like Higgs couplings*
  - ▶ lifting degeneracies in coupling space for expected uncertainties with adversarial machine learning
  - ▶ .....

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  - ▶ Higgs sector CP violation
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- ▶ *Constraining/observing the unexpected:*
  - ▶ Higgs sector CP violation
  - ▶ .....
- ▶ *Closing in on new physics in the Higgs sector*
  - ▶ *di-Higgs production as a probe of new physics*
  - ▶ .....

# “Yang-Mills theories had to be right”

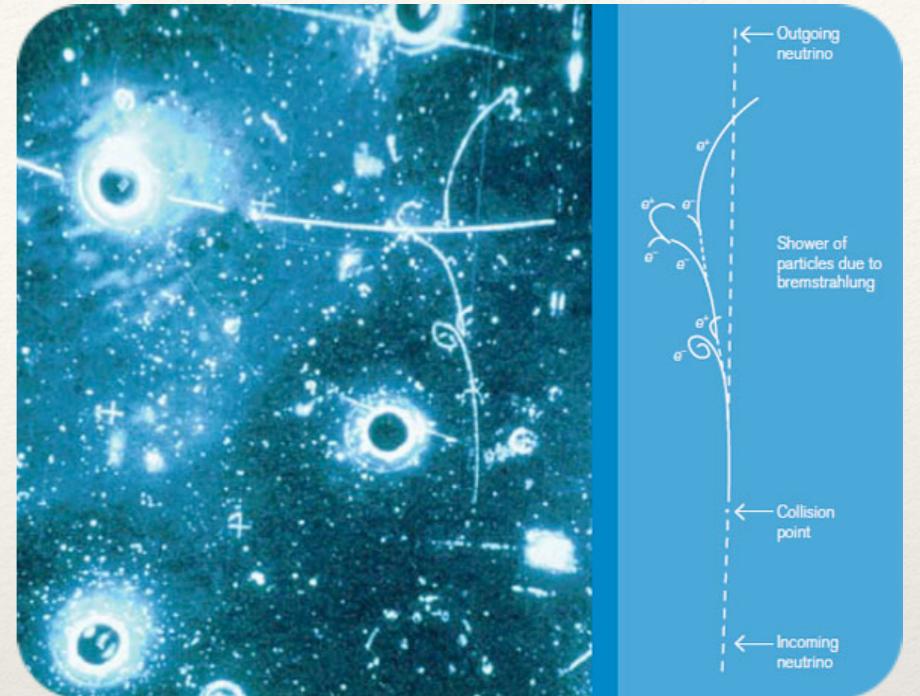
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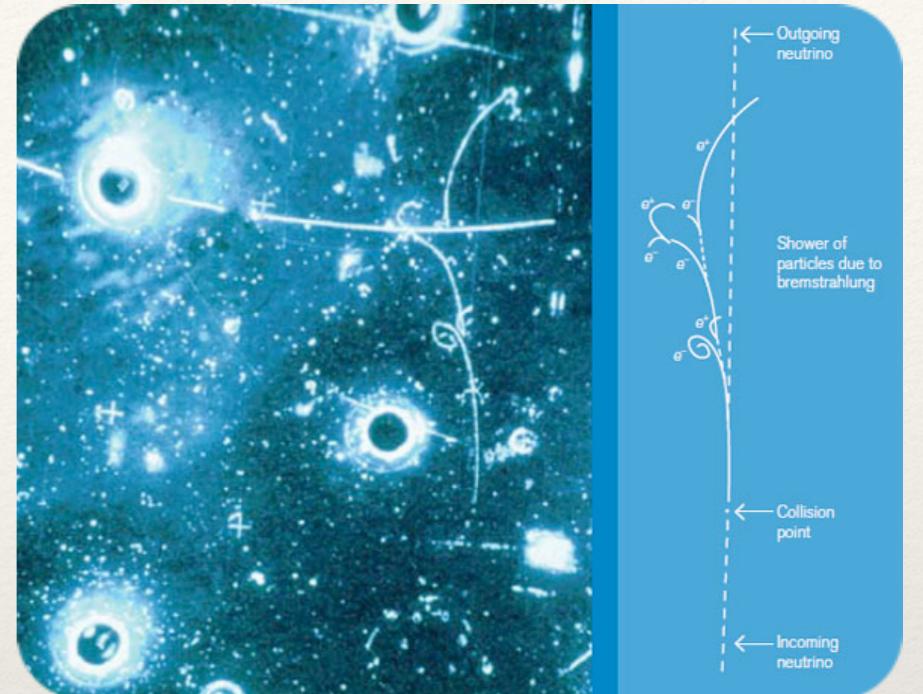
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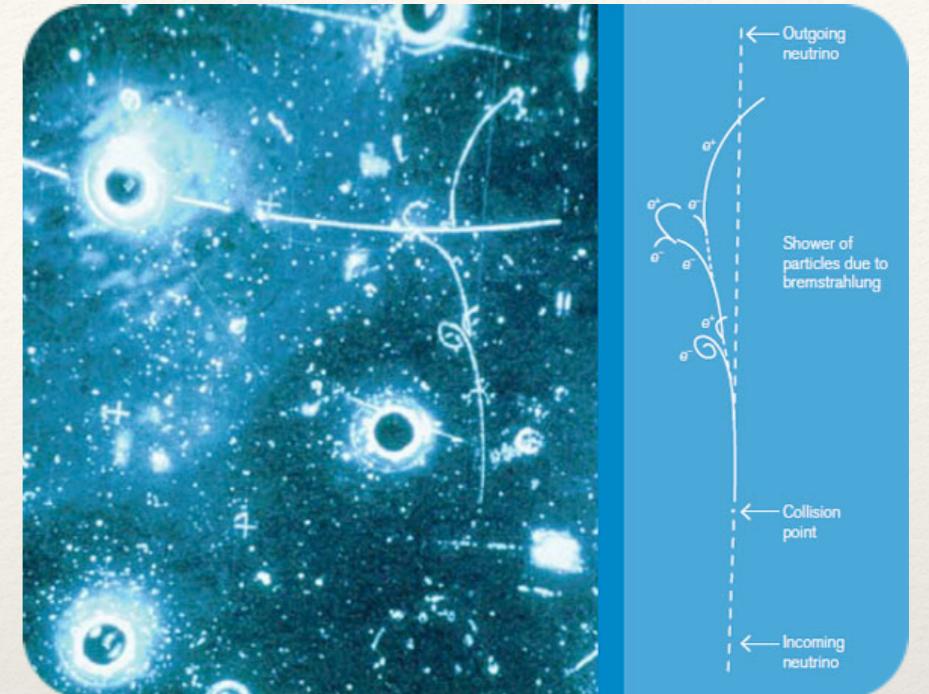
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  - 👉 “spontaneous” symmetry breaking



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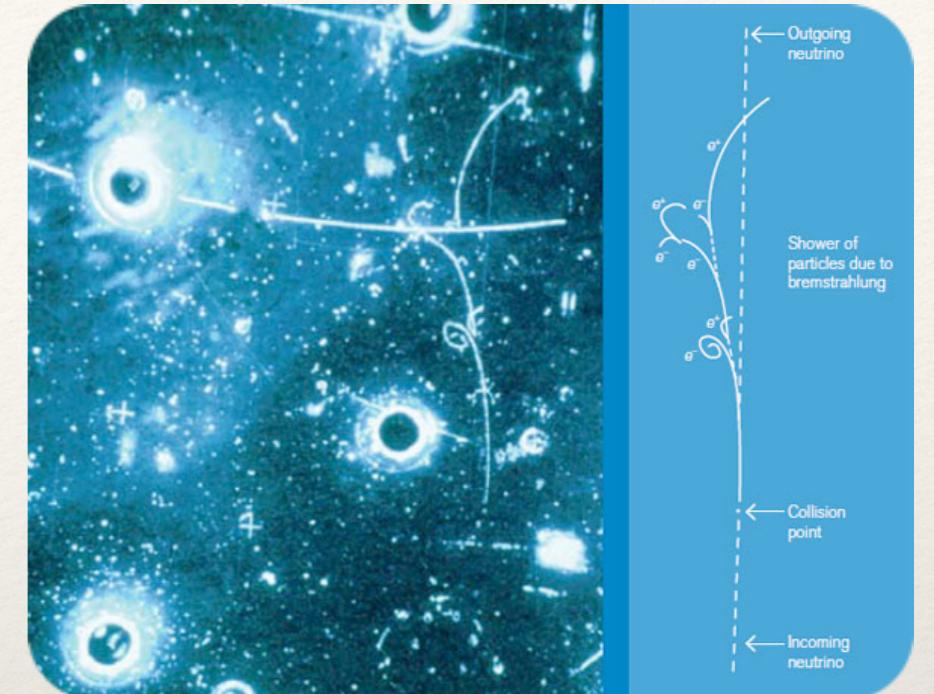
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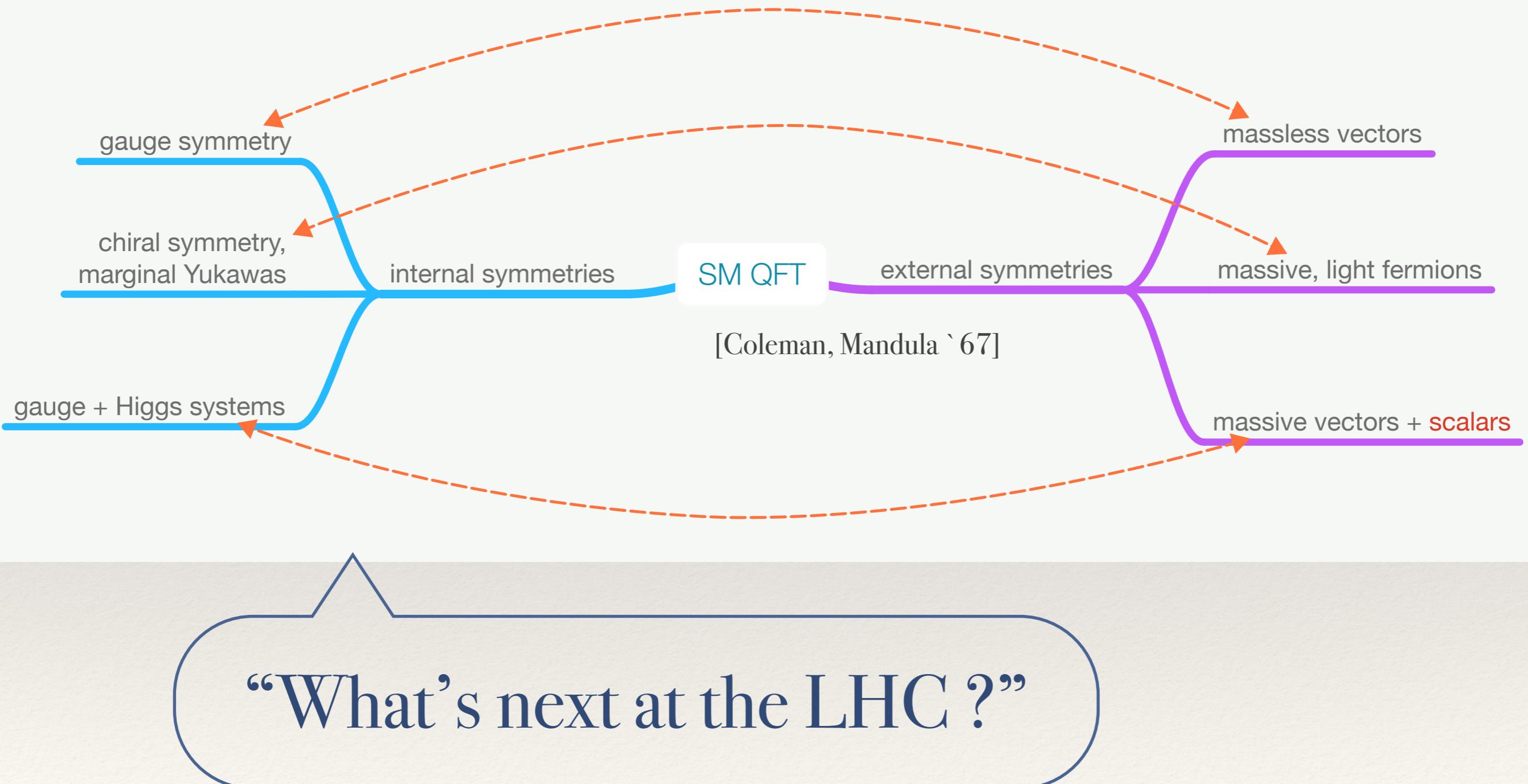
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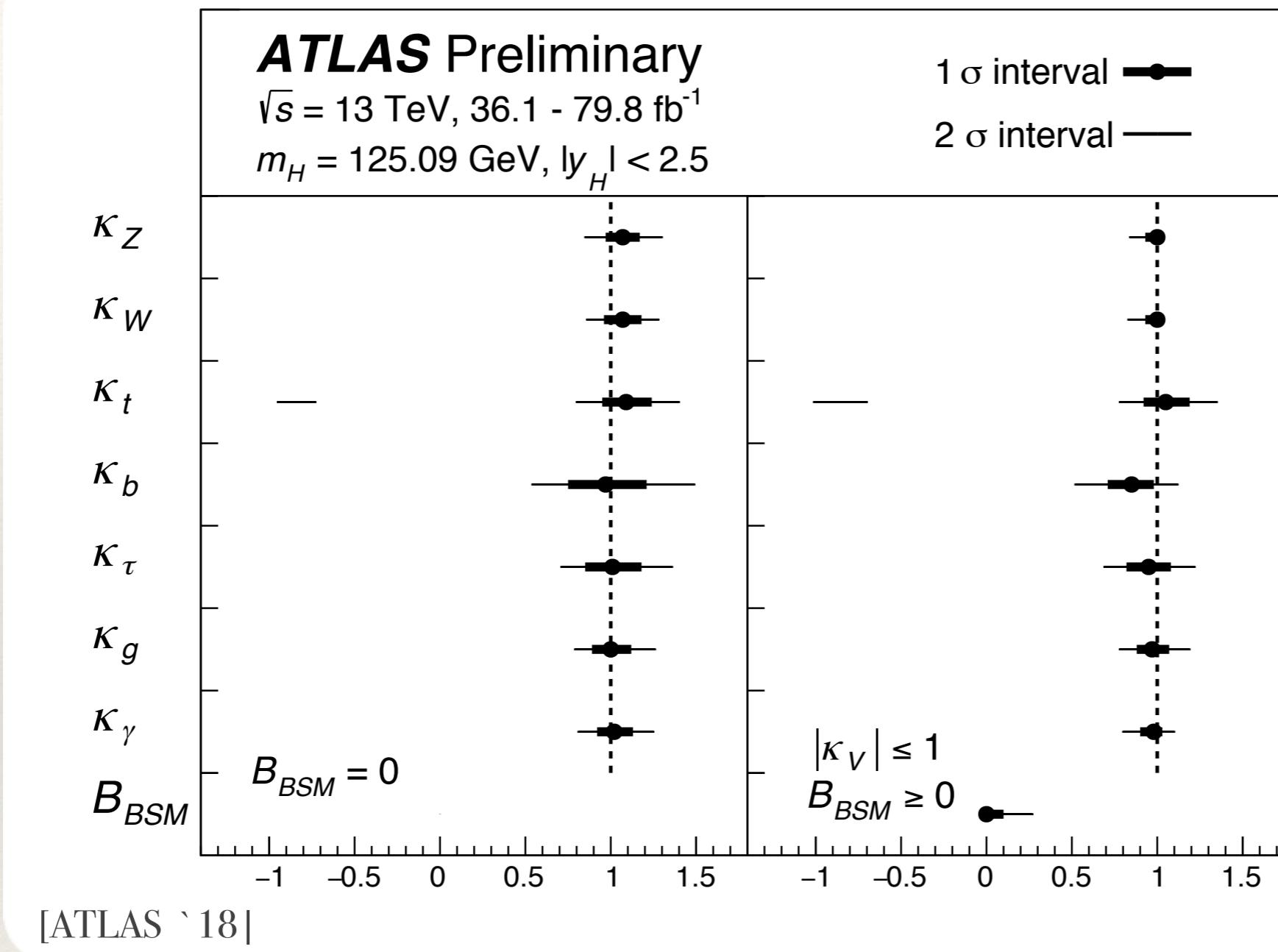
- non-linear realisation of gauge symmetry in a Yang Mills+scalar sector is compatible with  $\langle H \rangle \neq 0$ 
  - 👉 “spontaneous” symmetry breaking
- massive gauge bosons, but no ghost problems at small distances
  - 👉 renormalizability, probability conservation



# The Standard Model: taking stock



# Status of LHC measurements



- everything is consistent with the SM Higgs hypothesis (so far)  
but what are the implications for new physics?

# Fingerprinting the lack of new physics

the SM is flawed

no evidence for exotics

coupling/scale  
separated BSM physics

## Effective Field Theory

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \dots$$

[Buchmüller, Wyler '87]

[Hagiwara, Peccei, Zeppenfeld, Hikasa '87]

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59 B-conserving operators  $\otimes$  flavor  $\otimes$  h.c., d=6  
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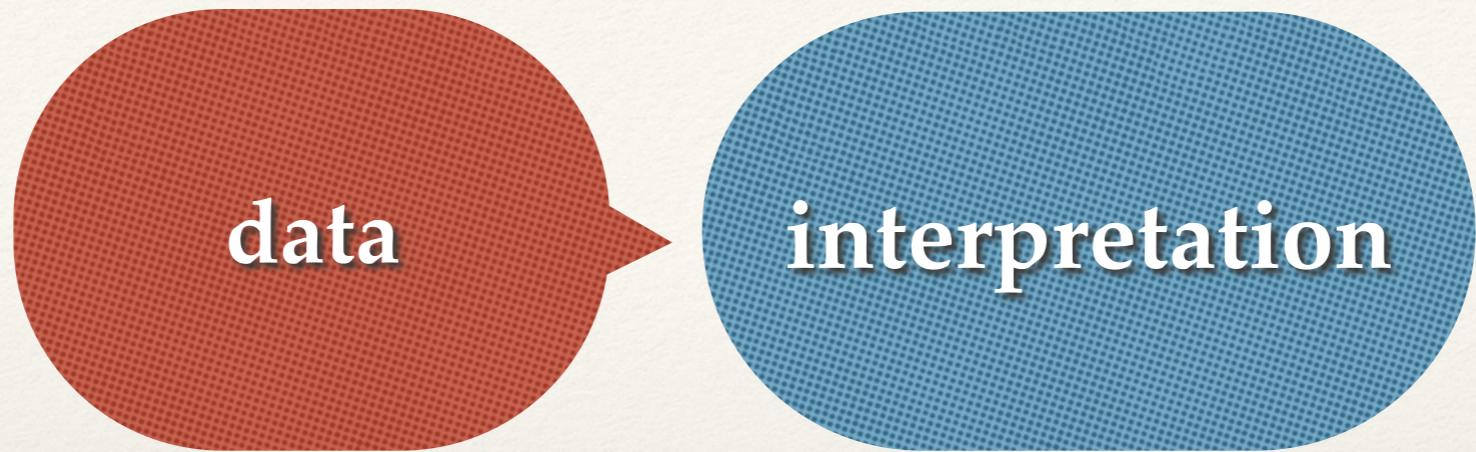
## concrete models

- extended SMEFT
- ( $\mathbb{C}$ ) Higgs portals
- 2HDMs
- (N)MSSM
- compositeness....

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# Effective theory

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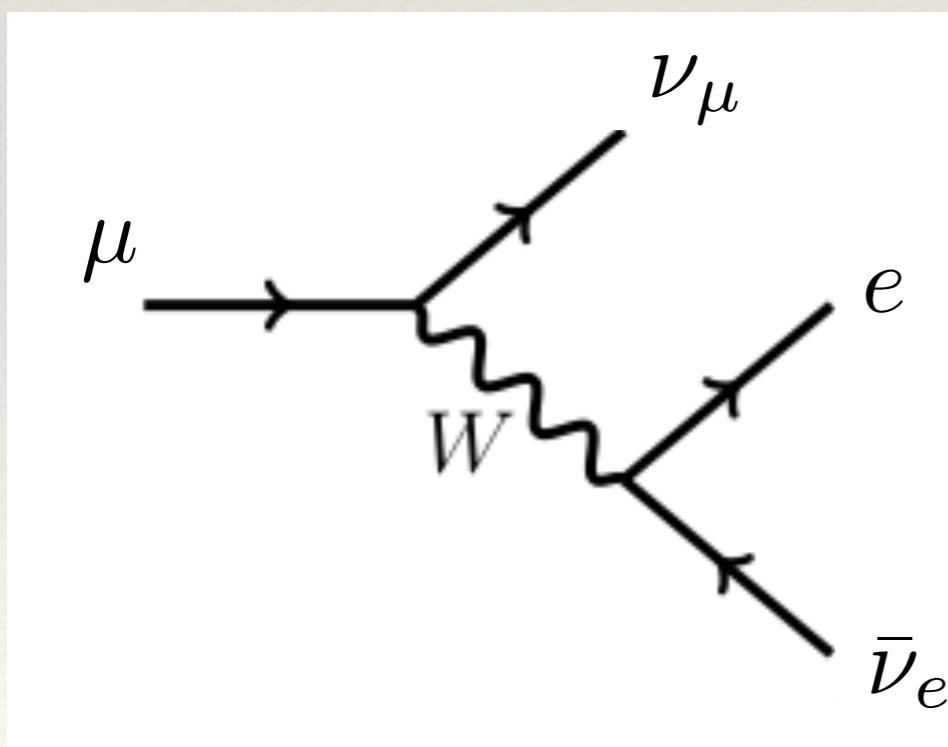


# Effective theory





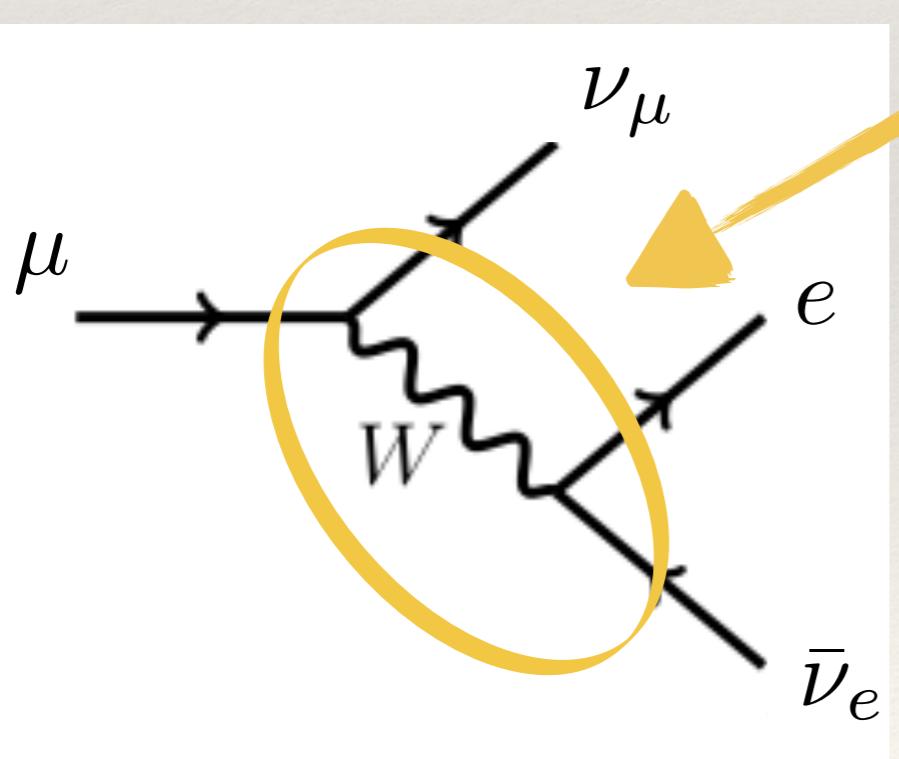
## Weak decay



# Effective theory



$$\frac{1}{t - m_W^2} = -\frac{1}{m_W^2} - \frac{t}{m_W^4} + \dots \quad |t| \ll m_W^2$$

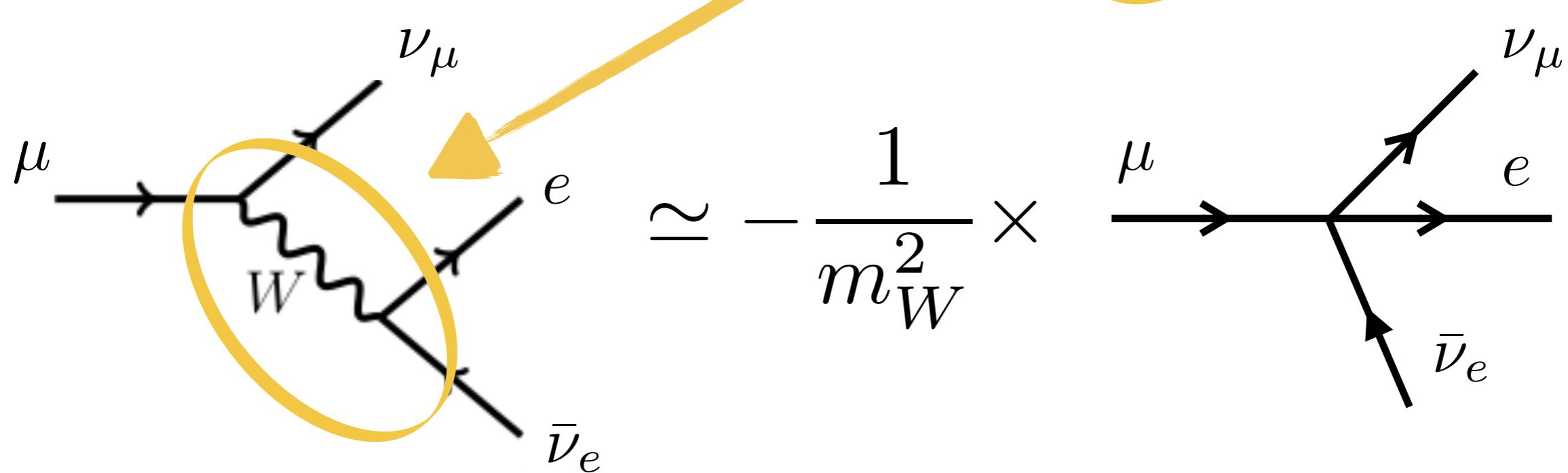


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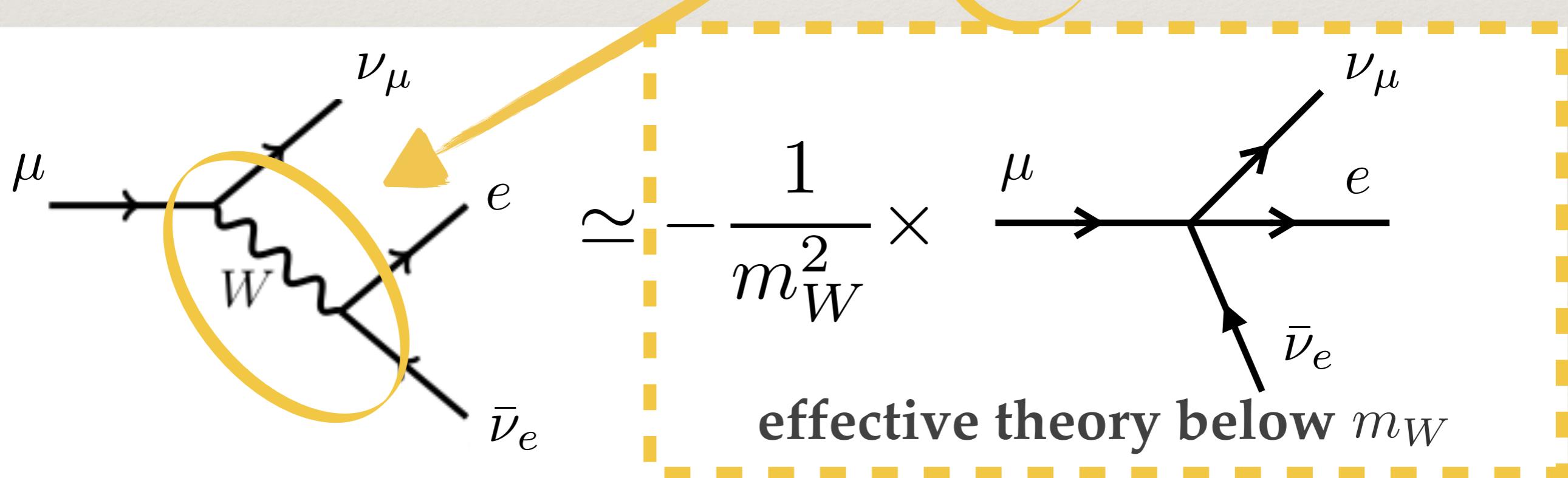


# Effective theory



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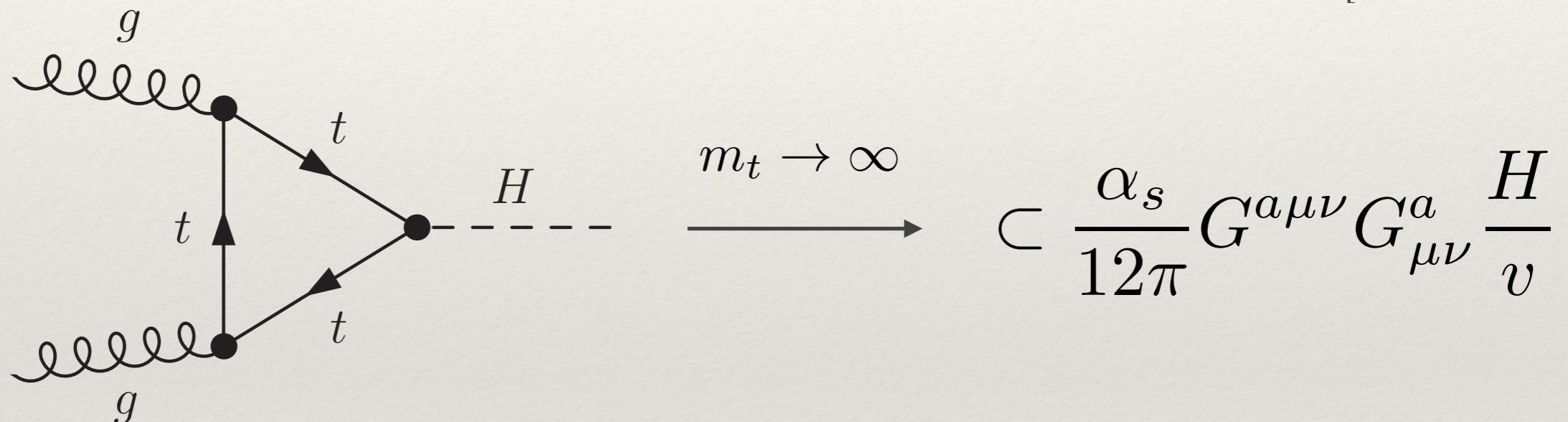
# SM-like couplings

- large number of unconstrained EFT parameters lead to phenomenological degeneracies = “blind directions”
- one of the most prominent and relevant for Higgs physics

[Ellis et al. '76]

[Vainstein et al. '70]

....



contact ggH interactions mask top Yukawa measurements

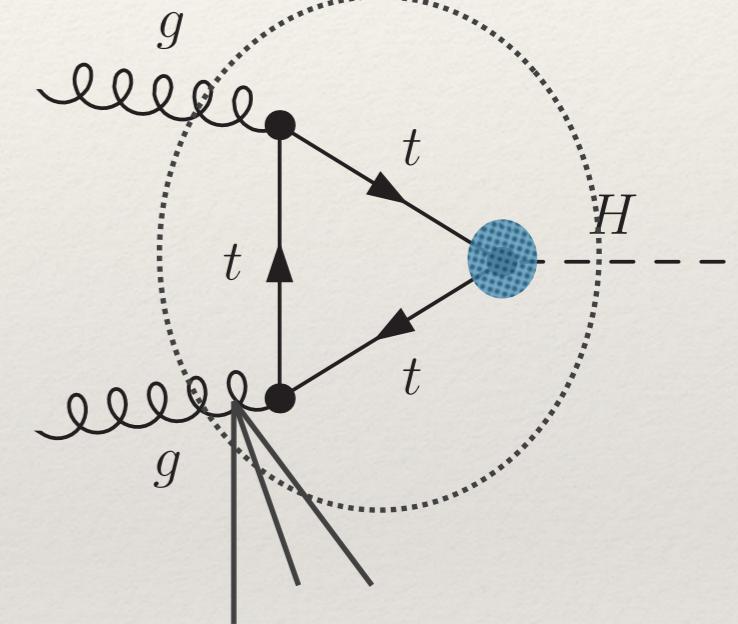
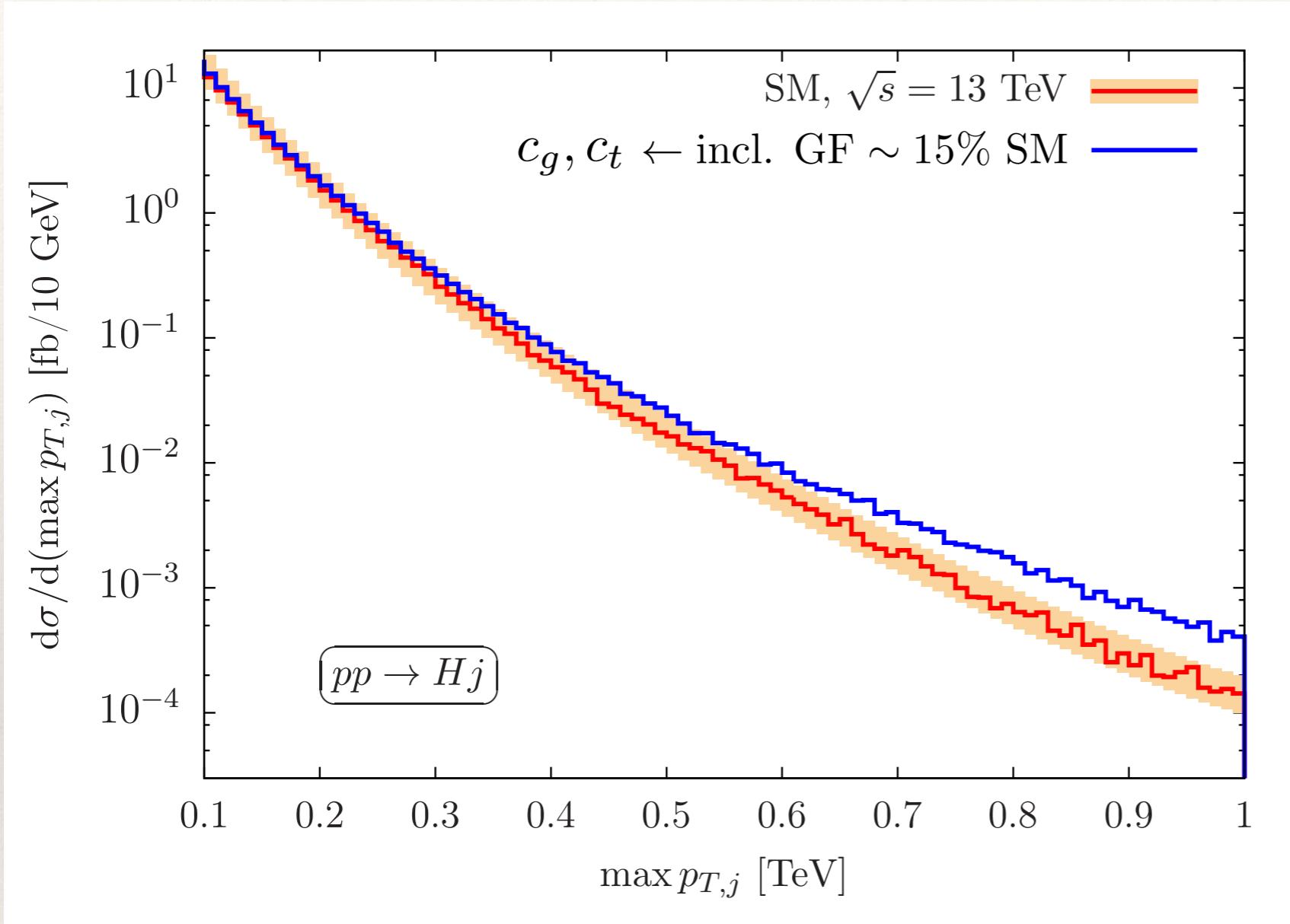
- way out: resolve  $C_0$  function for  $p_T(H) \gtrsim m_t$  with one or more jets

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# SM-like couplings

# A word of caution



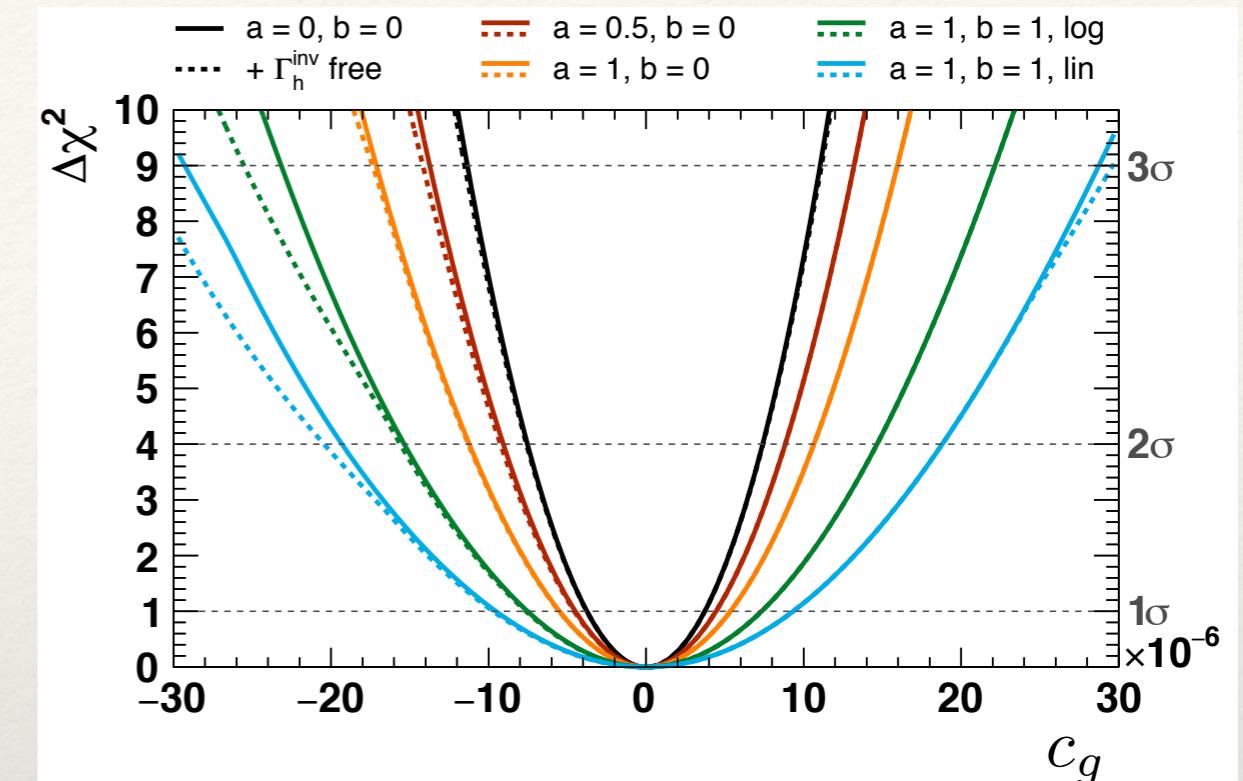
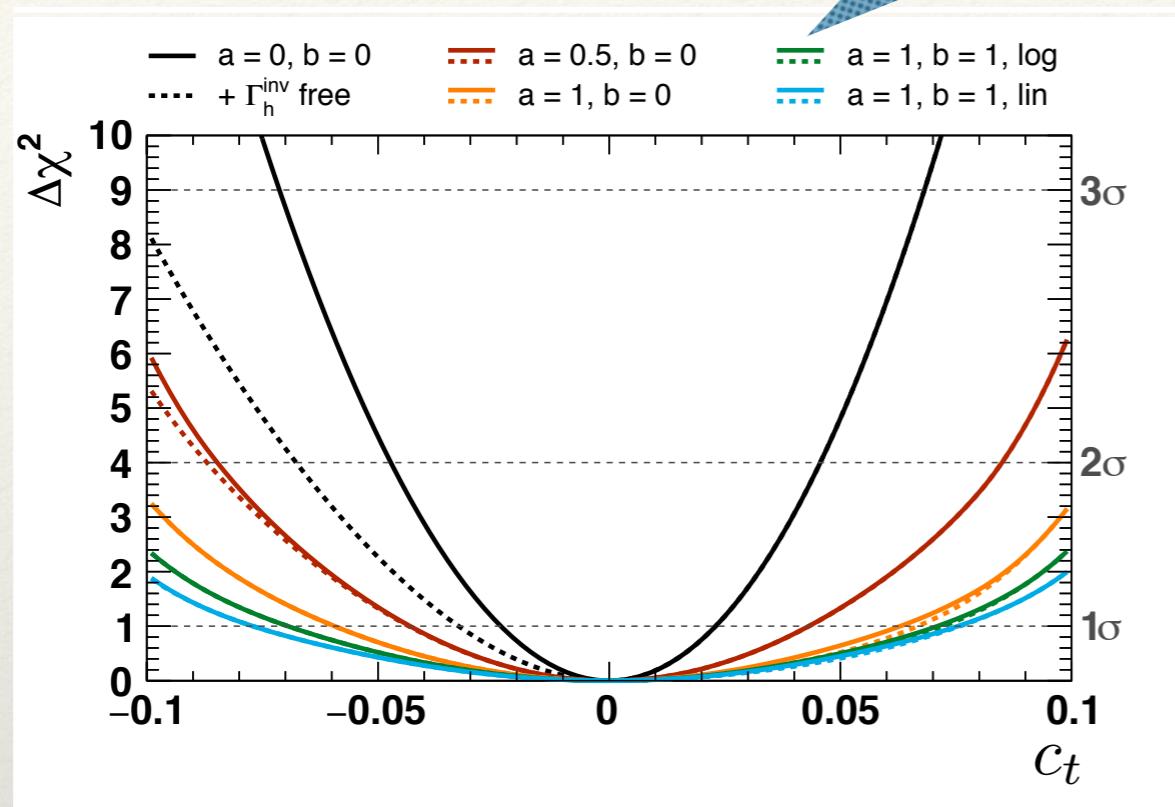
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steer  $p_T(H)$  shape  
uncertainty

[CE, Kogler, Schulz, Spannowsky '17]

## Role of uncertainties



- comparably small impact of tail uncertainties  
(lin vs log  $\sim 35\%$  different shape uncertainty at 150 GeV  $p_T$ )
- decoupled (non-resonant) new physics perturbatively constrained at relatively low transverse momentum**

“fit will always pick region where null hypothesis is under good control”

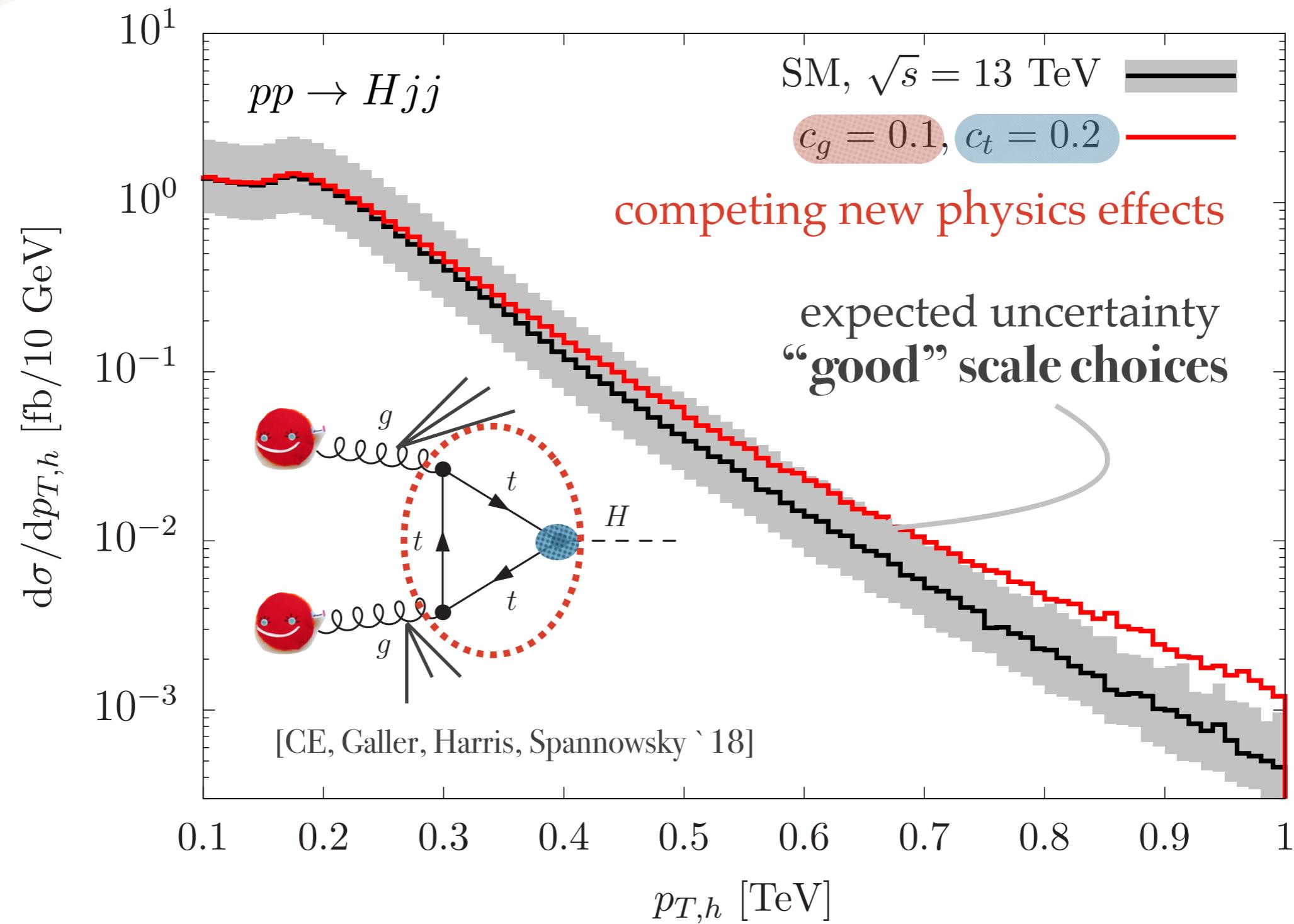
similar conclusion hold for more abundant top final states

large stats!

[CE, Moore, Nordstrom, Russell '16]

# SM-like couplings

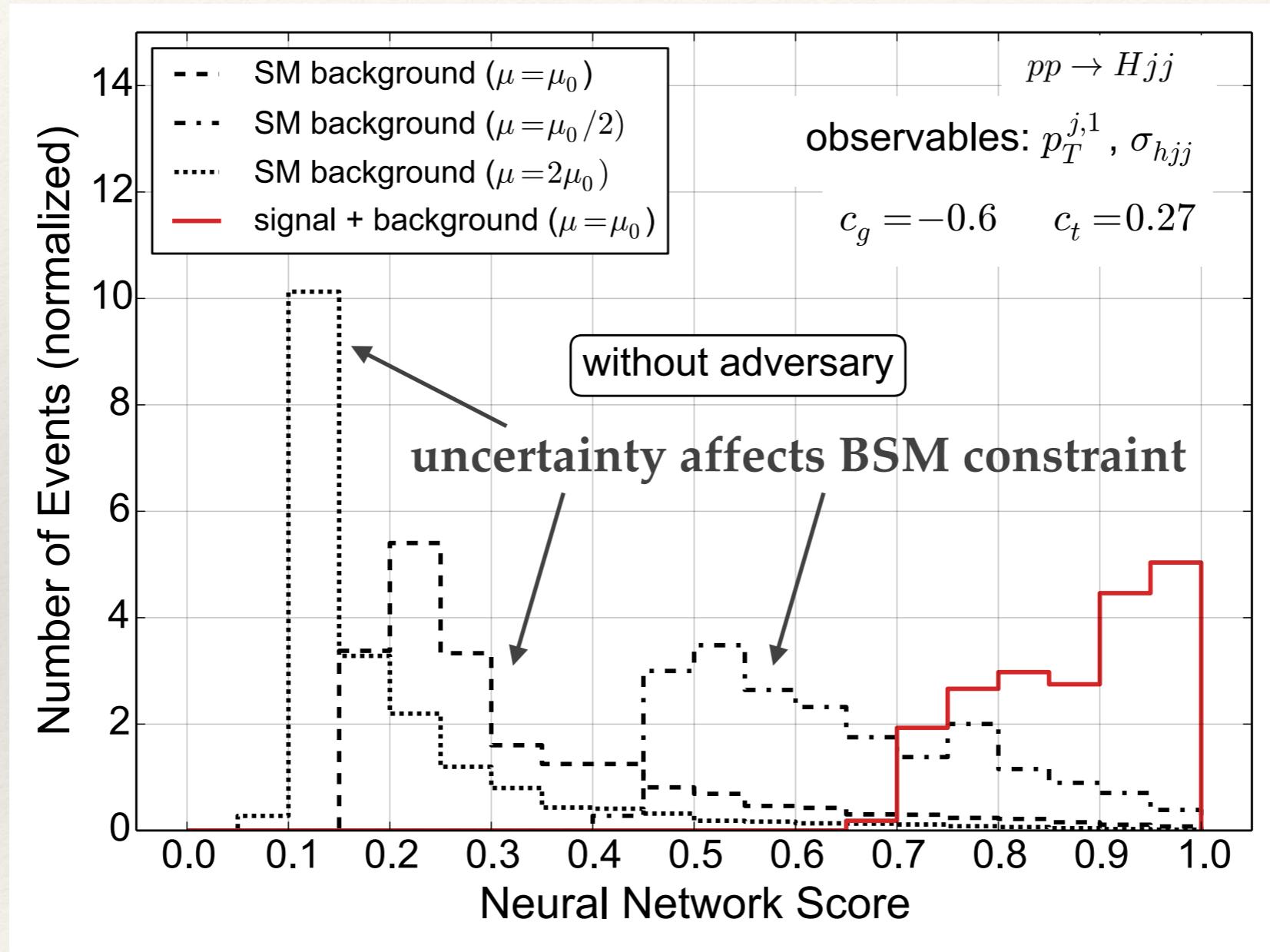
- more kinematic information for  $H+2j$ , which is particularly promising, unfortunately  $m_t=\infty$  SM limit accidentally good [Del Duca et al. '03]



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[CE, Galler, Harris, Spannowsky '18]

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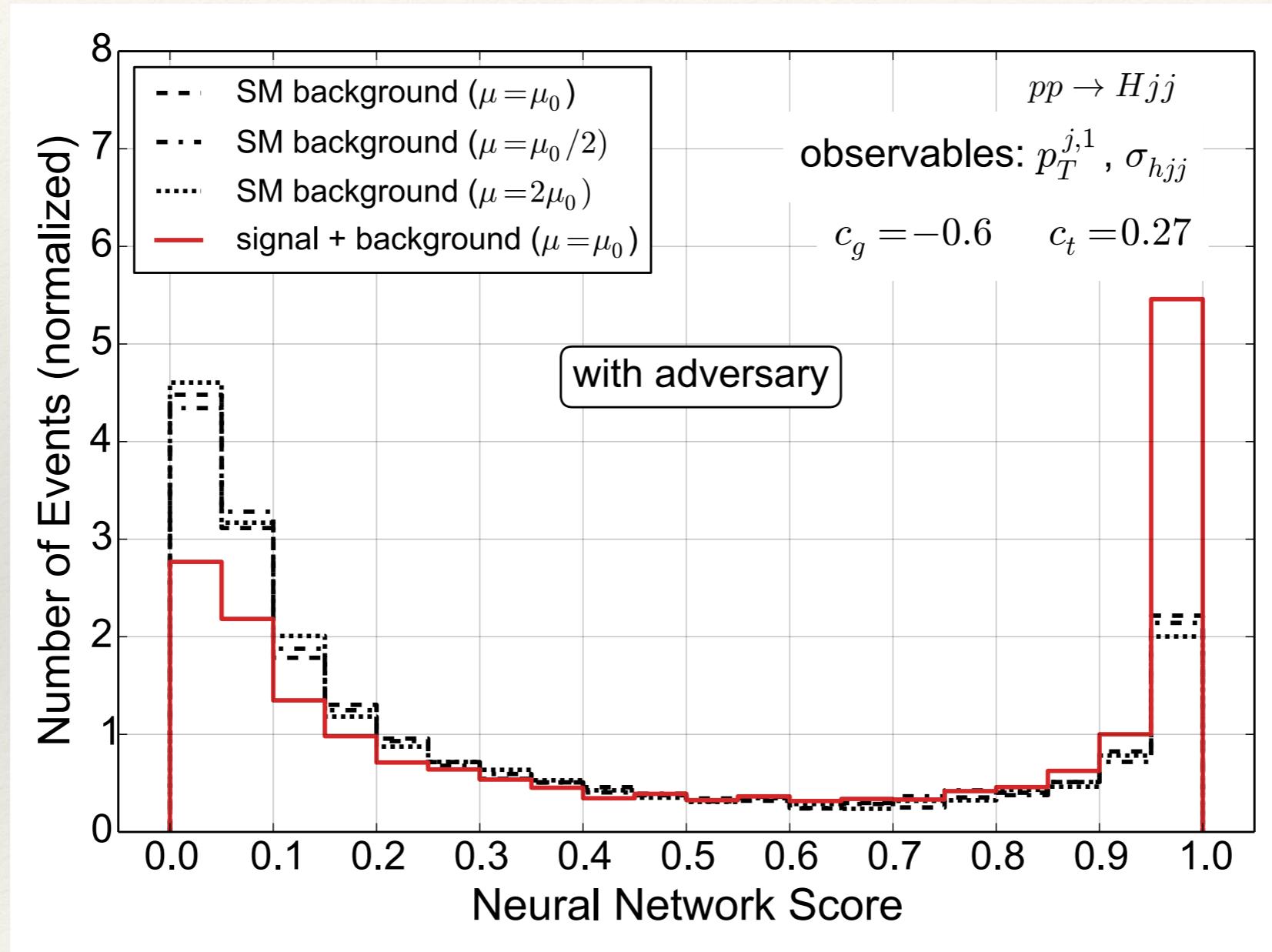


neural net learns regions that are sensitive to uncertainty....

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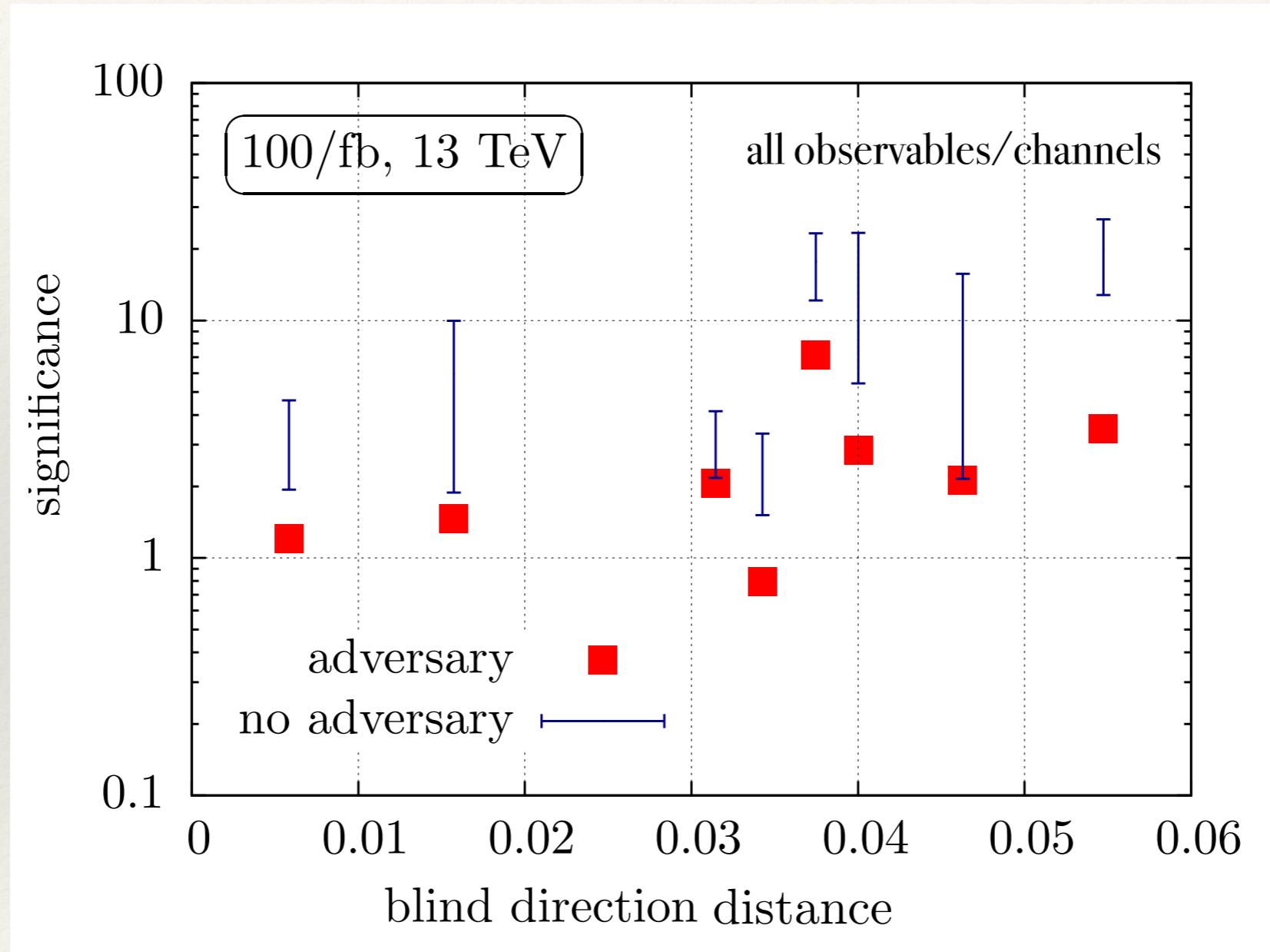


... and can be forced to avoid them → **most robust constraints**

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# CP violation

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MC  
perturbativity

unitarity...

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  - MC perturbativity
  - unitarity...
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- ▶ qualitatively different for CP-violation:

$$\frac{c_i}{\Lambda^2}$$

$\sim \text{dim } 6$

naive  
perturbative  
power counting

$$\frac{c_i^2}{\Lambda^4}$$

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- ▶ only genuinely CP-sensitive observables carry information

signed  $\Delta\phi_{jj}$ , asymmetries, ....

- ▶ every CP-even observable carries information

cross sections, widths, pT spectra...

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# CP violation

[Bernlochner, CE, Hays, Lohwasser, Mildner, Pilkington, Price, Spannowsky '18]

- the linearised upshot

$$\begin{aligned} O_{H\tilde{G}} &= H^\dagger H G^{a\mu\nu} \tilde{G}_{\mu\nu}^a, \\ O_{H\tilde{W}} &= H^\dagger H W^{a\mu\nu} \tilde{W}_{\mu\nu}^a, \\ O_{H\tilde{B}} &= H^\dagger H B^{\mu\nu} \tilde{B}_{\mu\nu}, \\ O_{H\tilde{W}B} &= H^\dagger \tau^a H B_{\mu\nu} \tilde{W}^{a\mu\nu} \end{aligned}$$

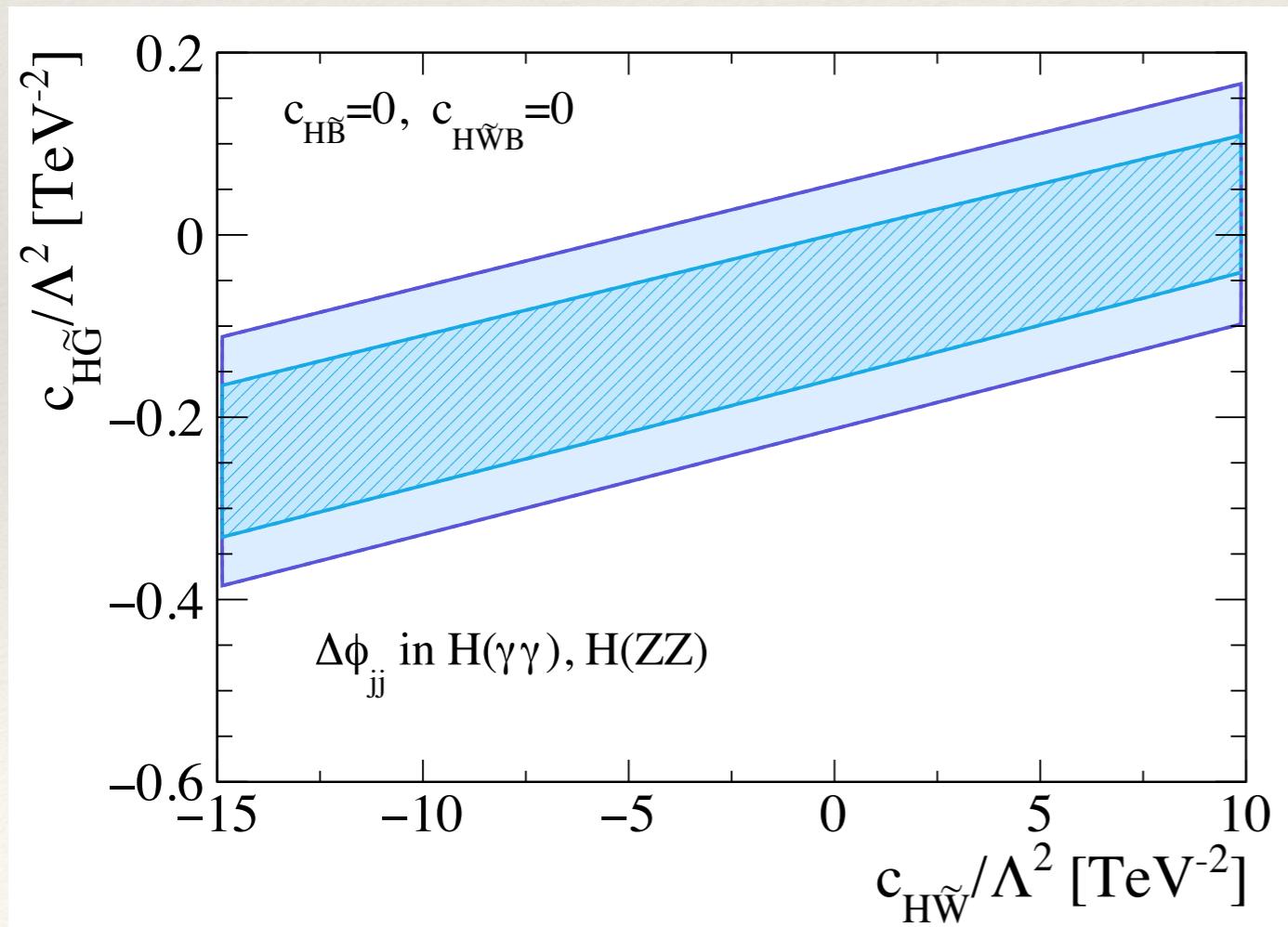
+

top quark

$$\sim \frac{\alpha_s}{8\pi v} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} h = \tilde{O}_G$$

Yukawa phases

*...ignore them for now...*



- fit uses ATLAS results for 4 leptons,  $\gamma\gamma$   
[ATLAS 1708.02810; 1802.04146]
- small stats/observables = blind directions for decay vs production
- non-significant asymmetry  $0.3 \pm 0.2$

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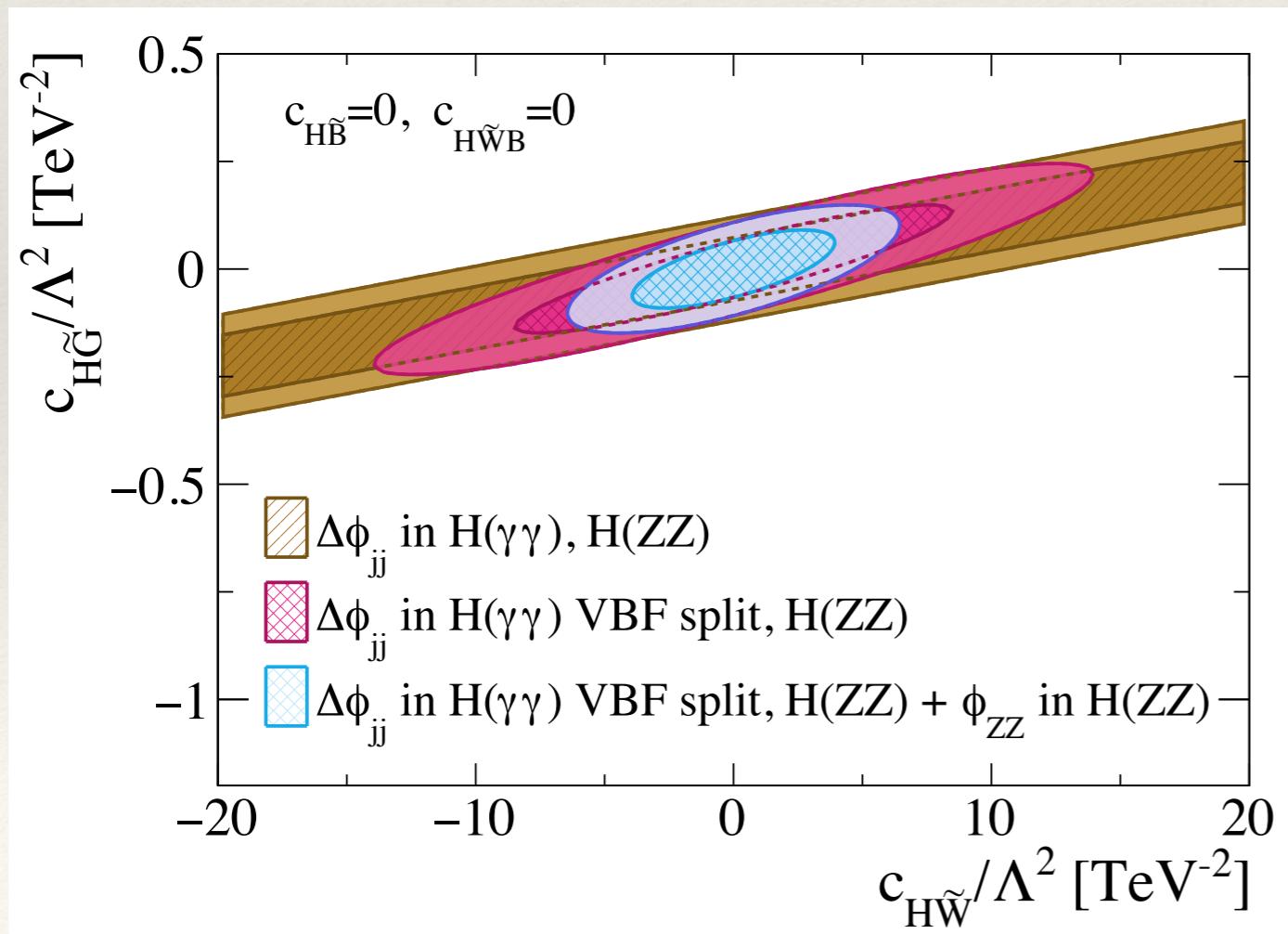
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- could already narrow this down at 36/fb by using
  - GF vs VBF selections
  - lepton decay plane angles

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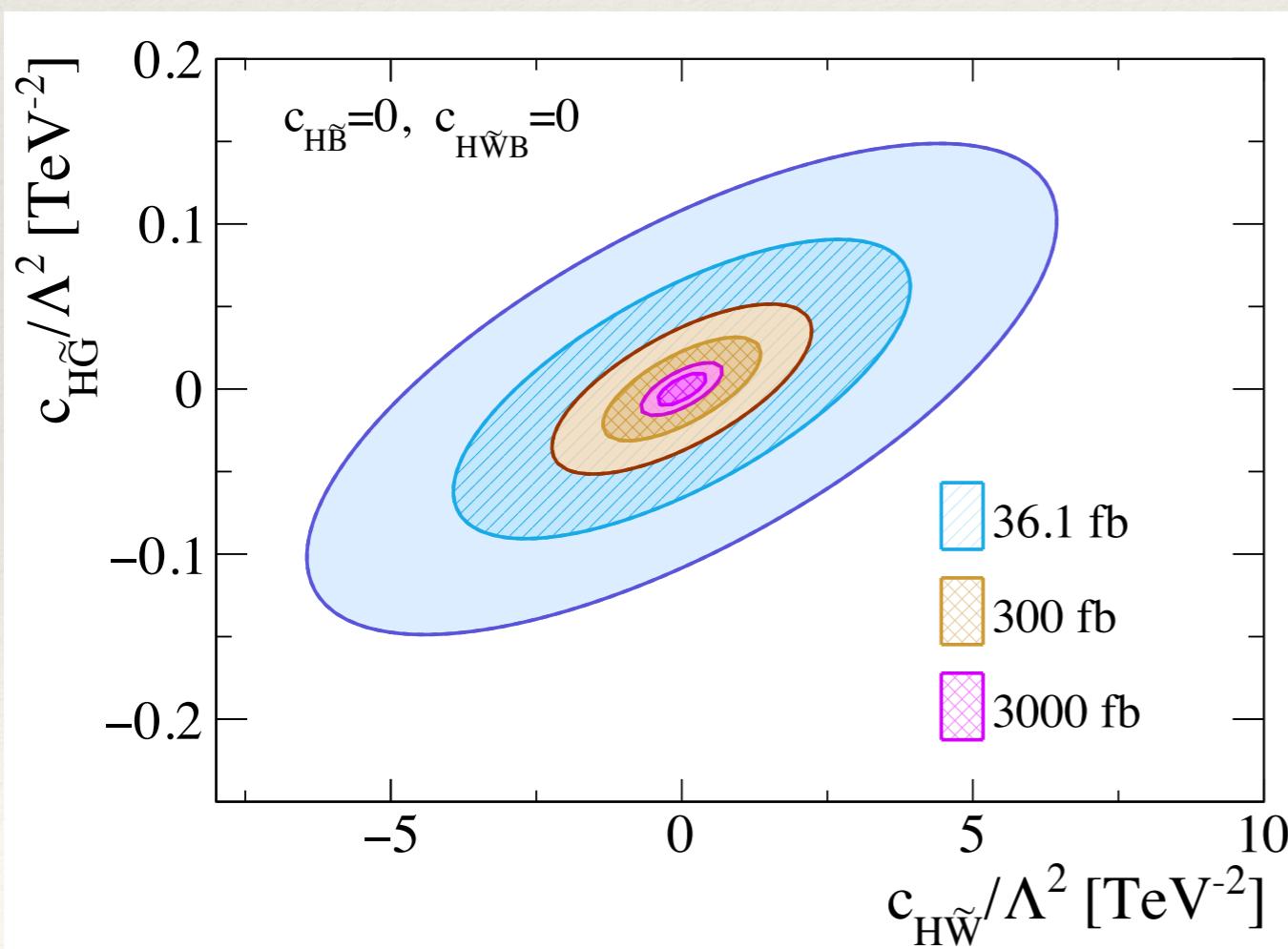
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LHC and HL-LHC  
extrapolations

Coefficient [TeV $^{-2}$ ]	36.1 fb $^{-1}$	300 fb $^{-1}$	3000 fb $^{-1}$
$c_{H\tilde{G}}/\Lambda^2$	[-0.19, 0.19]	[-0.067, 0.067]	[-0.021, 0.021]
$c_{H\tilde{W}}/\Lambda^2$	[-11, 11]	[-3.8, 3.8]	[-1.2, 1.2]
$c_{H\tilde{B}}/\Lambda^2$	[-5.9, 5.9]	[-2.1, 2.1]	[-0.65, 0.65]
$c_{H\tilde{W}B}/\Lambda^2$	[-14, 14]	[-4.9, 4.9]	[-1.5, 1.5]

# CP violation

[CE, Galler, Pilkington, Spannowsky in prep.]

- lifting top-specific blind directions

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top Yukawa phase

[Del Duca et al. '03]

- $m_t = \infty$  SM limit accidentally good
- split GF selection into  $m_t$ -related Higgs pT threshold  $\sim 150$  GeV

large stats / kin.  
coverage necessary

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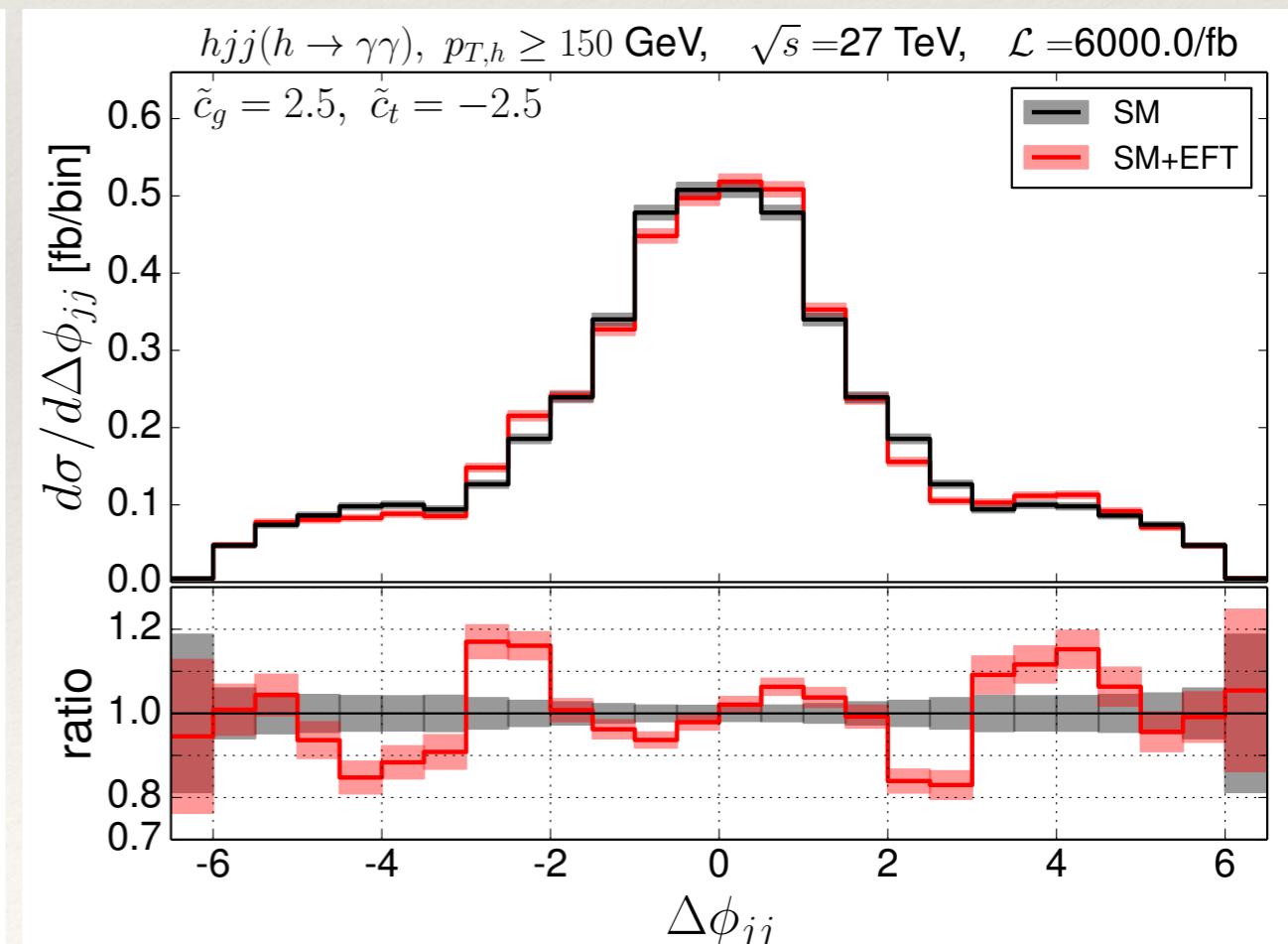
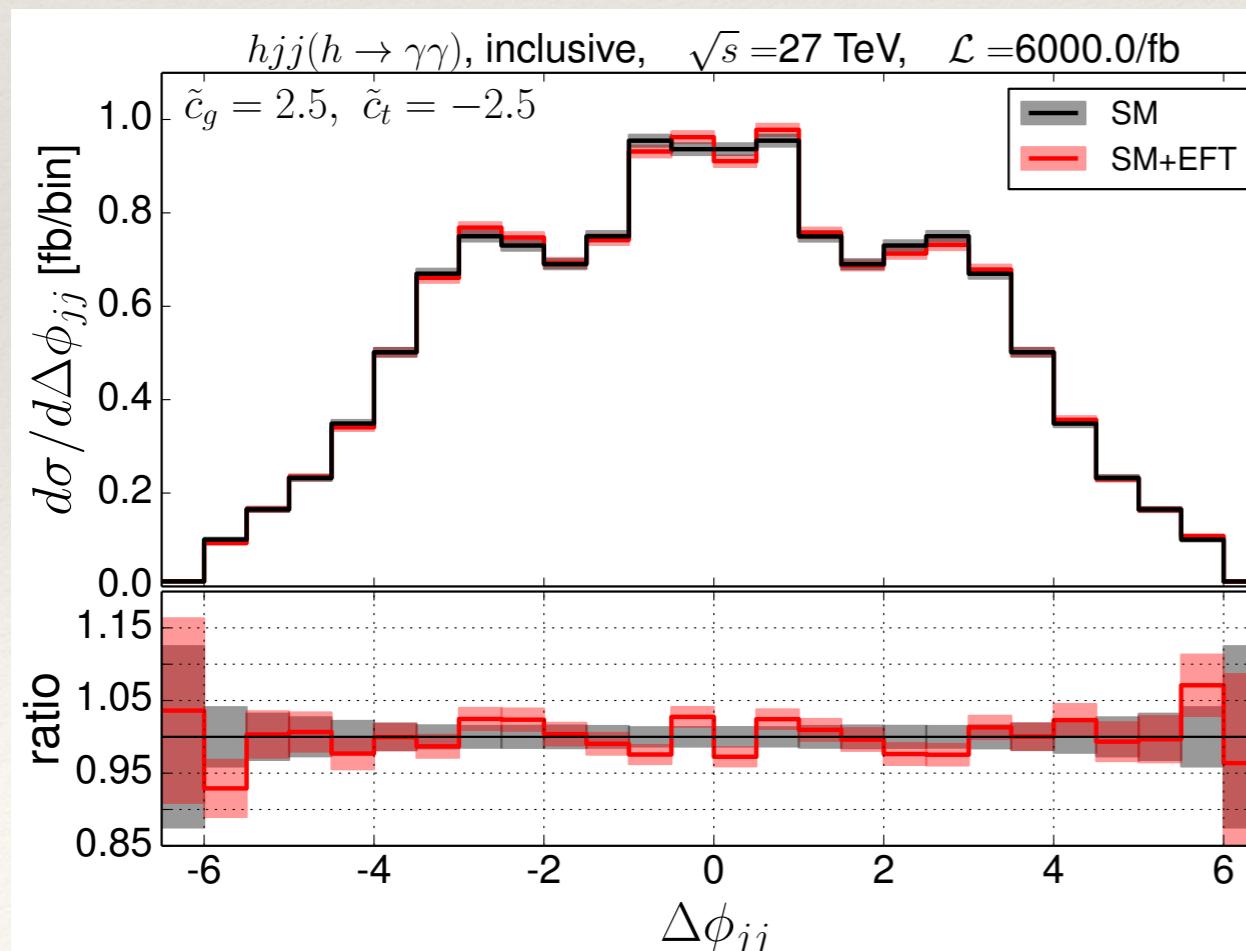
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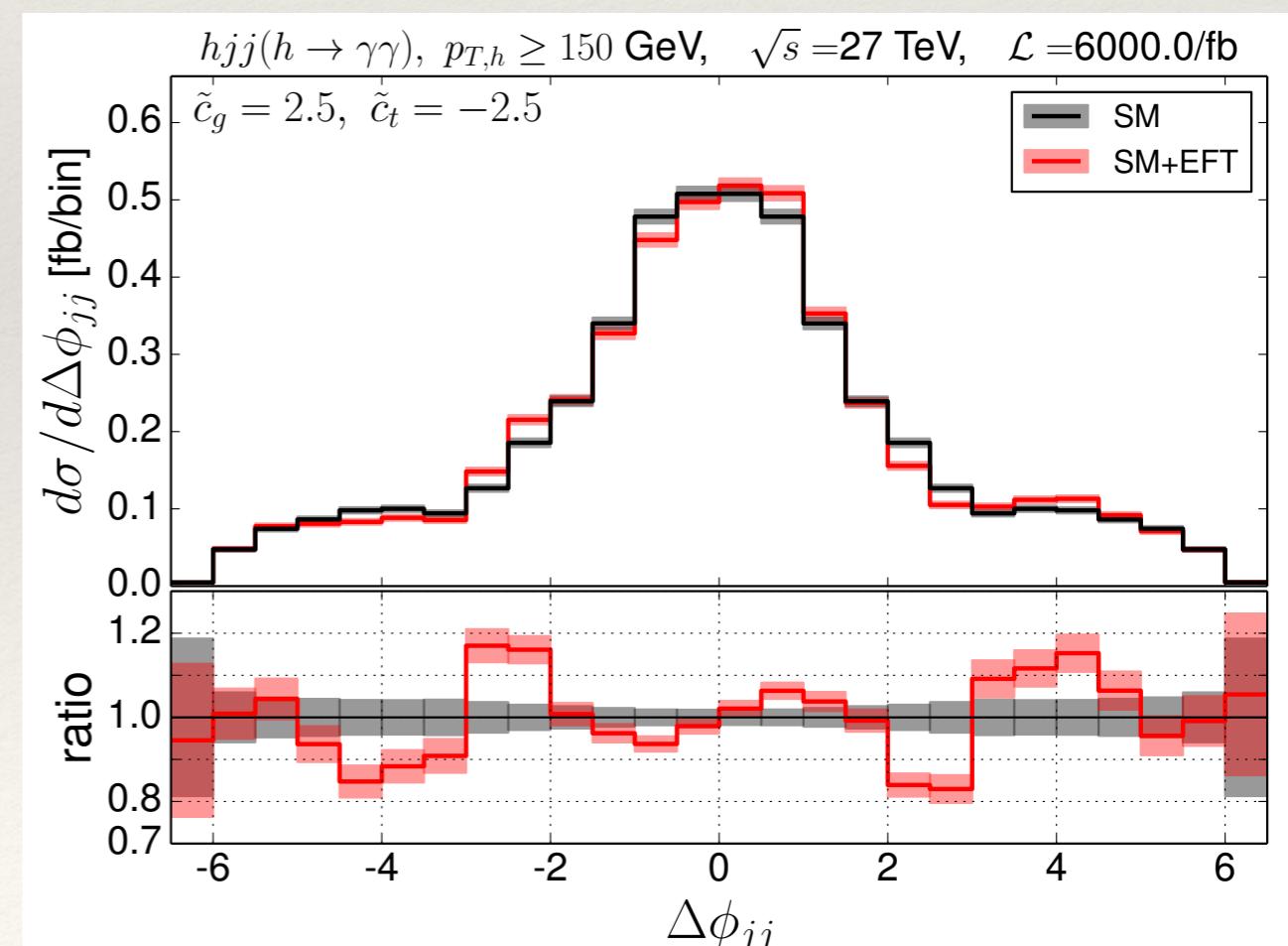
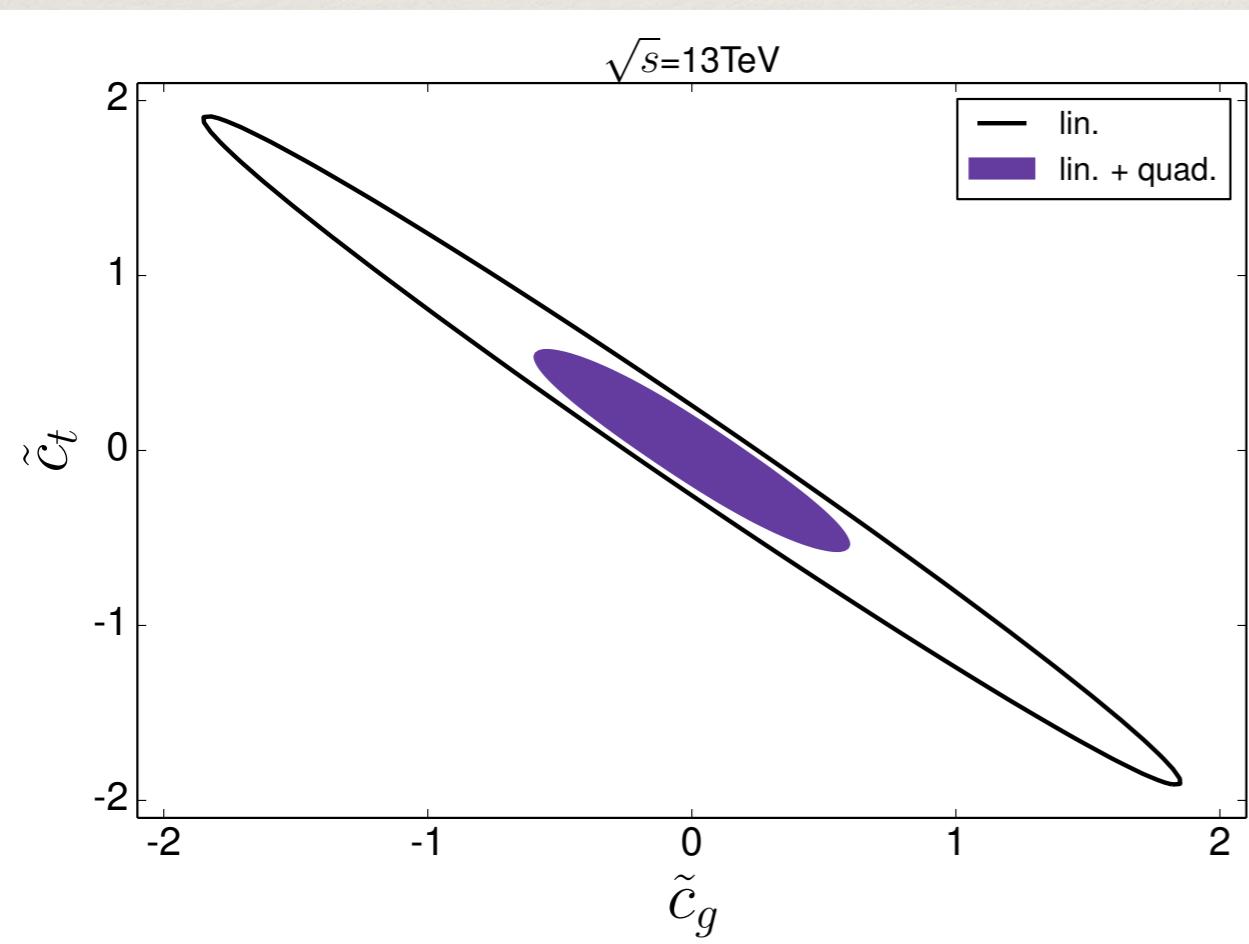
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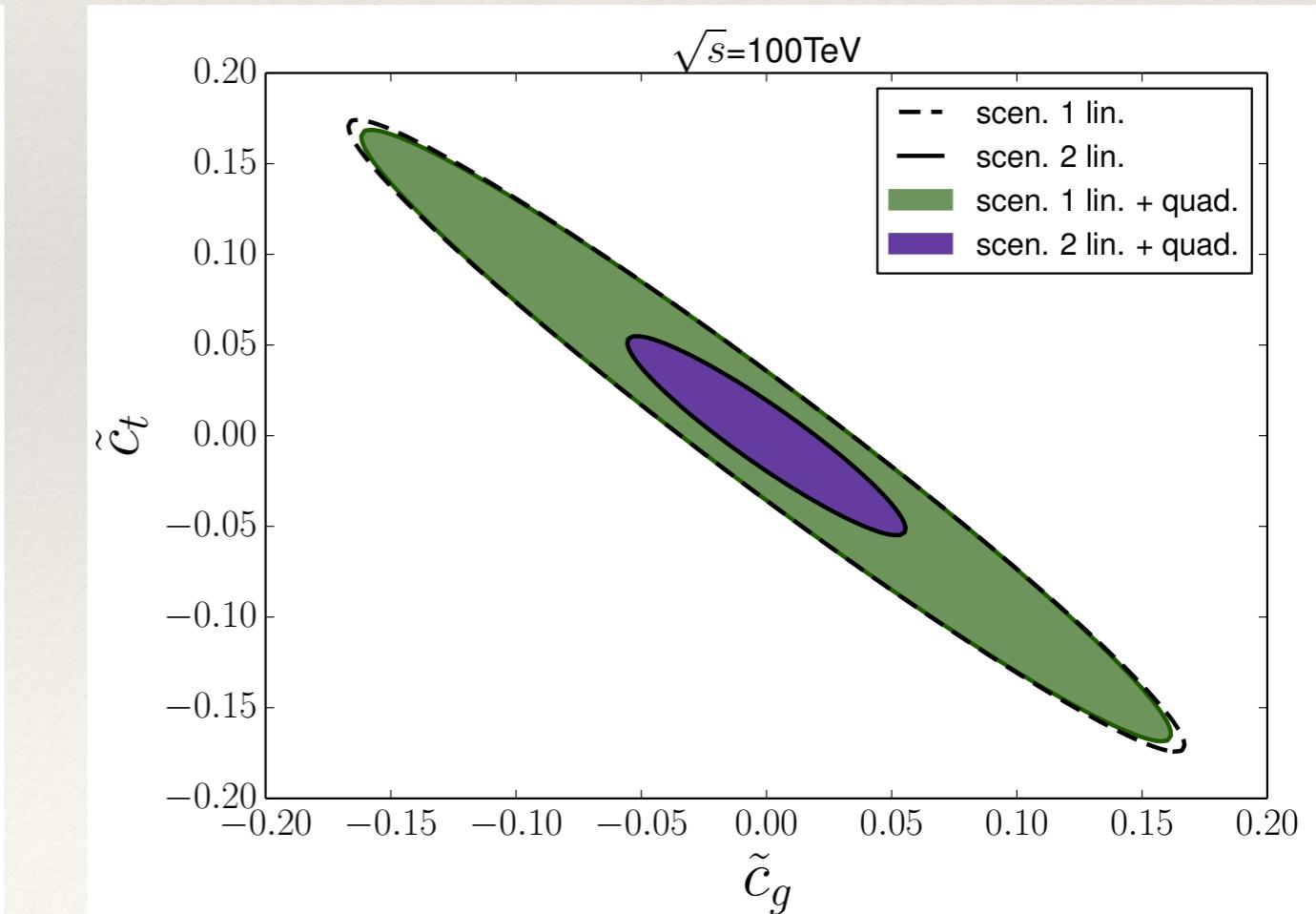
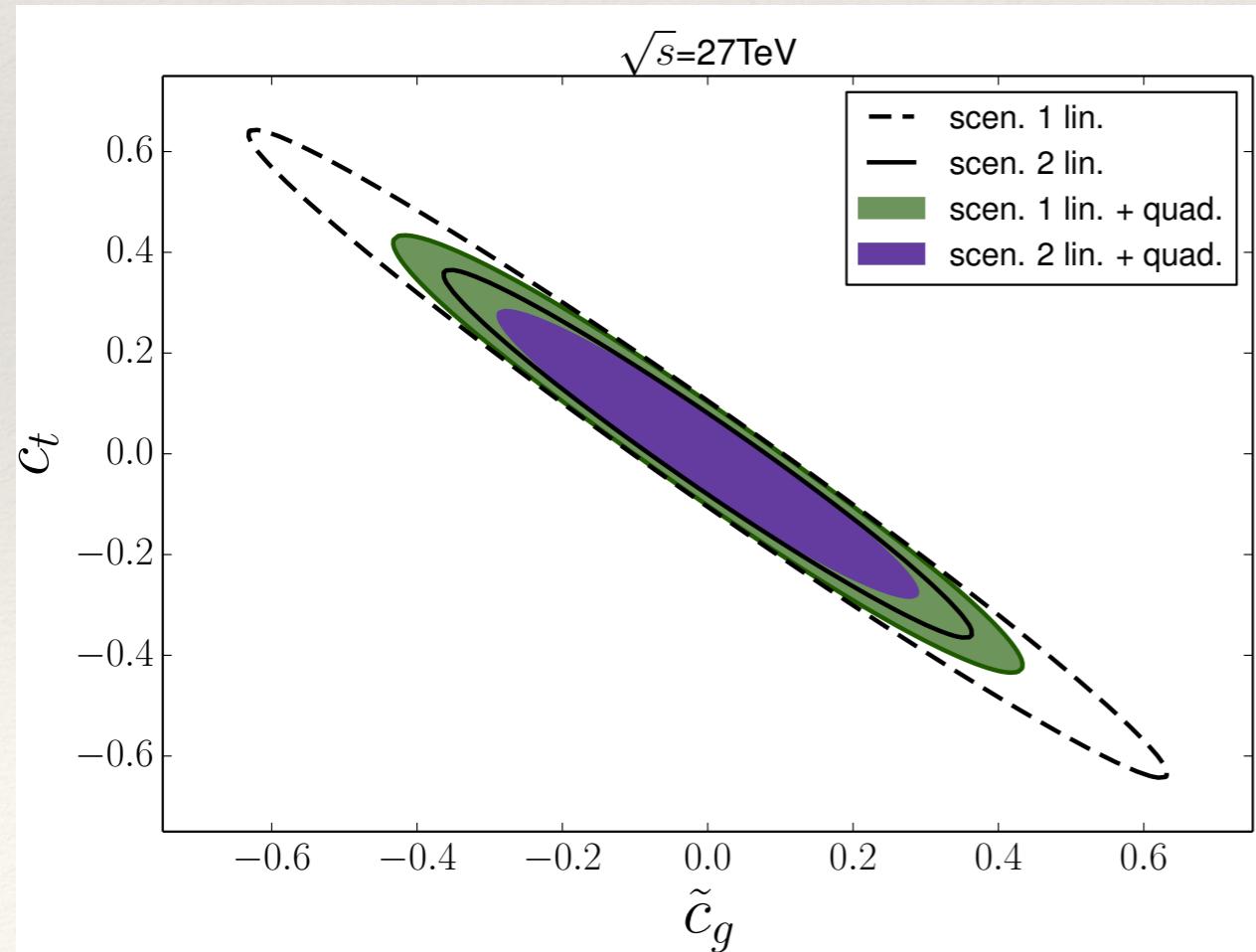
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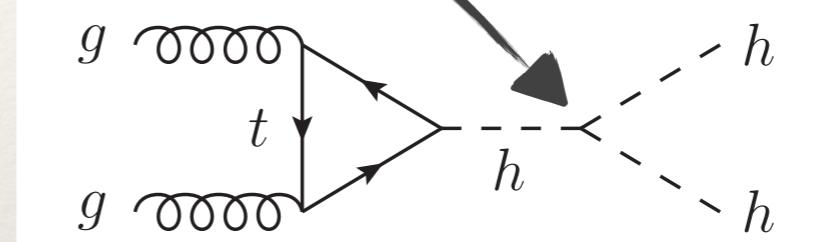
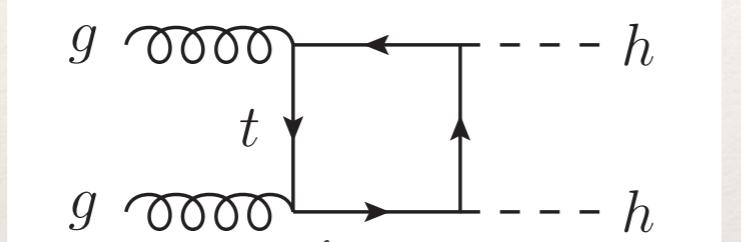


- dimension 6 deformations of the Higgs potential

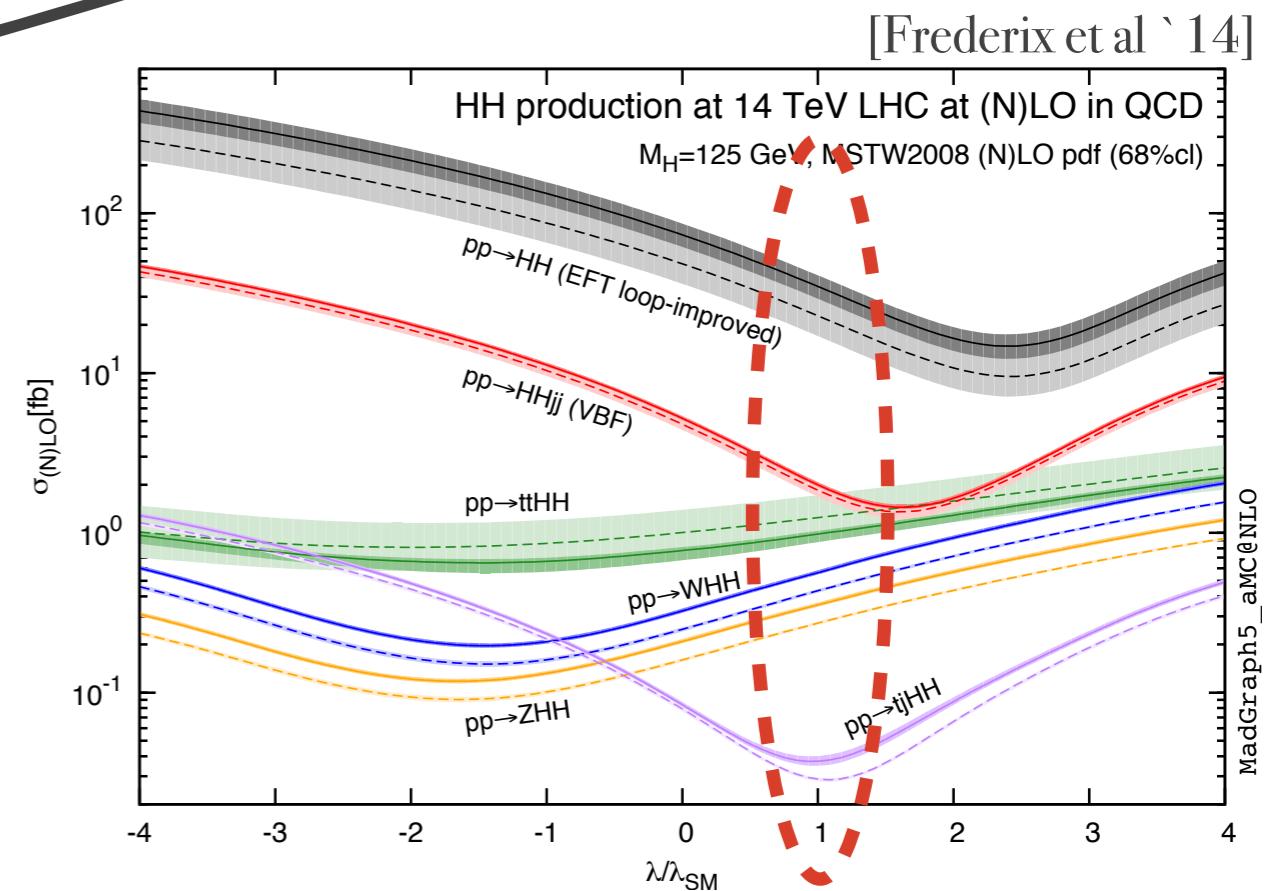
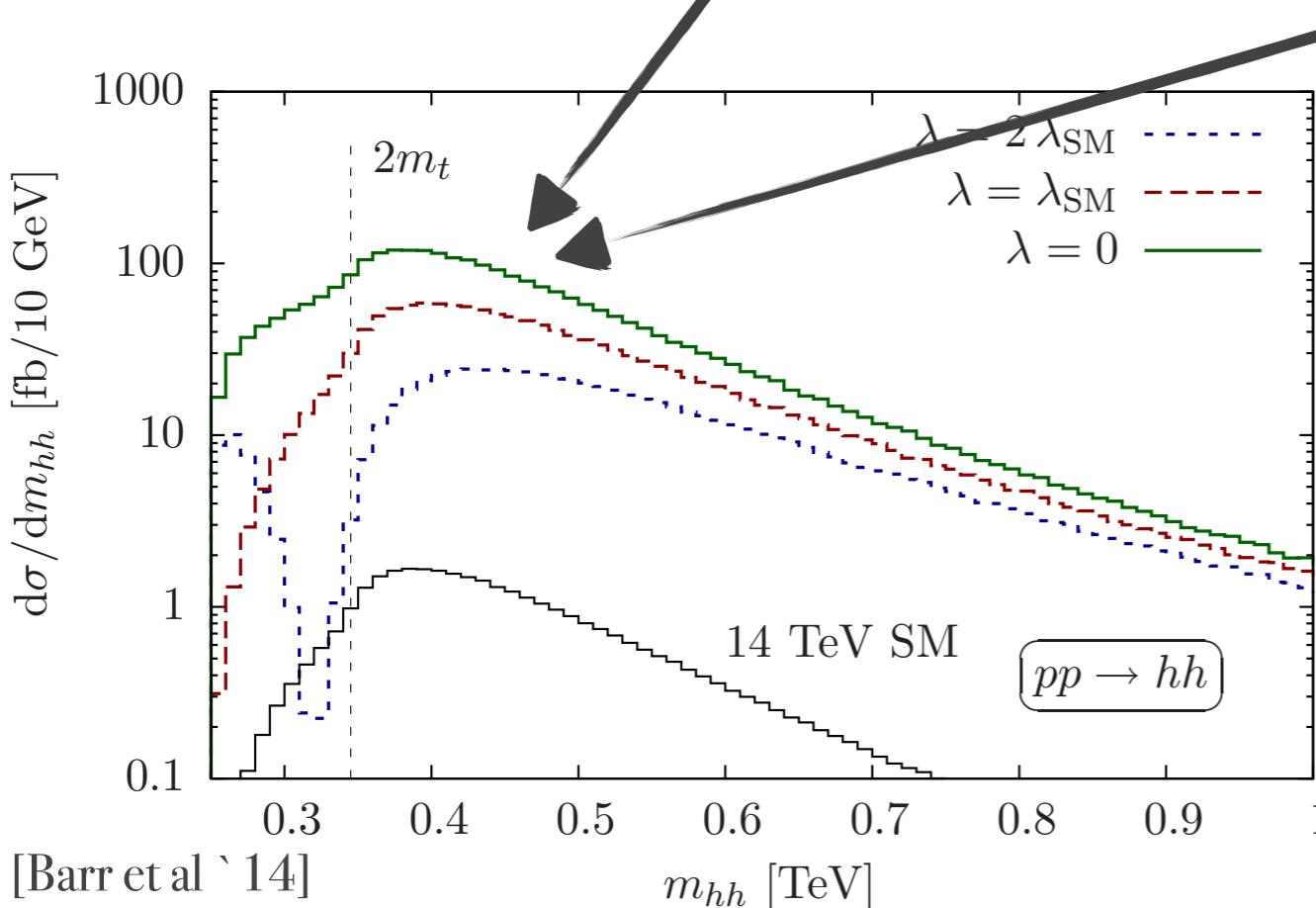
$$V(H^\dagger H)_6 \supset c_6/\Lambda^2 (H^\dagger H)^3$$

e.g.  
[Giudice, Grojean, Pomarol, Rattazzi '07]

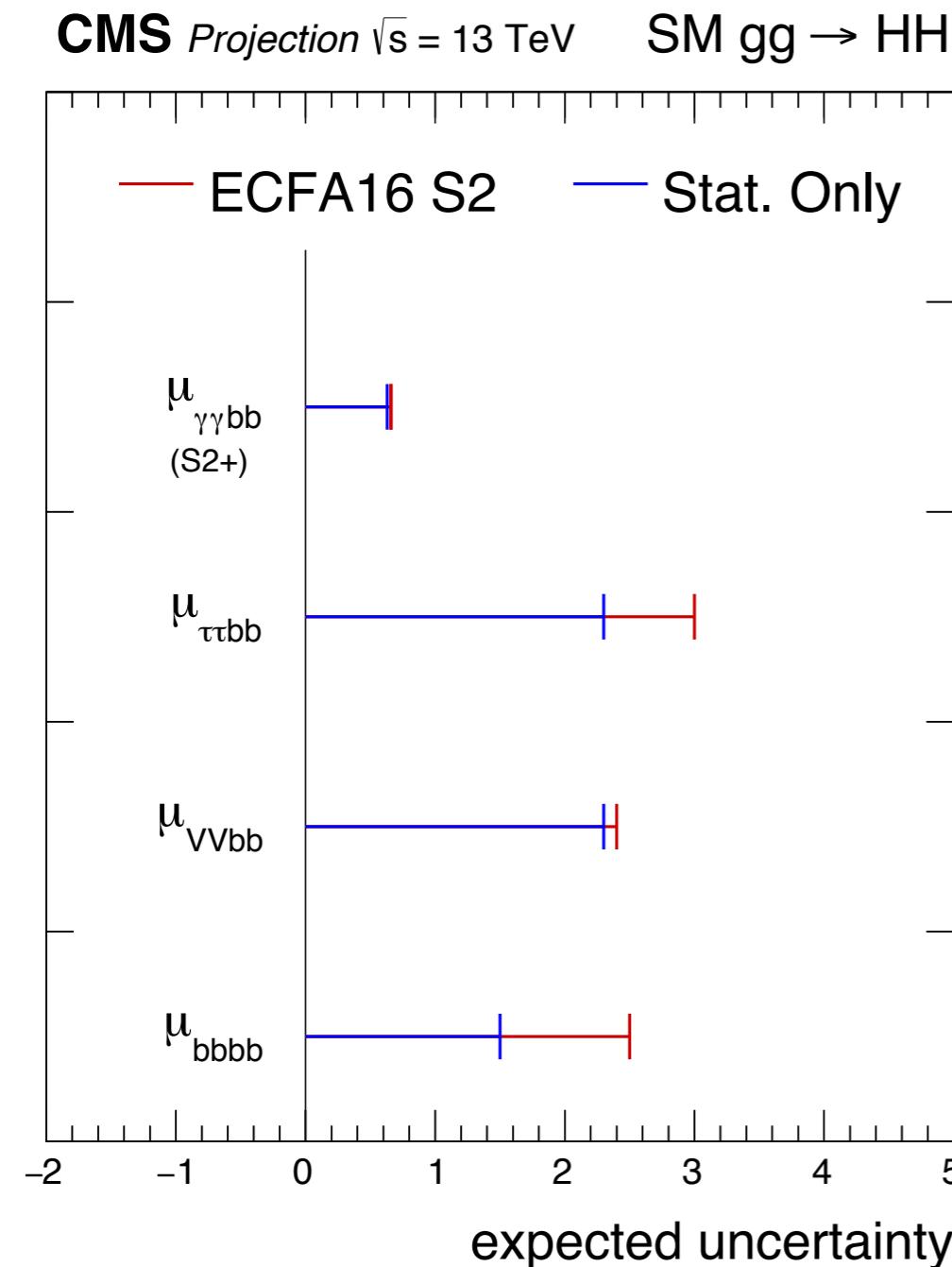
modify Higgs self-interactions. Large top-threshold interference.



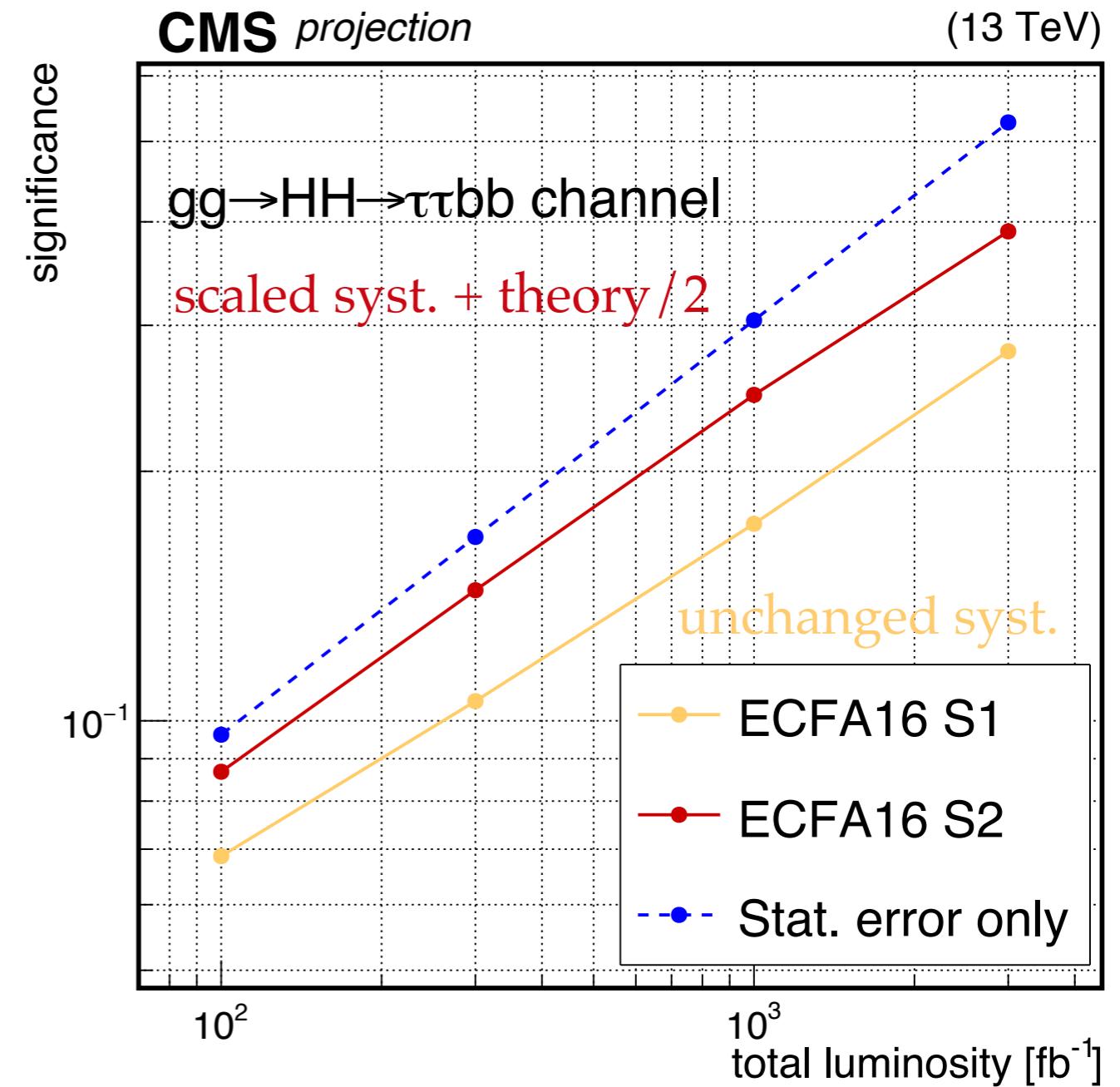
e.g.  
[Glover, van der Bij '88]



# HH pheno

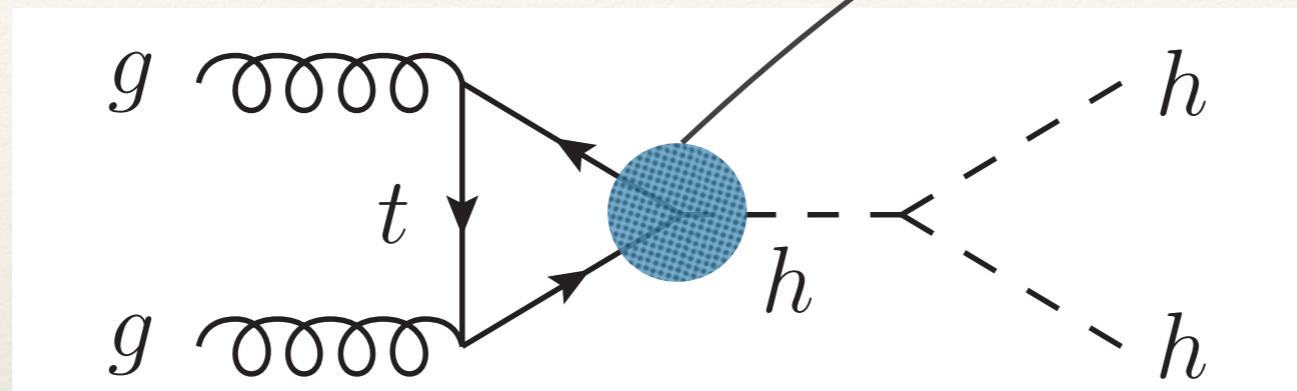


# LHC blind spots: Higgs potential



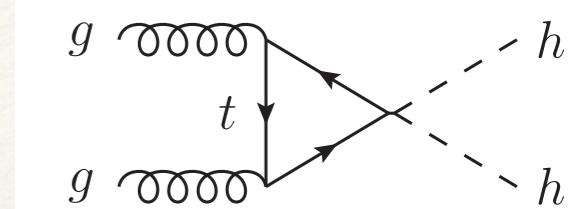
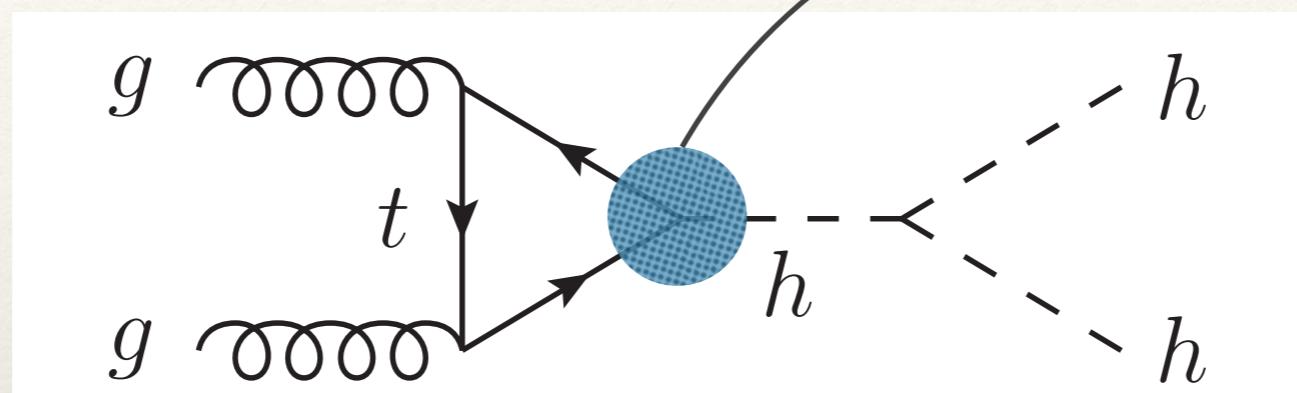
- however...

correlated with on-shell Higgs phenomenology



- however...

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broken by  $\sim \bar{t}th^2/\Lambda$ ....

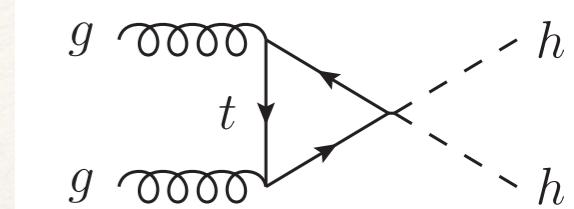
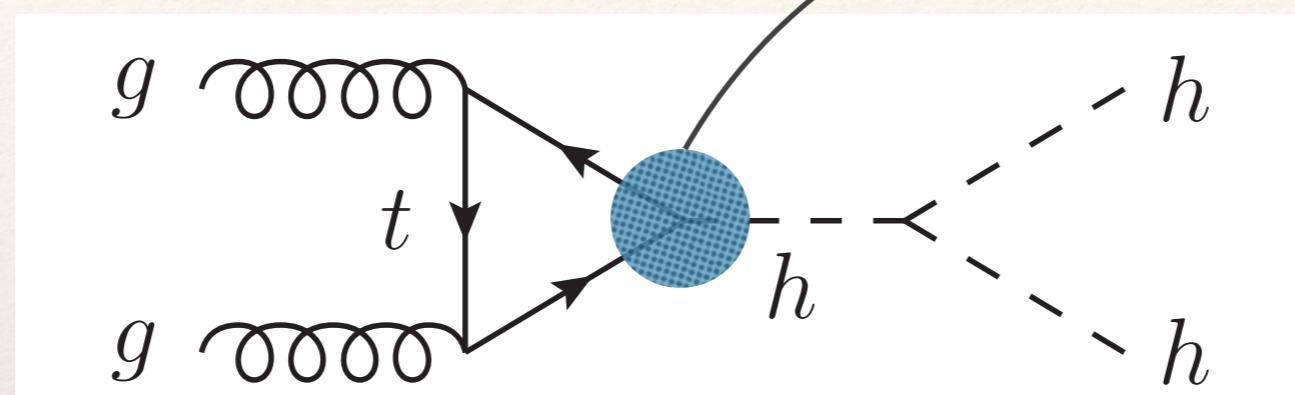


[Gröber, Mühlleitner '10]

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but can doubt physical relevance of such limits ( $\rightarrow$  matching)

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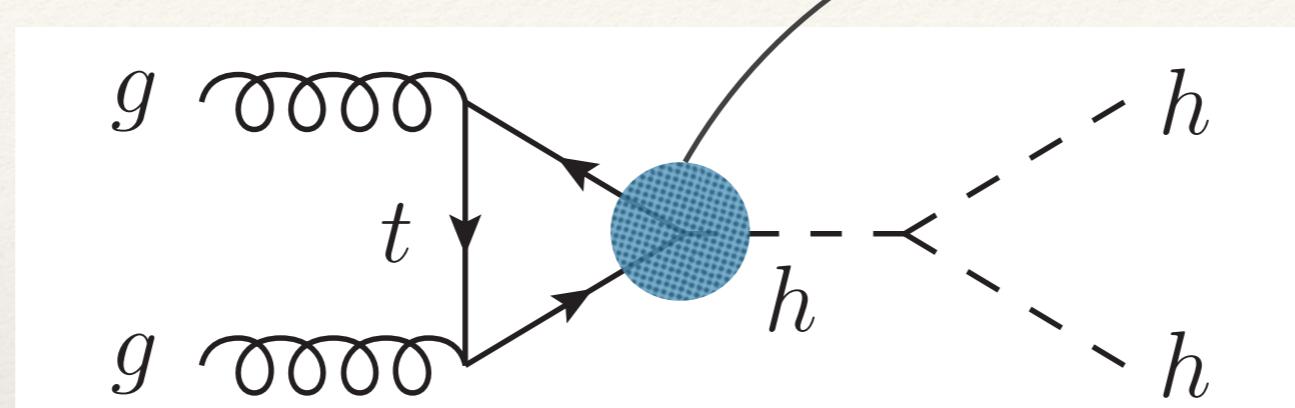


[Gröber, Mühlleitner '10]

- easy to arrange EFT coefficients in a way to get spectacular rates, but can doubt physical relevance of such limits ( $\rightarrow$  matching)
- use concrete Higgs sector extensions
  - extrapolate 125 GeV signal strengths
  - extrapolate exotic Higgs searches
  - additional constraints (*electron EDMs, flavor, perturbativity, ...*)

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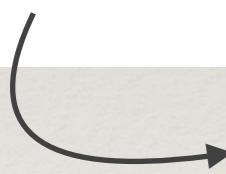
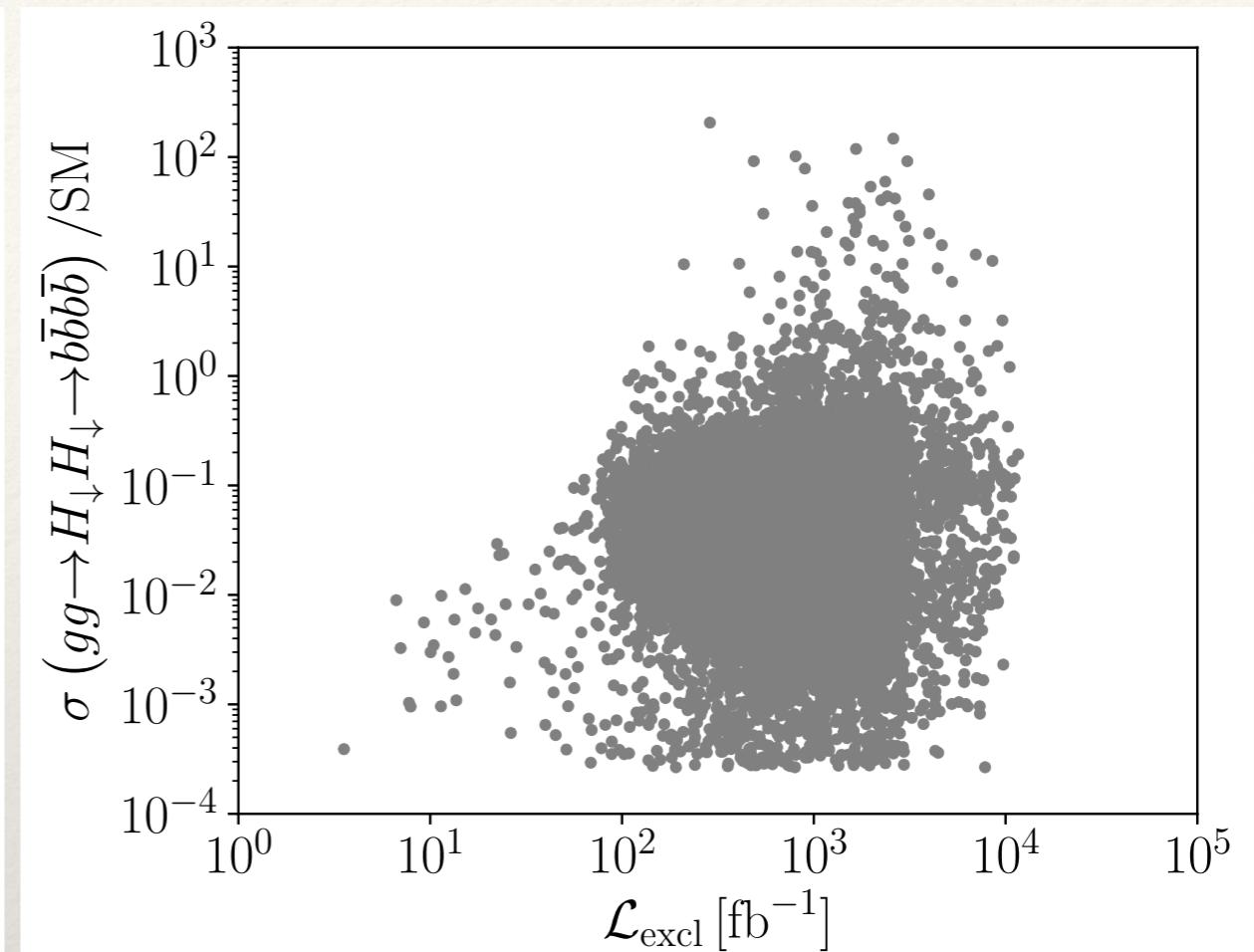
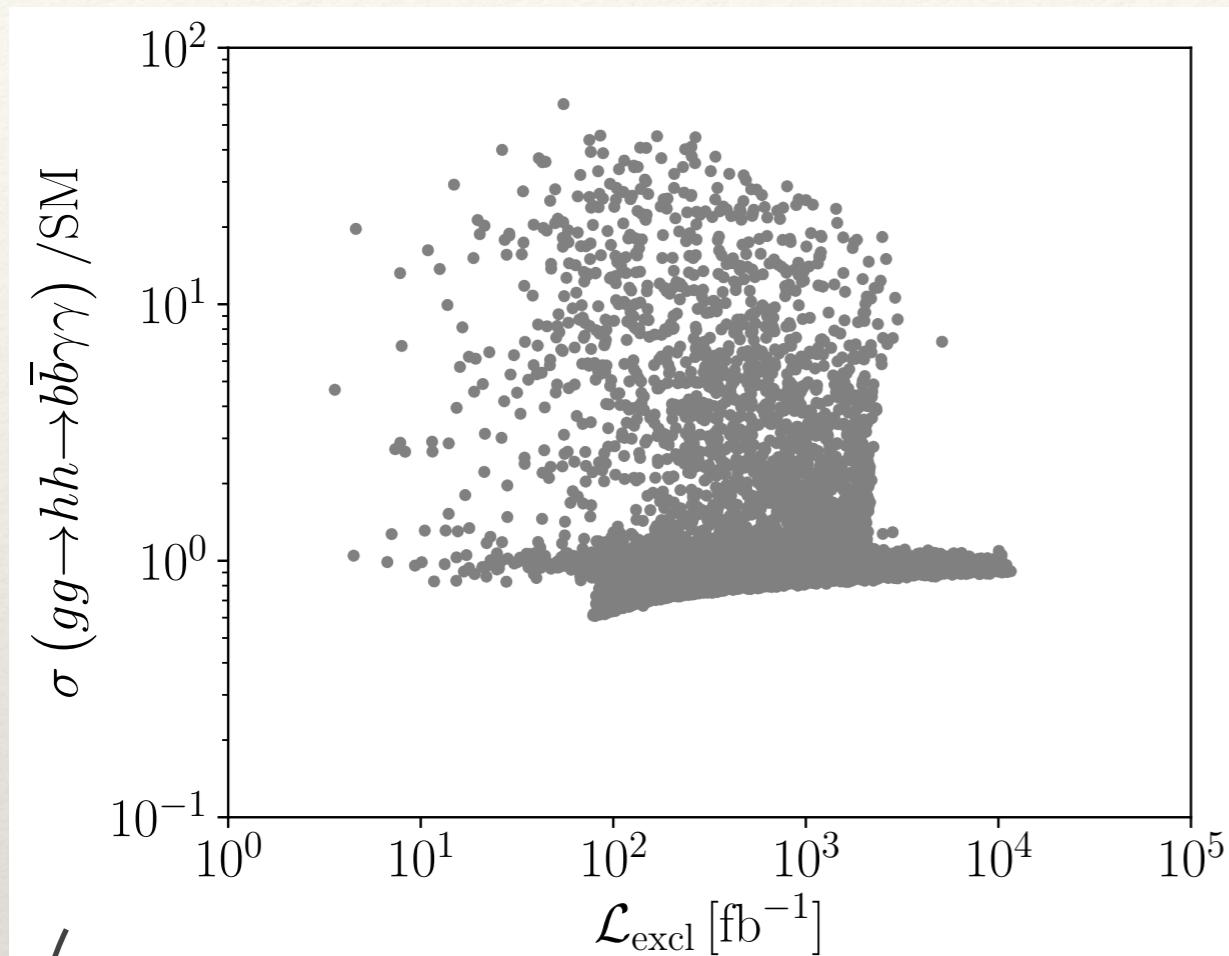
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in how far are di-Higgs final states still relevant at 3 / ab?

[Basler, Dawson, CE, Mühlleitner '18]



SM-like measurements can show a plethora of resonant anomalies  
diHiggs final states important for BSM discovery

...diHiggs final states quickly lose relevance when approaching EFT limit

## above Higgs pair threshold

- (multi) resonant diHiggs production (hh, hH,...)

[Basler, Dawson, CE, Mühlleitner '18]

Higgs interactions dominant

opportunity for diHiggs

exotics with large couplings to tops

top interactions dominant

## above top pair threshold

- tt final states preferred
- analysis highly model-dependent due to dedicated S-B interference

## below top pair threshold

- compressed spectra
- single Higgs competitive except b-final states (*trigger etc...*)

- ▶ *Technical advances have been extremely rapid*
  - ▶ matrix elements
  - ▶ jets
  - ▶ machine learning
- ▶ *Opportunity to link the Higgs sector to new physics*
  - ▶ cure SM shortcomings (CP violation...)
  - ▶ multi-Higgs production as a chance for BSM
  - ▶ LHC probably not be enough to achieve this in full glory