Flavour physics
as a test of the standard model
and a probe of new physics

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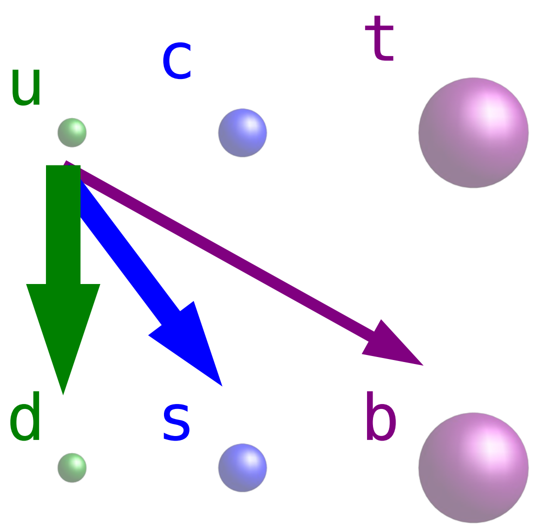
Outline

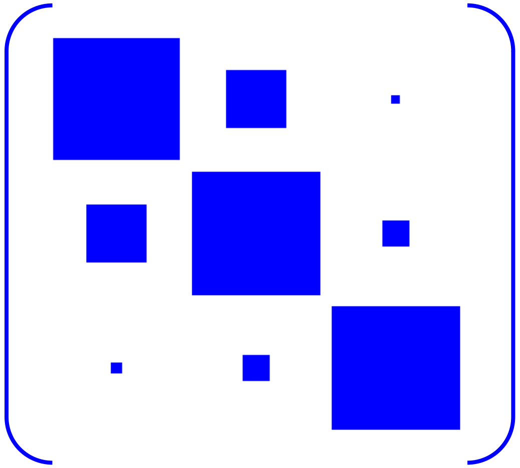
- very briefly:
  - introduction and motivations
  - the tool: the Unitarity Triangle fit
- Standard Model fit
  - Standard model constraints
  - checking for tensions
  - Standard Model predictions
- Beyond the Standard Model:
  - model-independent analysis
  - New-physics-specific constraints
  - New-physics scale analysis

Flavour mixing and CP violation in the Standard Model

- The CP symmetry is violated in any field theory having in the Lagrangian at least one phase that cannot be re-absorbed
- The **mass eigenstates** are not eigenstates of the weak interaction. This feature of the Standard Model Hamiltonian produces the (unitary) **mixing matrix** V_{CKM}.

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \approx \begin{pmatrix} 1-\frac{\lambda^2}{2} & \lambda & A\lambda^3(\overline{\rho}-i\overline{\eta}) \\ -\lambda & 1-\frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1-\overline{\rho}-i\overline{\eta}) & -A\lambda^2 & 1 \end{pmatrix}$$





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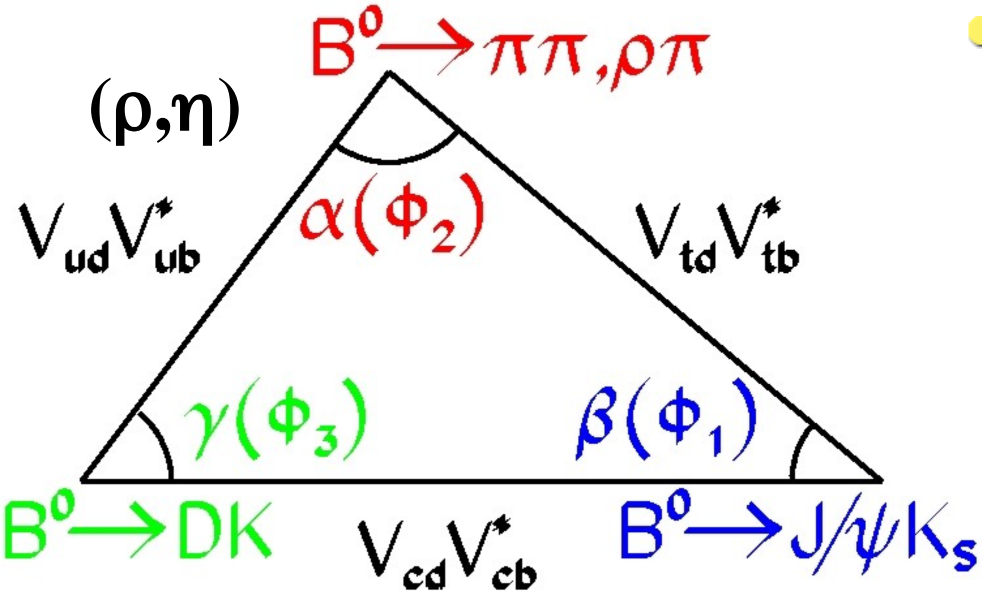


- With three families of quarks, there is one phase that allows CP violation in the SM. All the flavour mixing processes are related (through the unitarity of the V_{CKM}) to this phase.

Unitarity Triangle

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

All the angles are related to the CP asymmetries of specific B decays

CKM matrix and Unitarity Triangle

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

many observables functions of $\bar{\rho}$ and $\bar{\eta}$: overconstraining

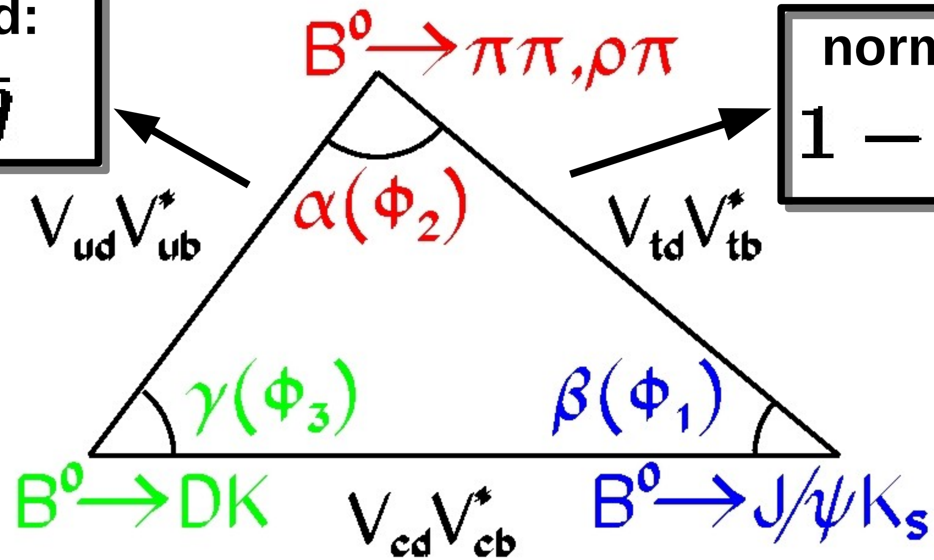
$$\alpha = \pi - \beta - \gamma$$

normalized:

$$\bar{\rho} + i\bar{\eta}$$

normalized:

$$1 - \bar{\rho} - i\bar{\eta}$$



$$\gamma = atan\left(\frac{\bar{\eta}}{\bar{\rho}}\right)$$

$$\beta = atan\left(\frac{\bar{\eta}}{(1-\bar{\rho})}\right)$$

www.utfit.org

RELOADED

M.Bona, M. Ciuchini, D. Derkach, F. Ferrari, E. Franco,
V. Lubicz, G. Martinelli, M. Pierini, L. Silvestrini,
S. Simula, C. Tarantino, V. Vagnoni, M. Valli, and L.Vittorio

Method and inputs:

$$f(\bar\rho,\bar\eta,X|c_1,...,c_m) \sim \prod_{j=1,m} f_j(\mathcal{C}|\bar\rho,\bar\eta,X) * \prod_{i=1,N} f_i(x_i) f_0(\bar\rho,\bar\eta)$$

Bayes Theorem

$$X \equiv x_1,...,x_n = m_t, B_K, F_B, ...$$

$$\mathcal{C} \equiv c_1,...,c_m = \epsilon, \Delta m_d/\Delta m_s, A_{CP}(J/\psi K_S), ...$$

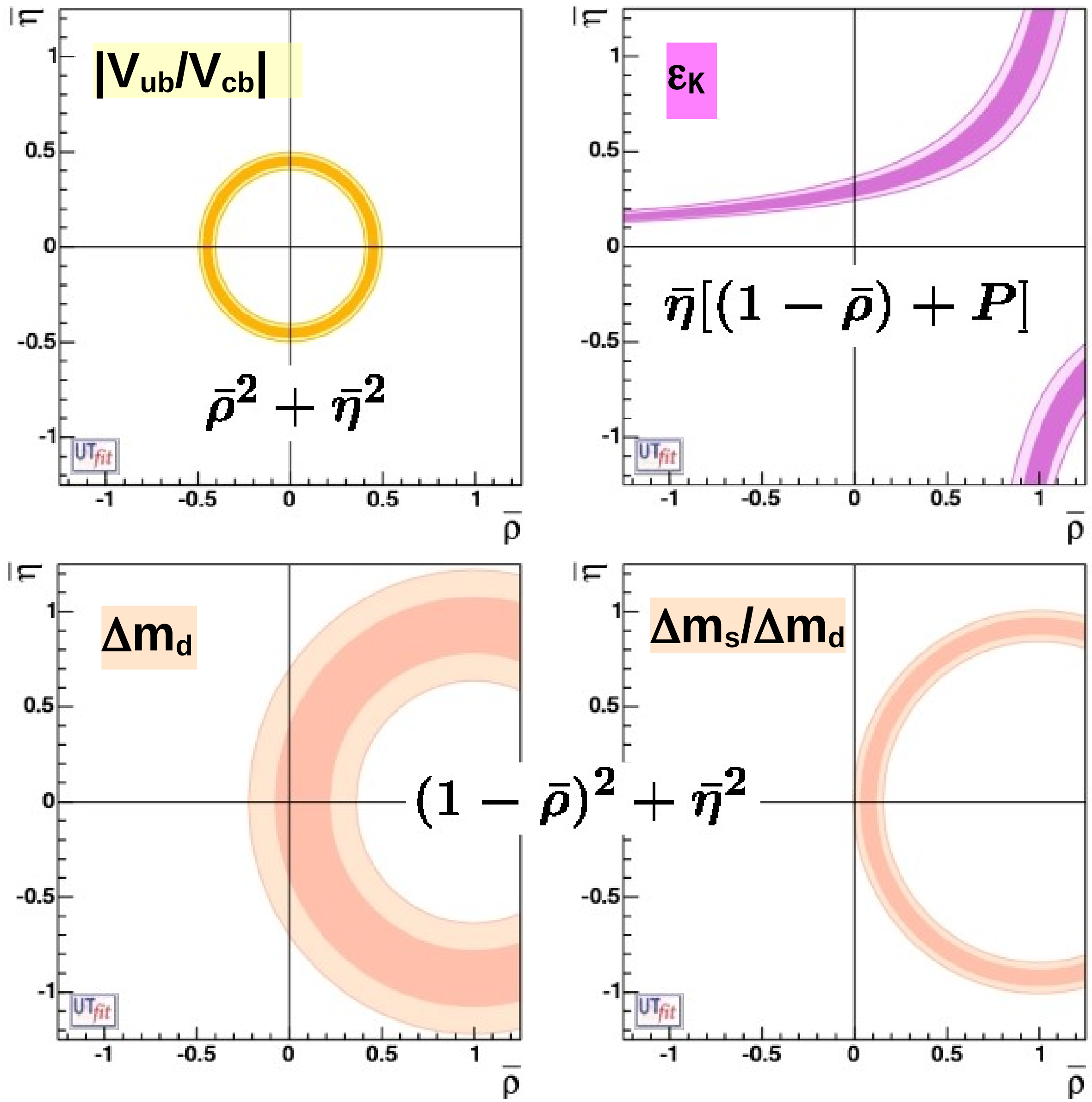
$(b\to u)/(b\to c)$	$\bar\rho^2+\bar\eta^2$	$\bar\Lambda, \lambda_1, F(1), \ldots$
ϵ_K	$\bar\eta[(1-\bar\rho)+P]$	B_K
Δm_d	$(1-\bar\rho)^2+\bar\eta^2$	$f_B^2 B_B$
$\Delta m_d/\Delta m_s$	$(1-\bar\rho)^2+\bar\eta^2$	ξ
$A_{CP}(J/\psi K_S)$	$\sin 2\beta$	

, $\mathbf{m_t}$

Standard Model +
OPE/HQET/
Lattice QCD
to go
from quarks
to hadrons

M. Bona *et al*. (UTfit Collaboration)
JHEP 0507:028,2005 hep-ph/0501199
M. Bona *et al*. (UTfit Collaboration)
JHEP 0603:080,2006 hep-ph/0509219

The LEP-style analysis in the $\bar{\rho}$-$\bar{\eta}$ plane:



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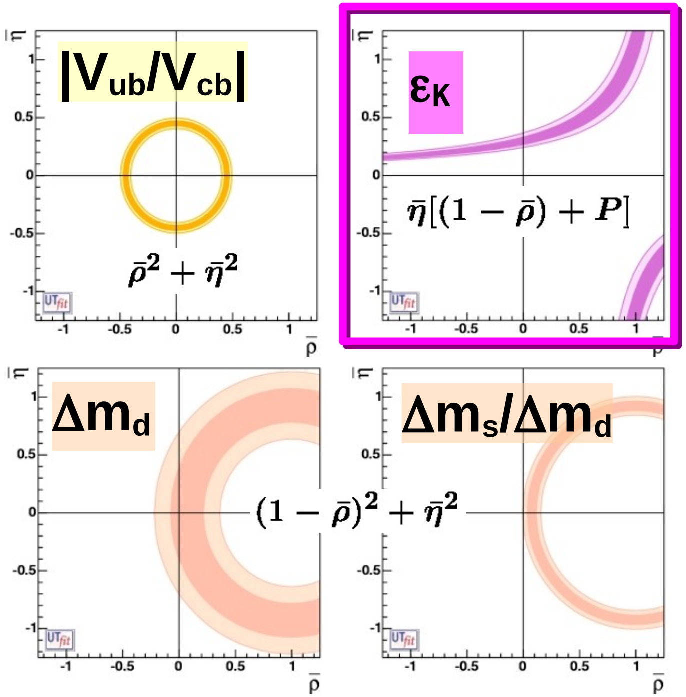
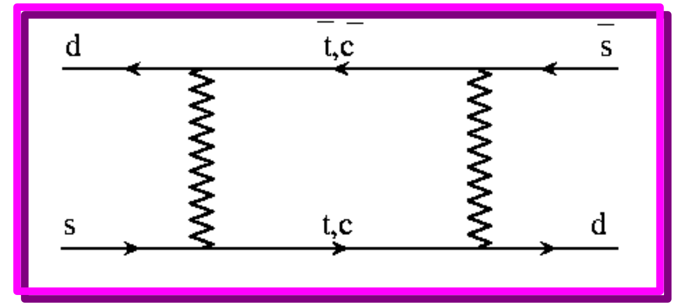


ε_K from K-$\bar{K}$ mixing



$\varepsilon_K = (2.228 \pm 0.011) \cdot 10^{-3}$ PDG

$$B_K = \frac{\langle K | J_\mu J^\mu | \overline{K} \rangle}{\langle K | J_\mu | 0 \rangle \langle 0 | J^\mu | \overline{K} \rangle}$$

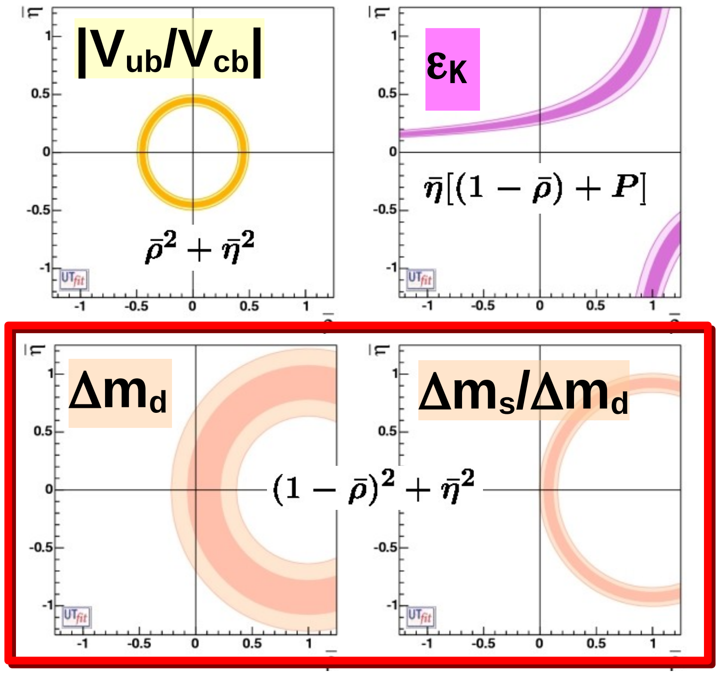
$B_K = 0.756 \pm 0.016$ FLAG 2019

from lattice QCD

$$|\varepsilon_K| = C_\varepsilon B_K A^2 \lambda^6 \bar{\eta} \{ -\eta_1 S_0(x_c)(1-\lambda^2/2) + \eta_3 S_0(x_c, x_t) + \eta_2 S_0(x_t) A^2 \lambda^4 (1-\bar{\rho}) \}$$

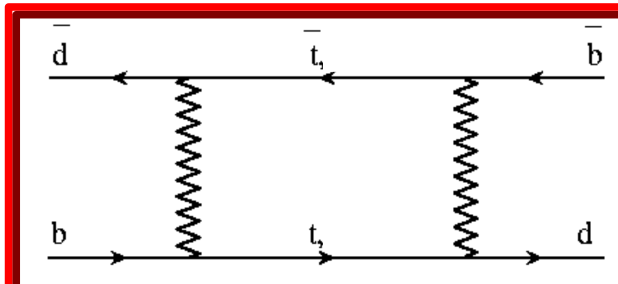
S_0 = Inami-Lim functions for c-c, c-t, e t-t contributions
(from perturbative calculations)

The LEP-style analysis in the $\bar{\rho}$-$\bar{\eta}$ plane:



Δm_q from B_q-$\bar{B}_q$ mixing

q=d,s



$\Delta m_d = 0.5065 \pm 0.0019\ ps^{-1}$ HFLAV

$\Delta m_s = 17.765 \pm 0.006\ ps^{-1}$ HFLAV

$$\Delta m_d = \frac{G_F^2}{6\pi^2} m_W^2 \eta_b S(x_t) m_{B_d} f_{B_d}^2 \hat{B}_{B_d} |V_{tb}|^2 |V_{td}|^2 =$$

$$= \frac{G_F^2}{6\pi^2} m_W^2 \eta_b S(x_t) m_{B_d} f_{B_d}^2 \hat{B}_{B_d} |V_{cb}|^2 \lambda^2 ((1-\bar{\rho})^2 + \bar{\eta}^2)$$

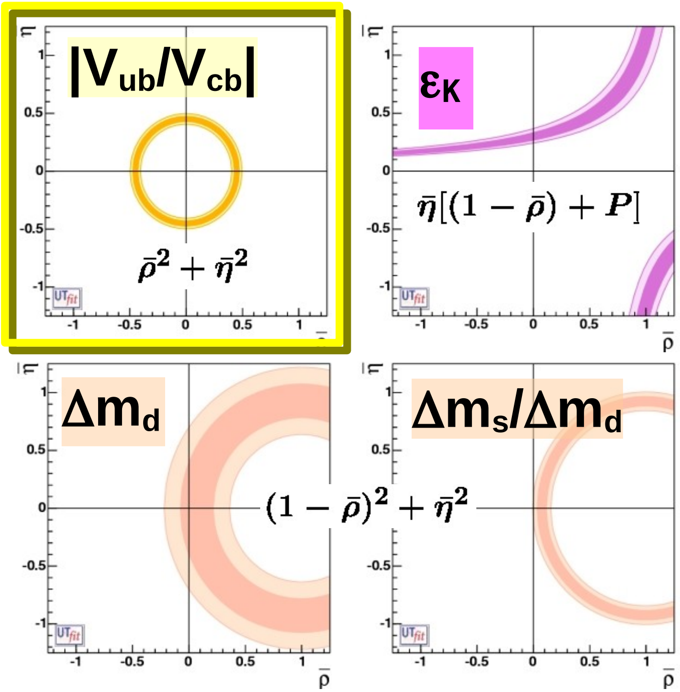
$$\Delta m_d \approx [(1-\rho)^2 + \eta^2] \frac{f_{B_s}^2 B_{B_s}}{\xi^2}$$

$$\Delta m_s \approx f_{B_s}^2 B_{B_s}$$

S = Inami-Lim function
for the t-t contribution
(from perturbative calculations)

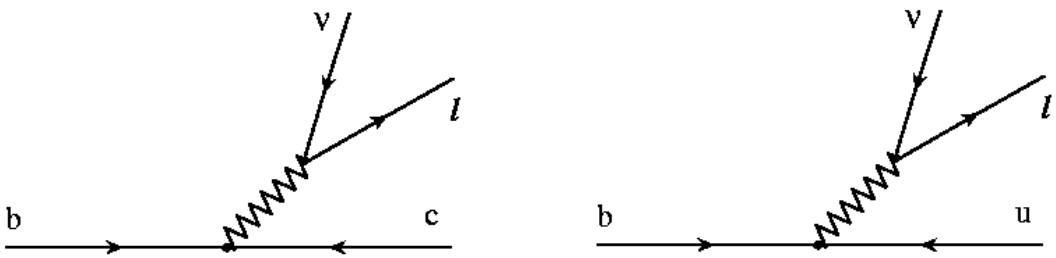
B_{B_q} and f_{B_q} from lattice QCD

The LEP-style analysis in the $\bar{\rho}$-$\bar{\eta}$ plane:



$$\left|\frac{V_{ub}}{V_{cb}}\right| = \frac{\lambda}{1-\frac{\lambda^2}{2}}\sqrt{\bar{\rho}^2+\bar{\eta}^2}$$

$|V_{ub}/V_{cb}|$



tree diagrams

b → c and b → u transition

- negligible new physics contributions
- inclusive and exclusive semileptonic B decay branching ratios

QCD corrections to be included

- inclusive measurements: OPE
- exclusive measurements: form factors from lattice QCD

V_{cb} and V_{ub}

from FLAG 2019 arXiv:1902.08191

$|V_{cb}|$ *(excl)* = (39.09 ± 0.68) 10^{-3}

$|V_{cb}|$ *(incl)* = (42.16 ± 0.50) 10^{-3}

from Bordone et al. arXiv:2107.00604

~2.8σ discrepancy

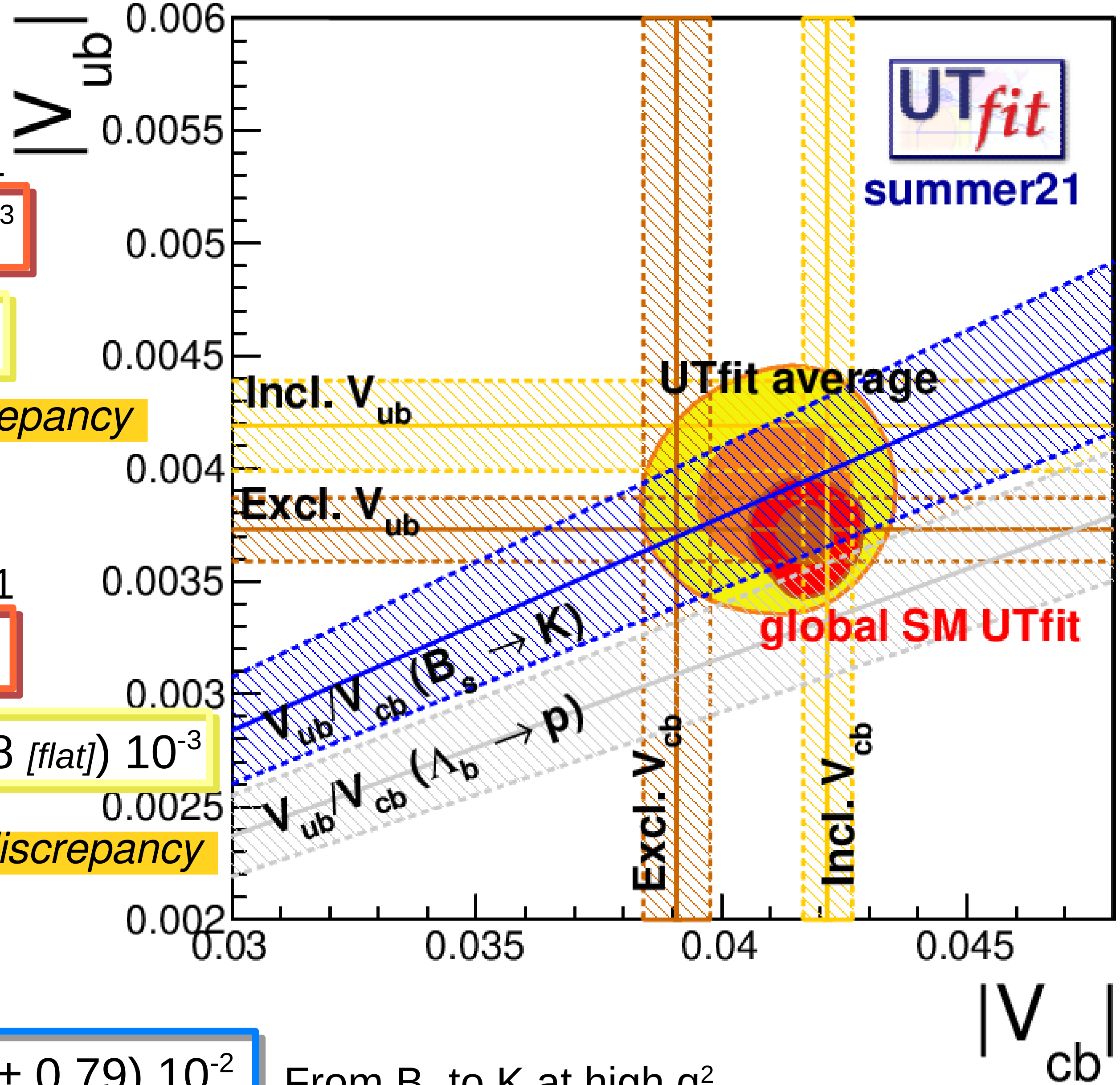
from FLAG 2019 arXiv:1902.08191

$|V_{ub}|$ *(excl)* = (3.73 ± 0.14) 10^{-3}

$|V_{ub}|$ *(incl)* = (4.19 ± 0.17 ± 0.18 *[flat]*) 10^{-3}

from GGOU HFLAV 2021 adding a flat uncertainty covering the spread of central values

~1.5σ discrepancy



$|V_{ub} / V_{cb}|$ *(LHCb)* = (9.46 ± 0.79) 10^{-2} — From B_s to K at high q^2

$|V_{ub} / V_{cb}|$ *(LHCb)* = (7.9 ± 0.6) 10^{-2} — From Λ_b, excluded following FLAG guidelines

V_{cb} and V_{ub}

A-la-D'Agostini two-dimensional average procedure:

$|V_{cb}| = (41.1 \pm 1.0)\ 10^{-3}$

uncertainty ~ 2.4%

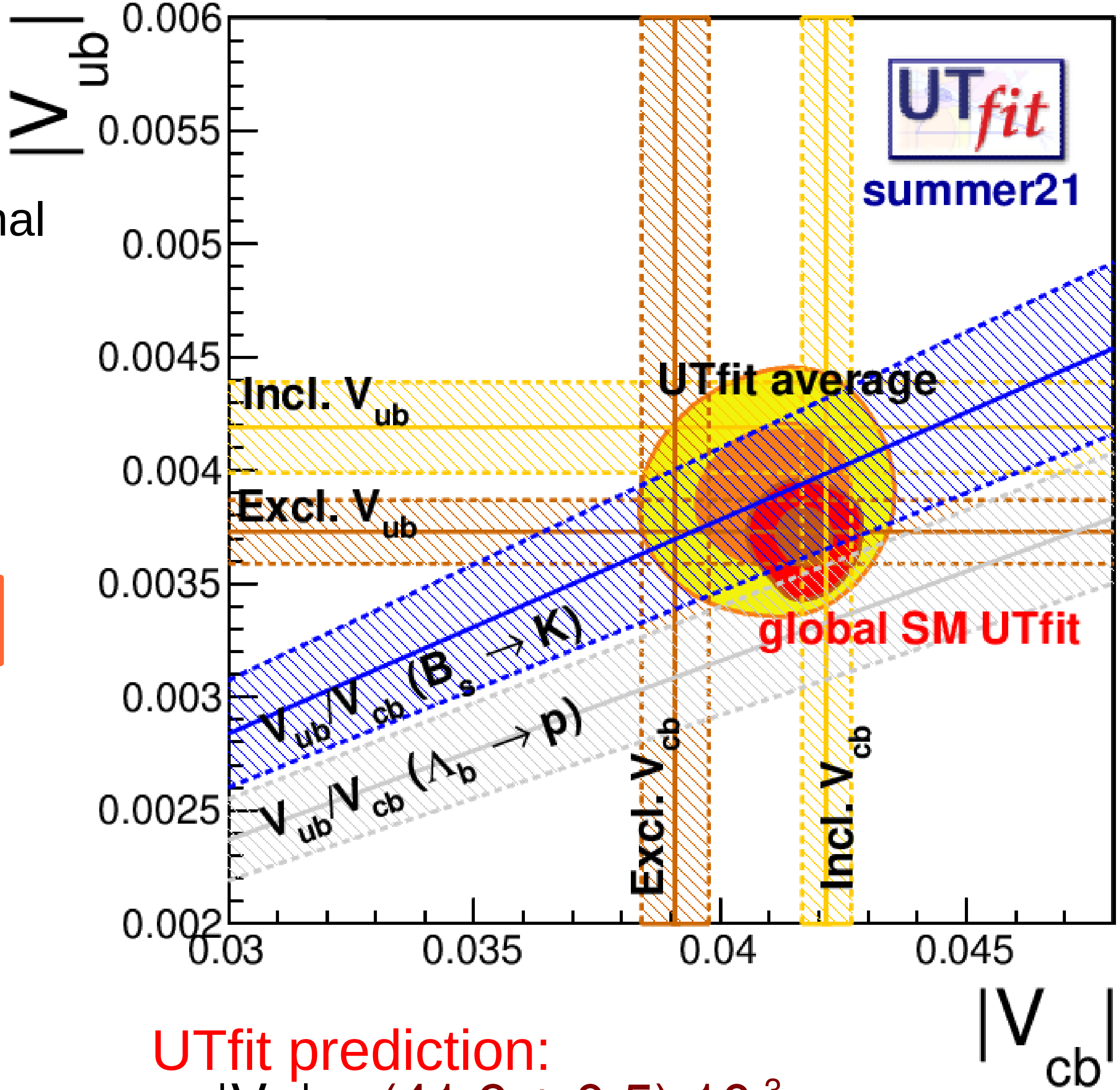
$|V_{ub}| = (3.89 \pm 0.21)\ 10^{-3}$

uncertainty ~ 5.4%

From global SM fit

$|V_{cb}| = (41.7 \pm 0.4)\ 10^{-3}$

$|V_{ub}| = (3.70 \pm 0.10)\ 10^{-3}$

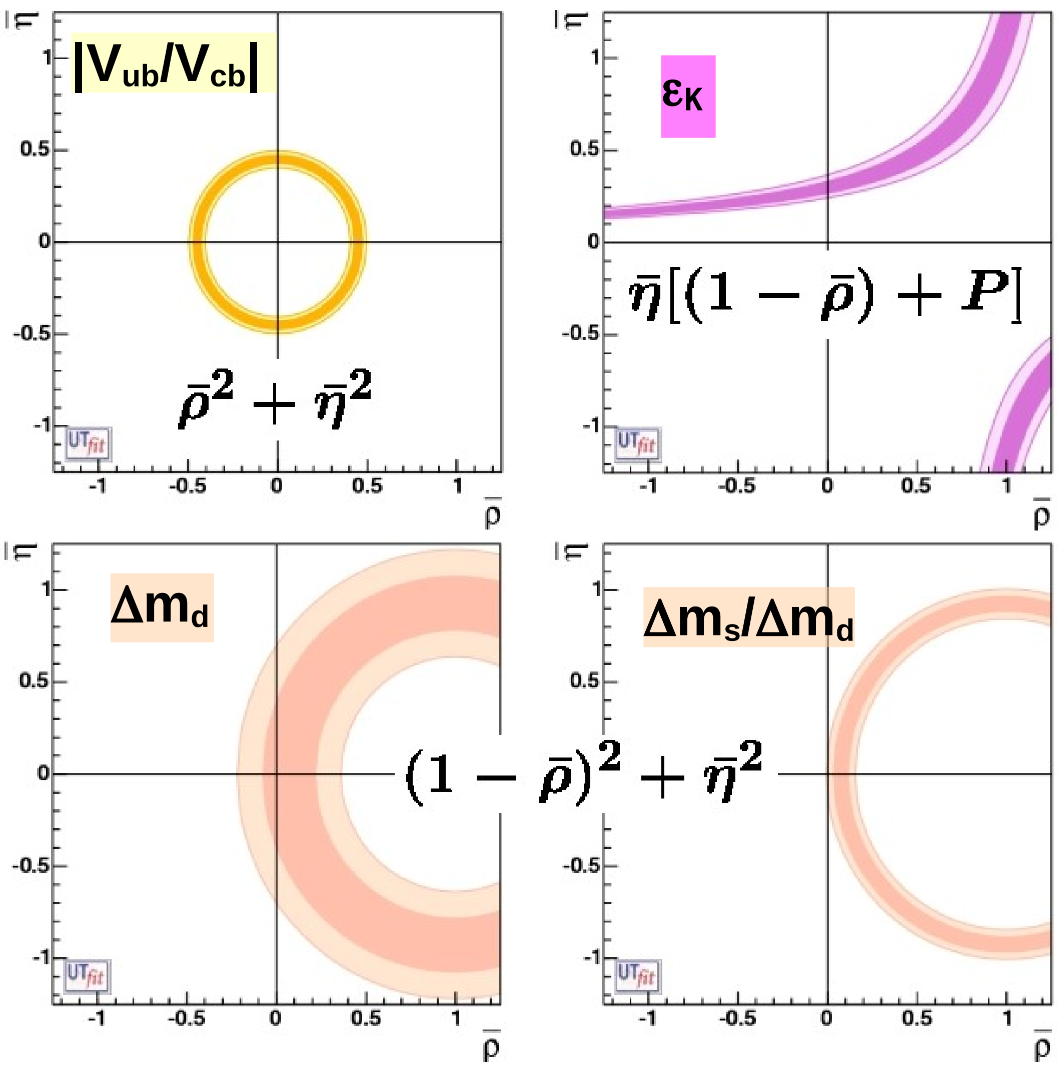


UTfit prediction:

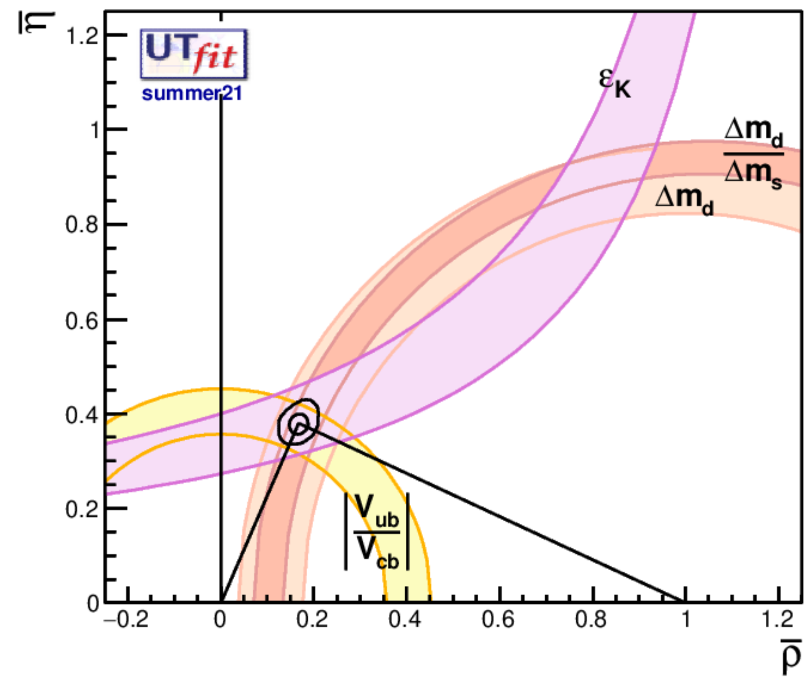
$|V_{cb}| = (41.9 \pm 0.5)\ 10^{-3}$

$|V_{ub}| = (3.68 \pm 0.10)\ 10^{-3}$

The LEP-style analysis in the $\bar{\rho}$-$\bar{\eta}$ plane:



levels @
95% Prob

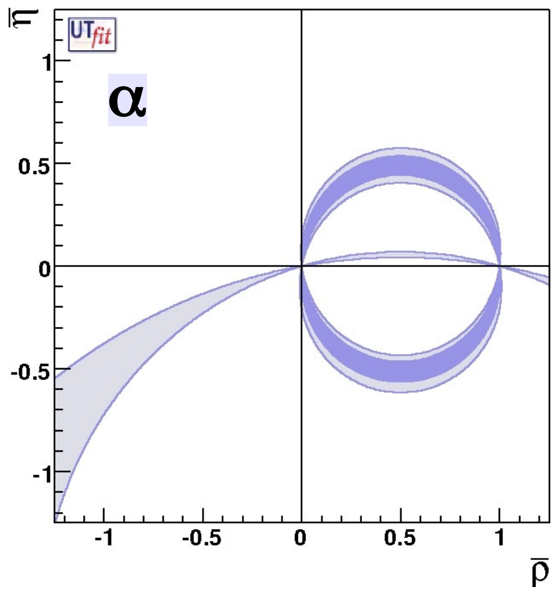


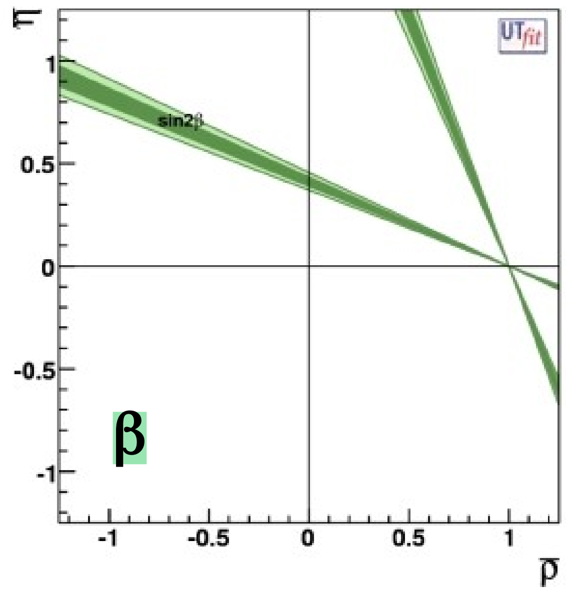
~10%

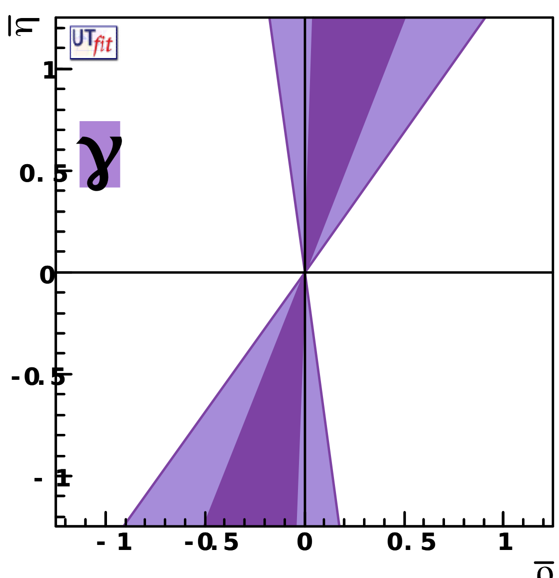
$\bar{\rho}$ = 0.169 ± 0.017
$\bar{\eta}$ = 0.383 ± 0.025

~7%

angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:

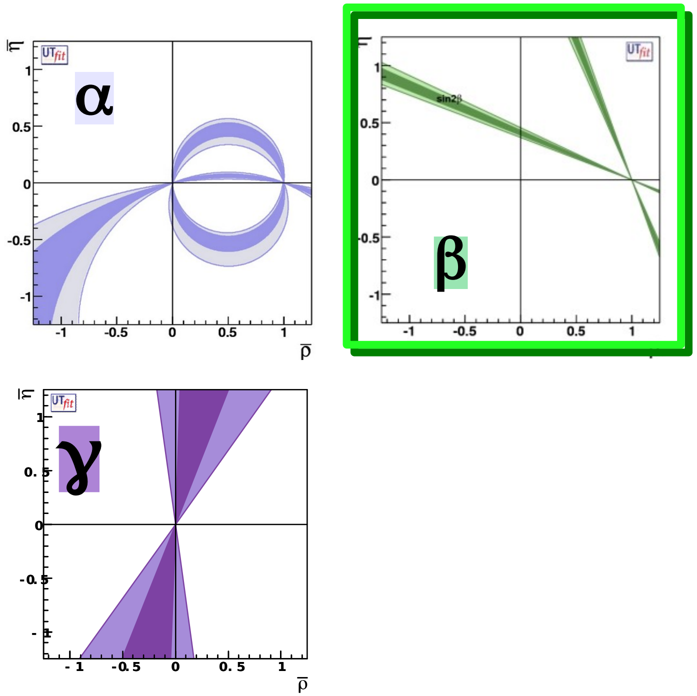




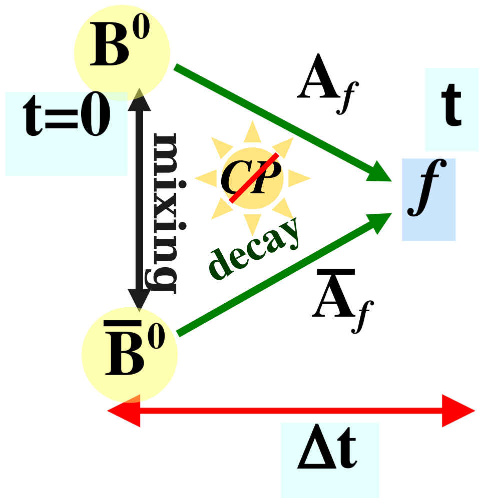


B factories
+ LHCb

angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:



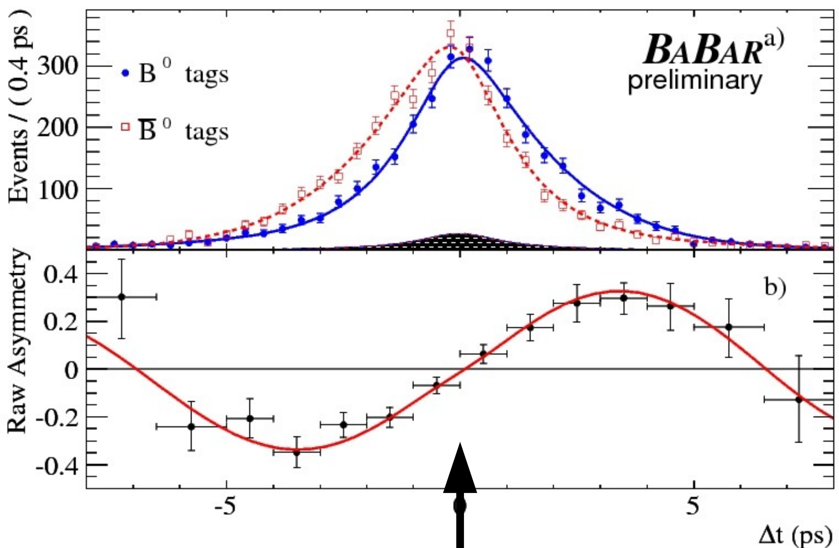
sin2β from time-dependent A_{CP} in $B \to J/\psi K$



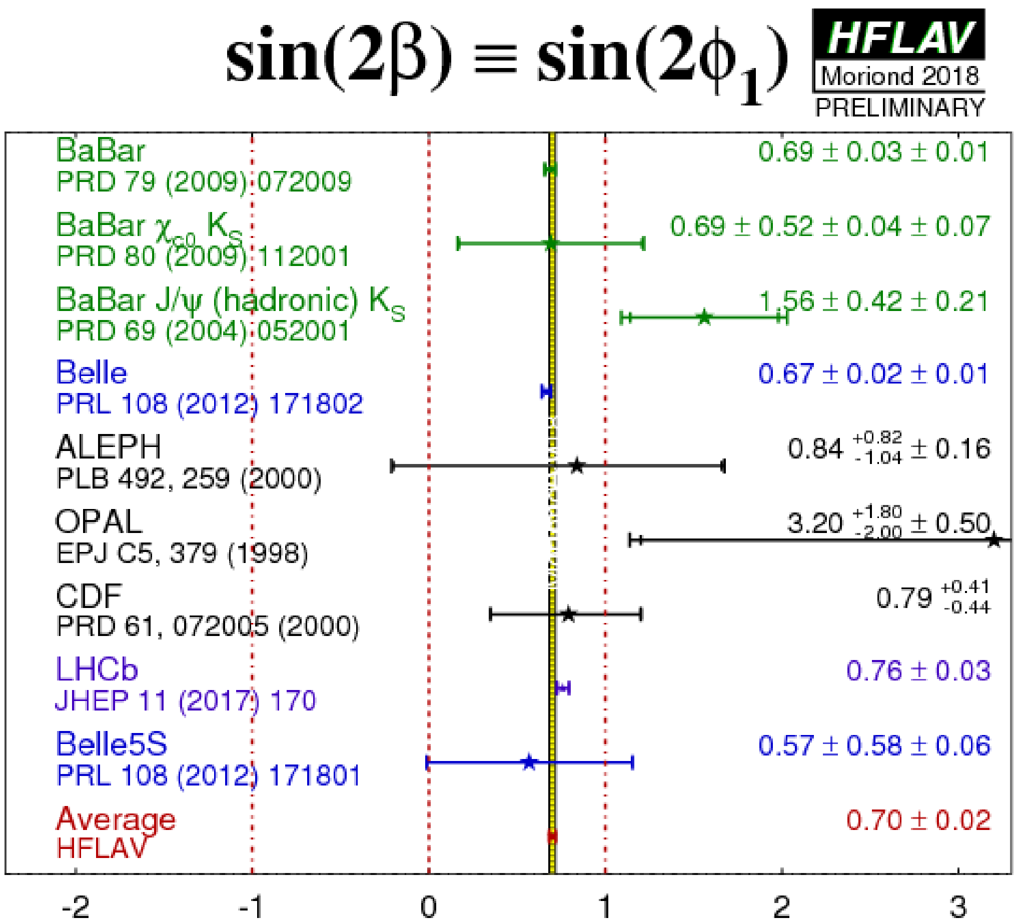
$$a_{f_{CP}}(t) = \frac{\text{Prob}(B^0(t) \to f_{CP}) - \text{Prob}(\overline{B^0}(t) \to f_{CP})}{\text{Prob}(\overline{B^0}(t) \to f_{CP}) + \text{Prob}(B^0(t) \to f_{CP})} = C_f \cos\Delta m_d t \; + \; S_f \sin\Delta m_d t$$

$$a_{f_{CP}}(t) = -\eta_{CP} \sin\Delta m_d \Delta t \sin 2\beta$$

Latest sin2β results:



raw asymmetry
as function of Δt



$\sin 2\beta(J/\psi K^0) = 0.698 \pm 0.017$

HFLAV

data-driven theoretical uncertainty

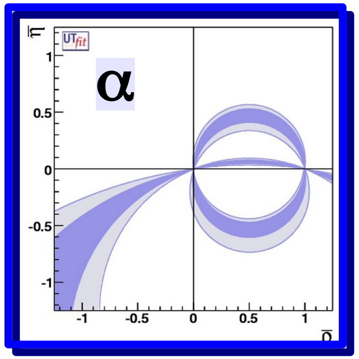
$\Delta S = -0.01 \pm 0.01$

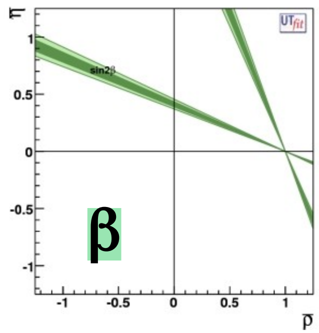
$\sin 2\beta(J/\psi K^0) = 0.688 \pm 0.020$

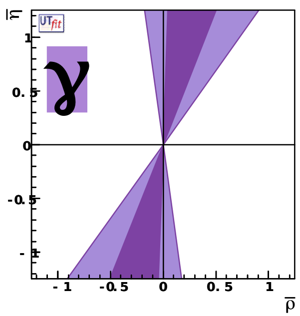
UTfit input

M.Ciuchini, M.Pierini, L.Silvestrini
Phys. Rev. Lett. 95, 221804 (2005)

angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:







α: CP violation in $B^0 \rightarrow \pi^+\pi^-$

- considering the tree (T) only:

  $\lambda_{\pi\pi} = e^{2i\alpha}$

  $C_{\pi\pi} = 0$

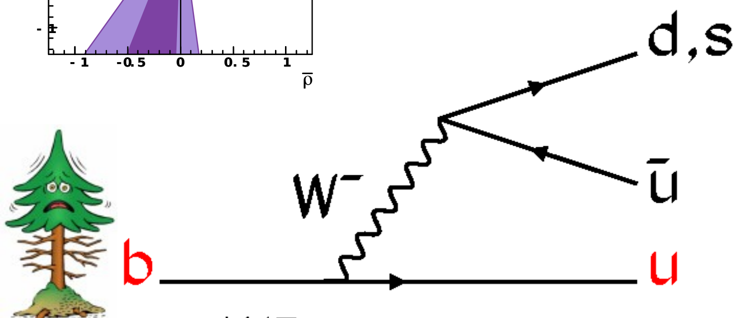
  $S_{\pi\pi} = \sin(2\alpha)$

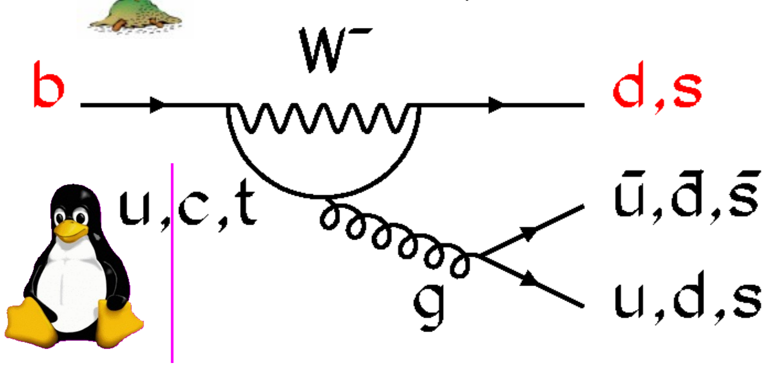
- adding the penguins (P):

$$\lambda_{\pi\pi} = e^{2i\alpha}\frac{1 + |P/T|e^{i\delta}e^{i\gamma}}{1 + |P/T|e^{i\delta}e^{-i\gamma}}$$

$$C_{\pi\pi} \propto \sin(\delta)$$

$$S_{\pi\pi} = \sqrt{1 - C^2_{\pi\pi}}\sin(2\alpha_{eff})$$





angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:

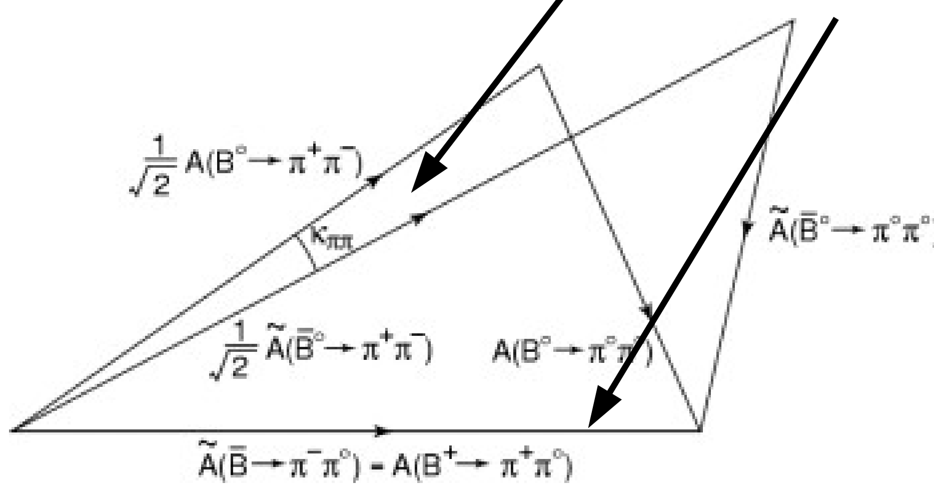
from α_{eff} to α: isospin analysis

- $B \to \pi^+\pi^-$, $\pi^+\pi^0$, $\pi^0\pi^0$ decays are connected from isospin relations
  - $\pi\,\pi$ states can have I = 2 or I = 0
  - the gluonic penguins contribute only to the I = 0 state (ΔI=1/2)
  - $\pi^+\pi^0$ is a pure I = 2 state (ΔI = 3/2) and it gets contribution only from the tree diagram
  - triangular relations allow for the determination of the phase difference induced on a:

$$2\alpha_{eff} = 2\alpha + \kappa_{\pi\pi}$$

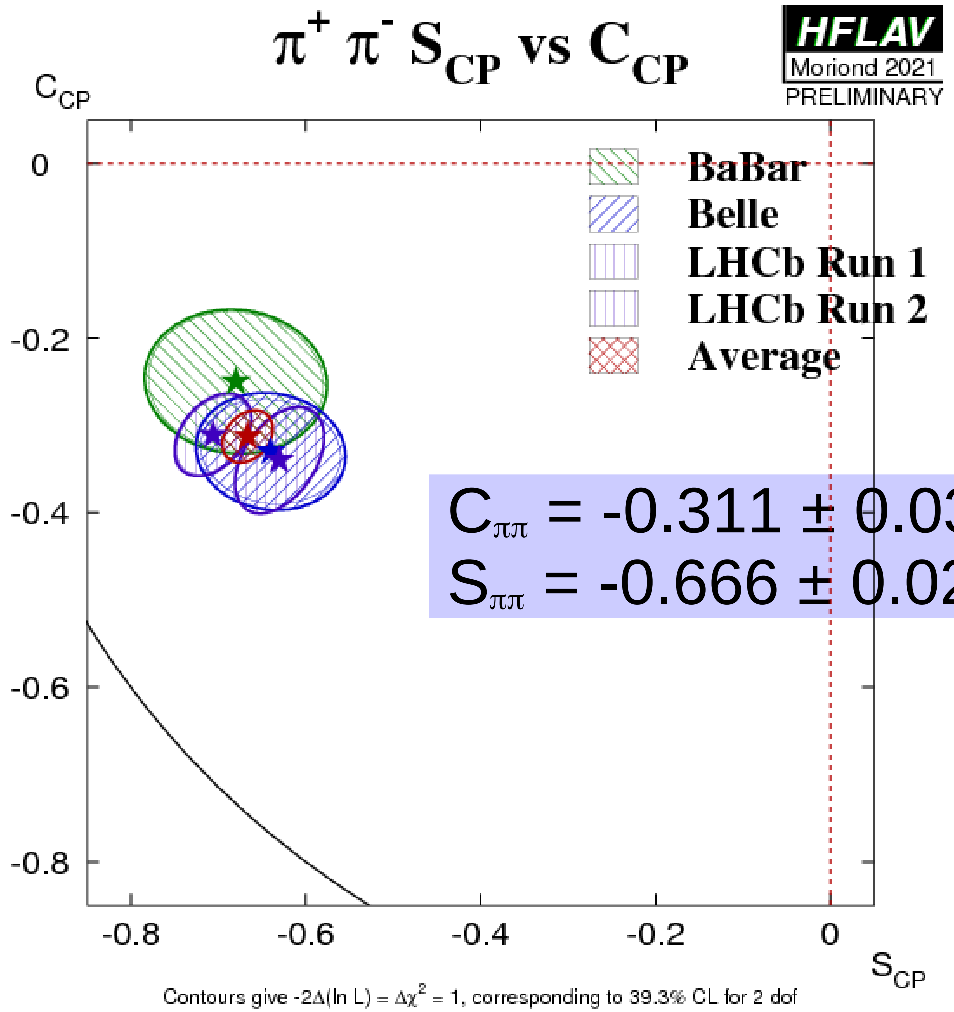
$$|A^{+0}|=|A^{-0}|$$

Both BR(B^0) and BR($\bar{B}^0$) have to be measured in all the $\pi\pi$ channels

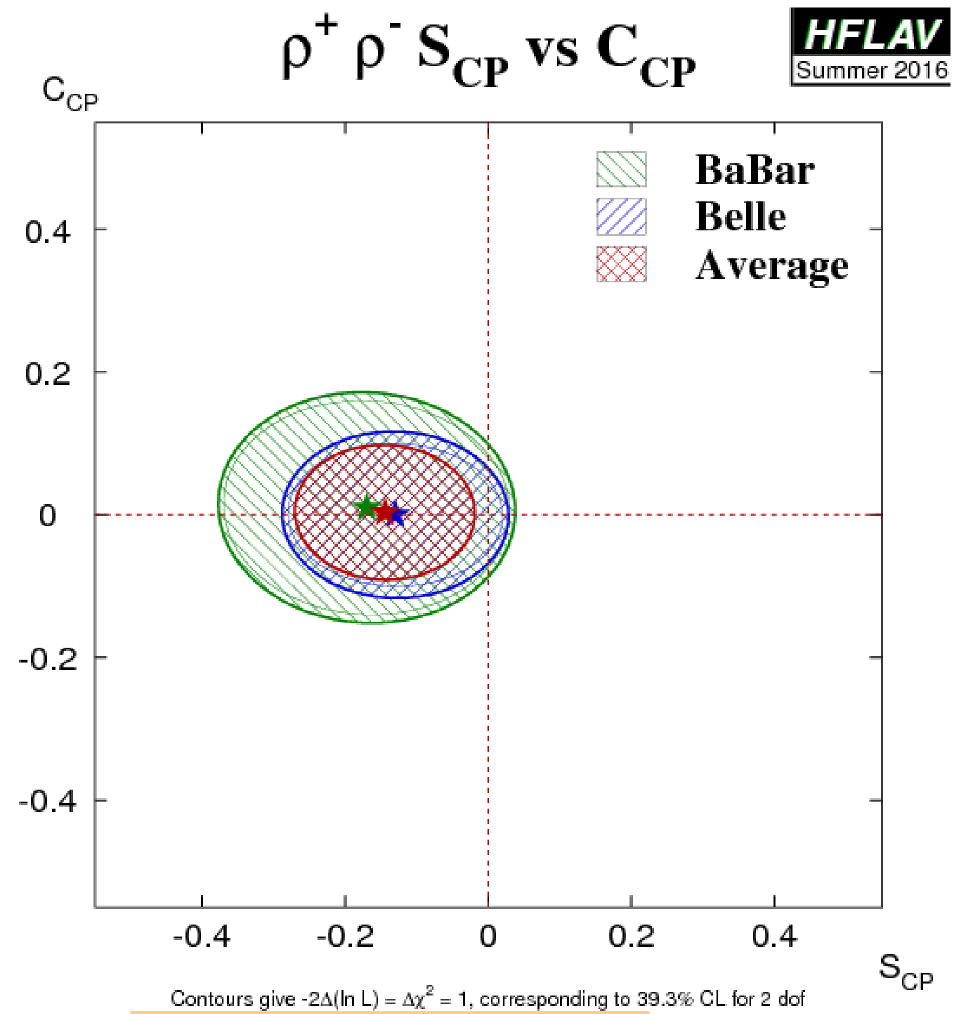


angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:

α result for $\pi^+\pi^-$ and $\rho^+\rho^-$

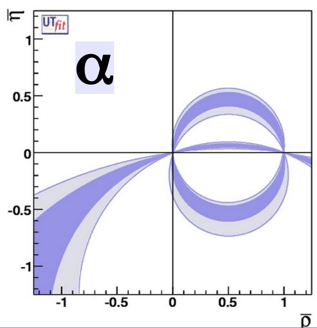


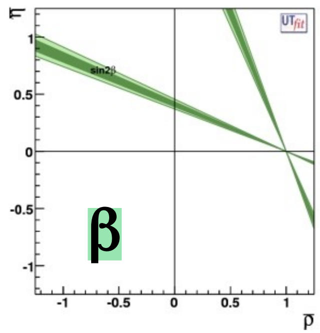
$C_{\pi\pi}$ = -0.311 ± 0.030
$S_{\pi\pi}$ = -0.666 ± 0.029

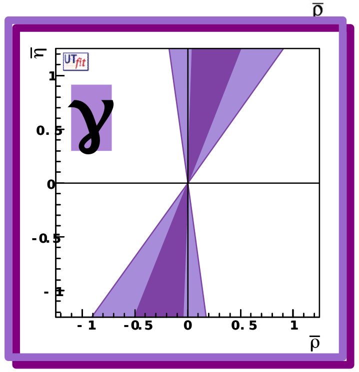


$C_{\rho\rho}$ = 0.00 ± 0.09
$S_{\rho\rho}$ = -0.14 ± 0.13

angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:

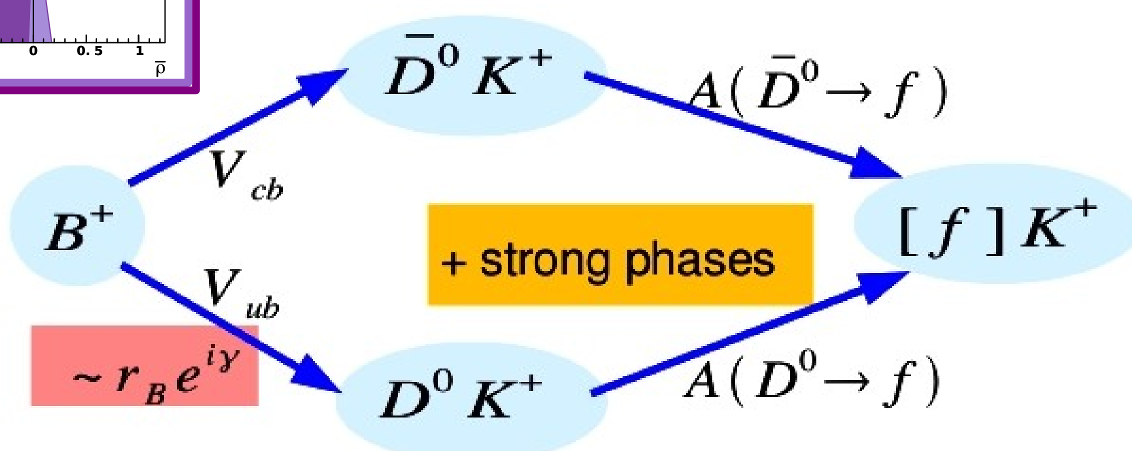


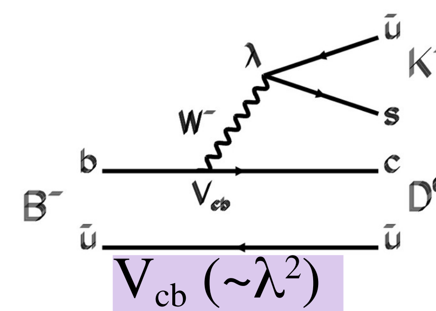


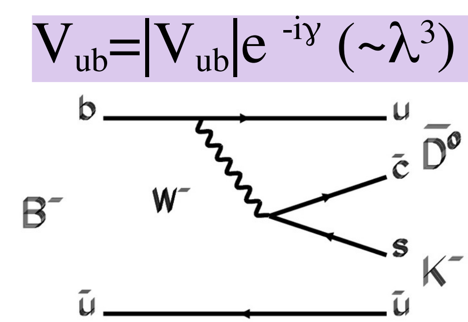


γ and DK trees

- $D^{(*)}K^{(*)}$ decays: from BRs and BR ratios, no time-dependent analysis, just rates
- the phase γ is measured exploiting interferences: two amplitudes leading to the same final states
- some rates can be really small: ~ 10^{-7}

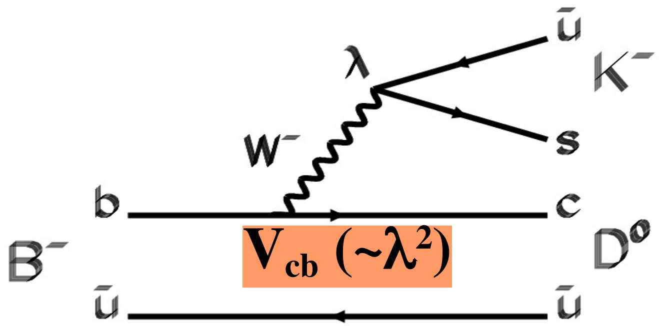


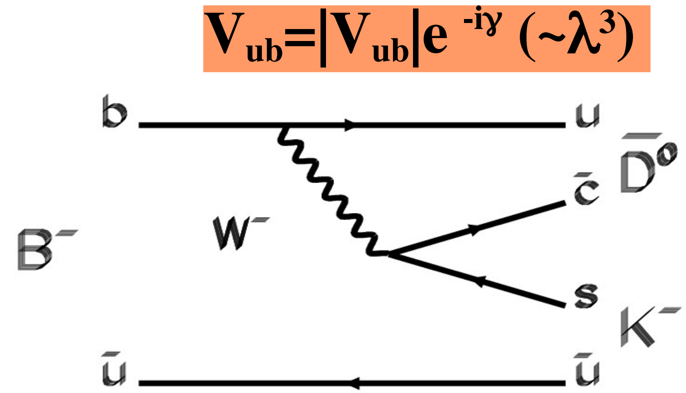




$B \to D^{(*)0}(\bar{D}^{(*)0})K^{(*)}$ decays can proceed both through V_{cb} and V_{ub} amplitudes

sensitivity to γ: the ratio r_B





δ_B = strong phase diff.

$$A(B^- \to D^0 K^-) = A_B \qquad A(B^- \to \bar{D}^0 K^-) = A_B r_B e^{i(\delta_B - \gamma)}$$

$$A(B^+ \to \bar{D}^0 K^+) = A_B \qquad A(B^+ \to D^0 K^+) = A_B r_B e^{i(\delta_B + \gamma)}$$

r_B = amplitude ratio

$$r_B = \left| \frac{B^- \to \bar{D}^0 K^-}{B^- \to D^0 K^-} \right| = \sqrt{\bar{\eta}^2 + \bar{\rho}^2} \times F_{CS}$$

~0.36

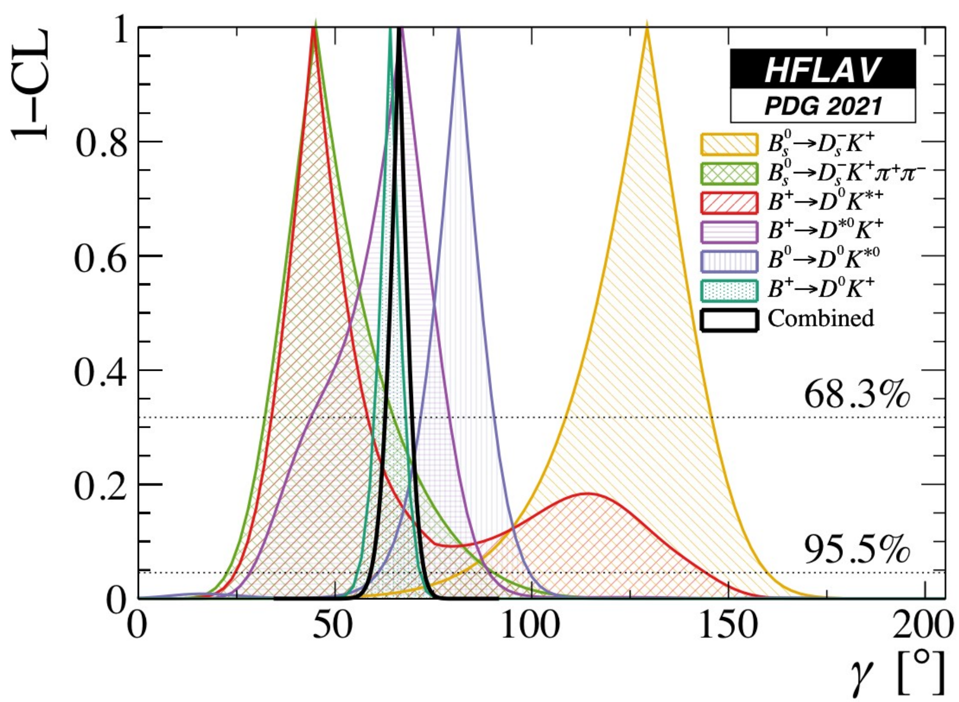
hadronic contribution channel-dependent

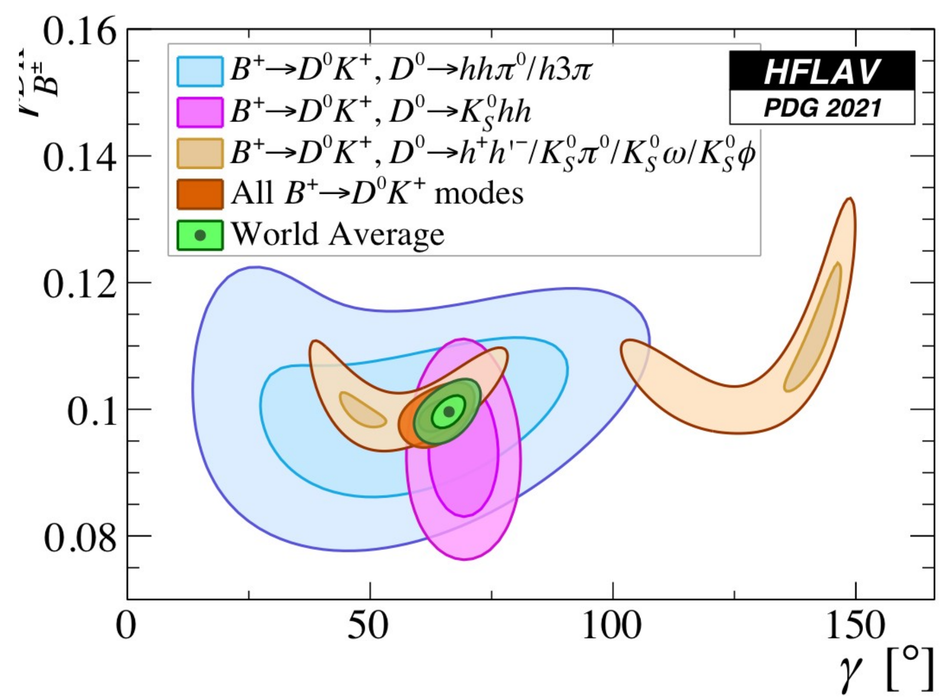
- in B^+ -> $D^{(*)0}K^+$: r_B is ~0.1
- while in B^0 to $D^{(*)0}K^0$ r_B is ~0.25
- Also measured: $r_B(DK)$, $r^*_B(D^*K)$ and $r^s_B(DK^*)$

angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:

γ and DK trees

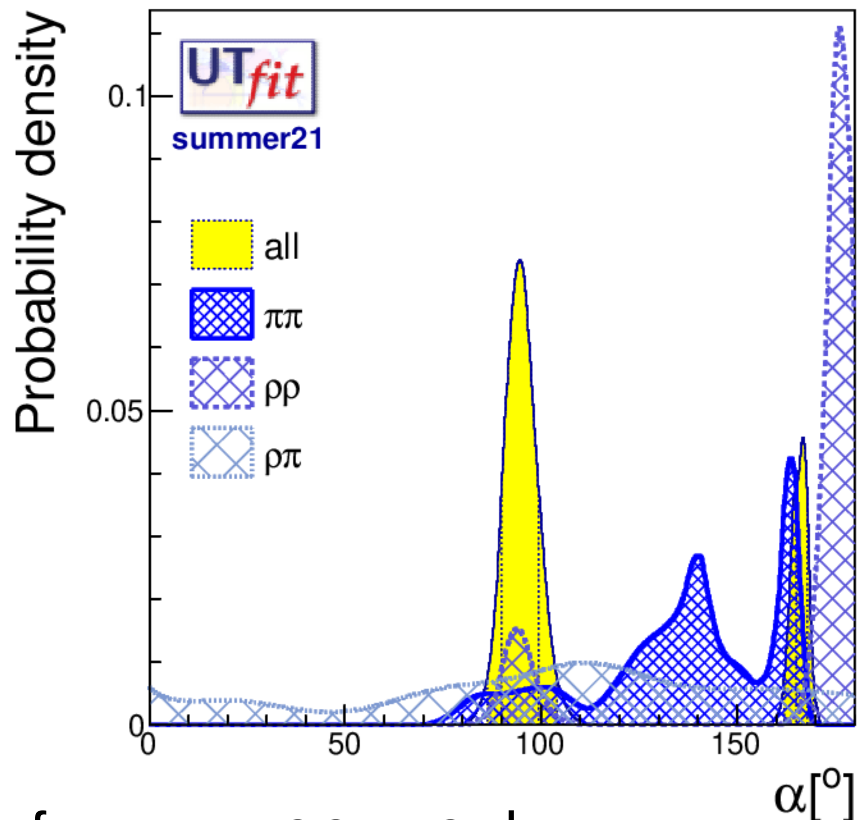
Parameter: $\gamma \equiv \varphi_3$ from all B → DK and similar b → cu-bar s & b → uc-bar s modes	
$\gamma \equiv \varphi_3$	$(66.2\ ^{+3.4}_{-3.6})°$
$r_B(DK^+) = 0.0996 \pm 0.0026$	$\delta_B(DK^+) = (128.0\ ^{+3.8}_{-4.0})°$
$r_B(D^*K^+) = 0.104\ ^{+0.013}_{-0.014}$	$\delta_B(D^*K^+) = (314.9\ ^{+7.8}_{-10.0})°$
$r_B(DK^{*+}) = 0.101\ ^{+0.016}_{-0.037}$	$\delta_B(DK^{*+}) = (49\ ^{+61}_{-16})°$
$r_B(DK^{*0}) = 0.257\ ^{+0.021}_{-0.022}$	$\delta_B(DK^{*0}) = (194\ ^{+9.5}_{--8.8})°$



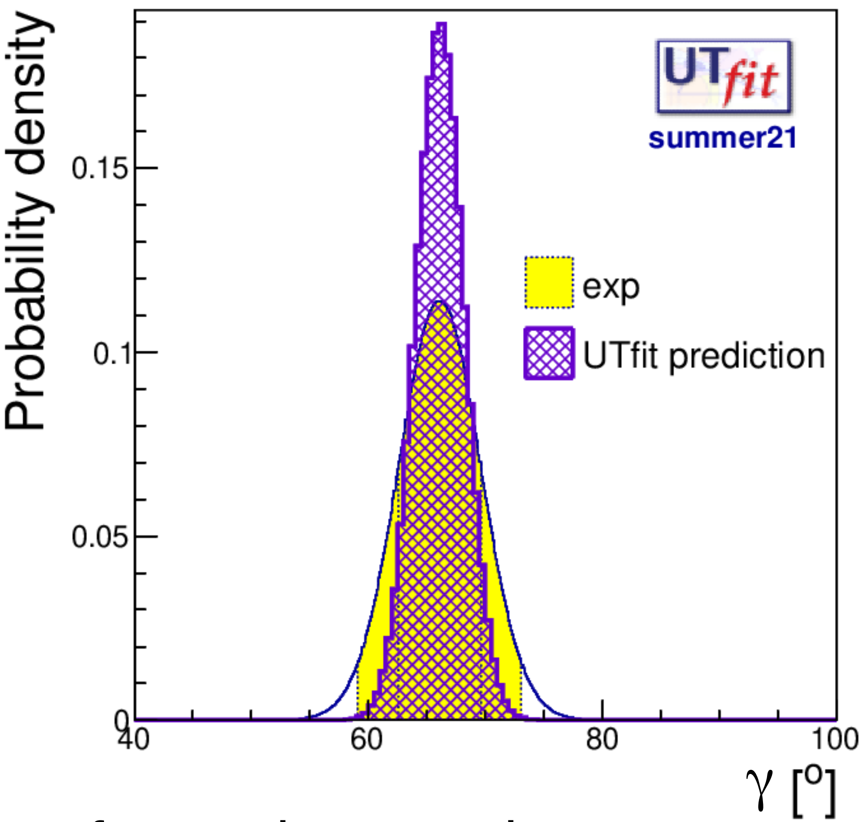


sin2α (ϕ_2) and γ (ϕ_3)

α updated with latest ππ/ρρ BR and C/S results



γ updated with all the latest results (LHCb)

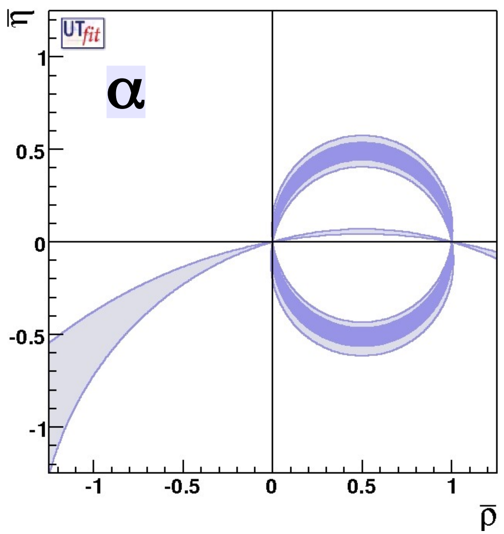


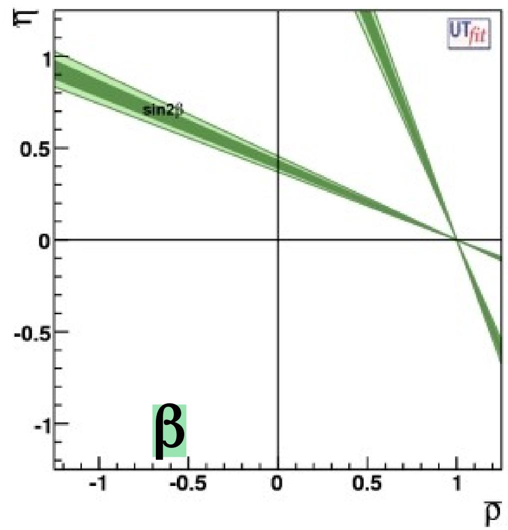
α from ππ, ρρ, πρ decays:
combined SM: (93.6 ± 4.2)°
UTfit prediction: (90.5 ± 2.1)°

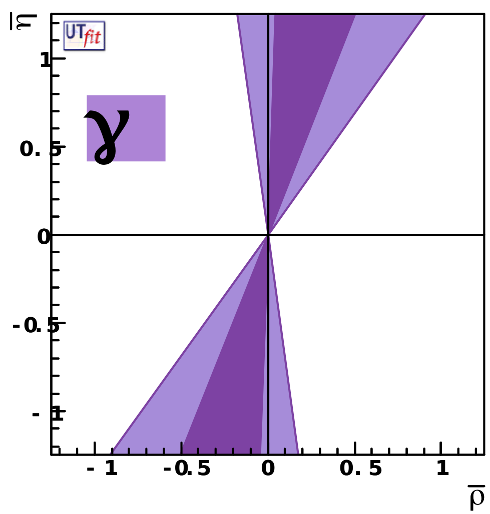
α from HFLAV: 85.5 ± 4.6

γ from B into DK decays:
HFLAV: (66.1 ± 3.5)°
UTfit prediction: (66.1 ± 2.1)°

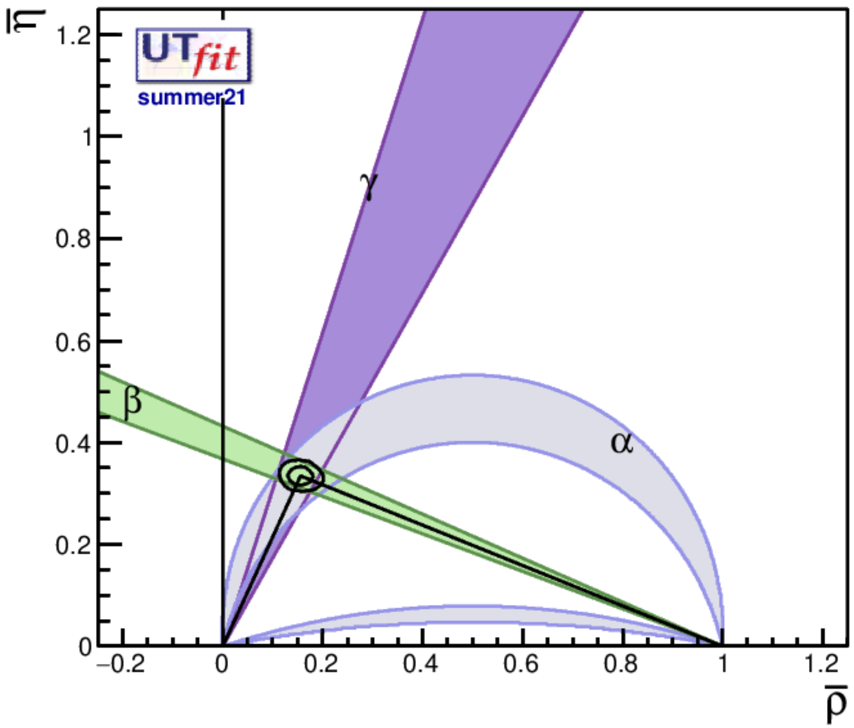
angle constraints in the $\bar{\rho}$-$\bar{\eta}$ plane:







levels @
95% Prob

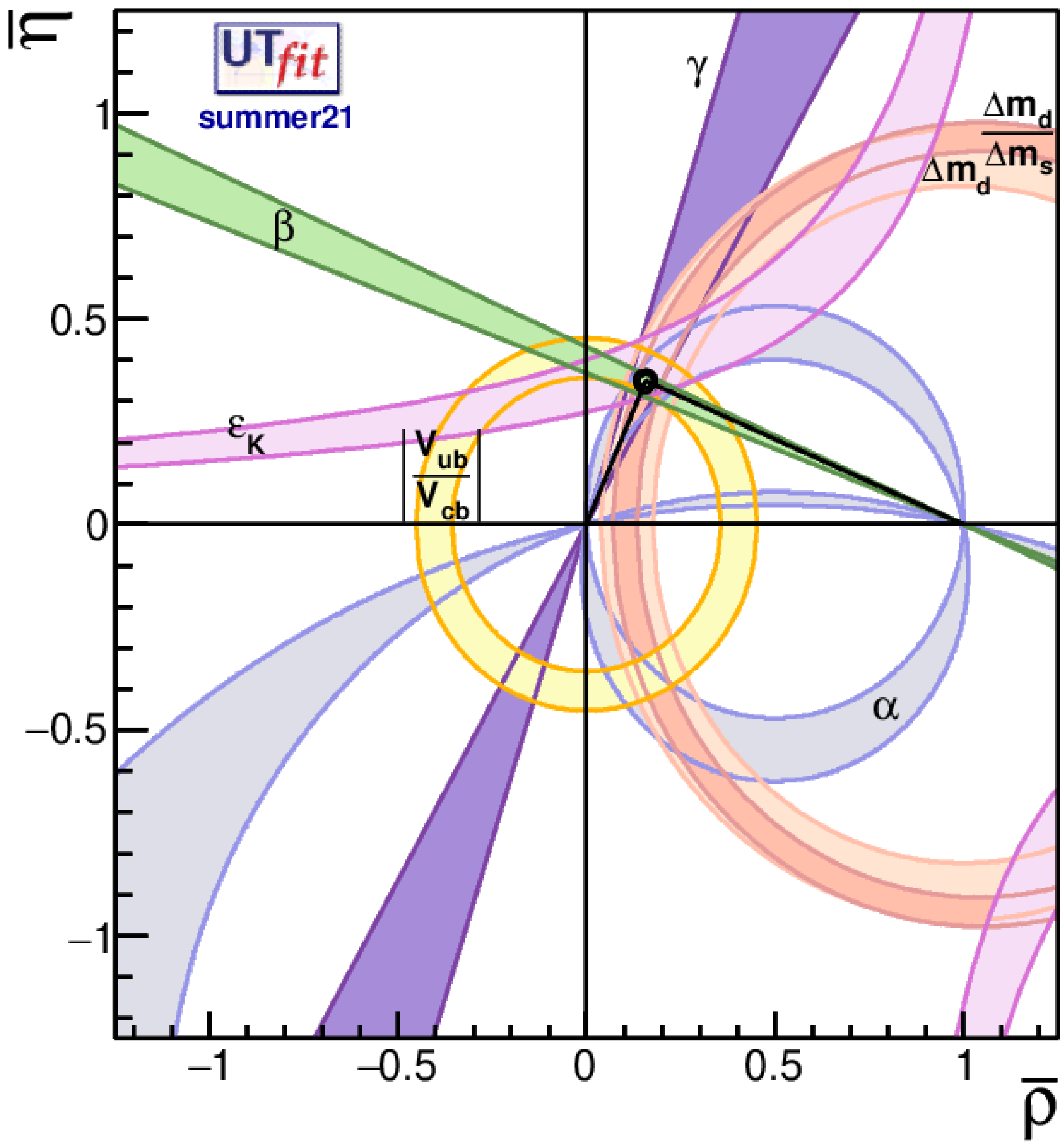


~12%

$\bar{\rho} = 0.156 \pm 0.018$
$\bar{\eta} = 0.335 \pm 0.018$

~5%

Unitarity Triangle analysis in the SM:



levels @
95% Prob

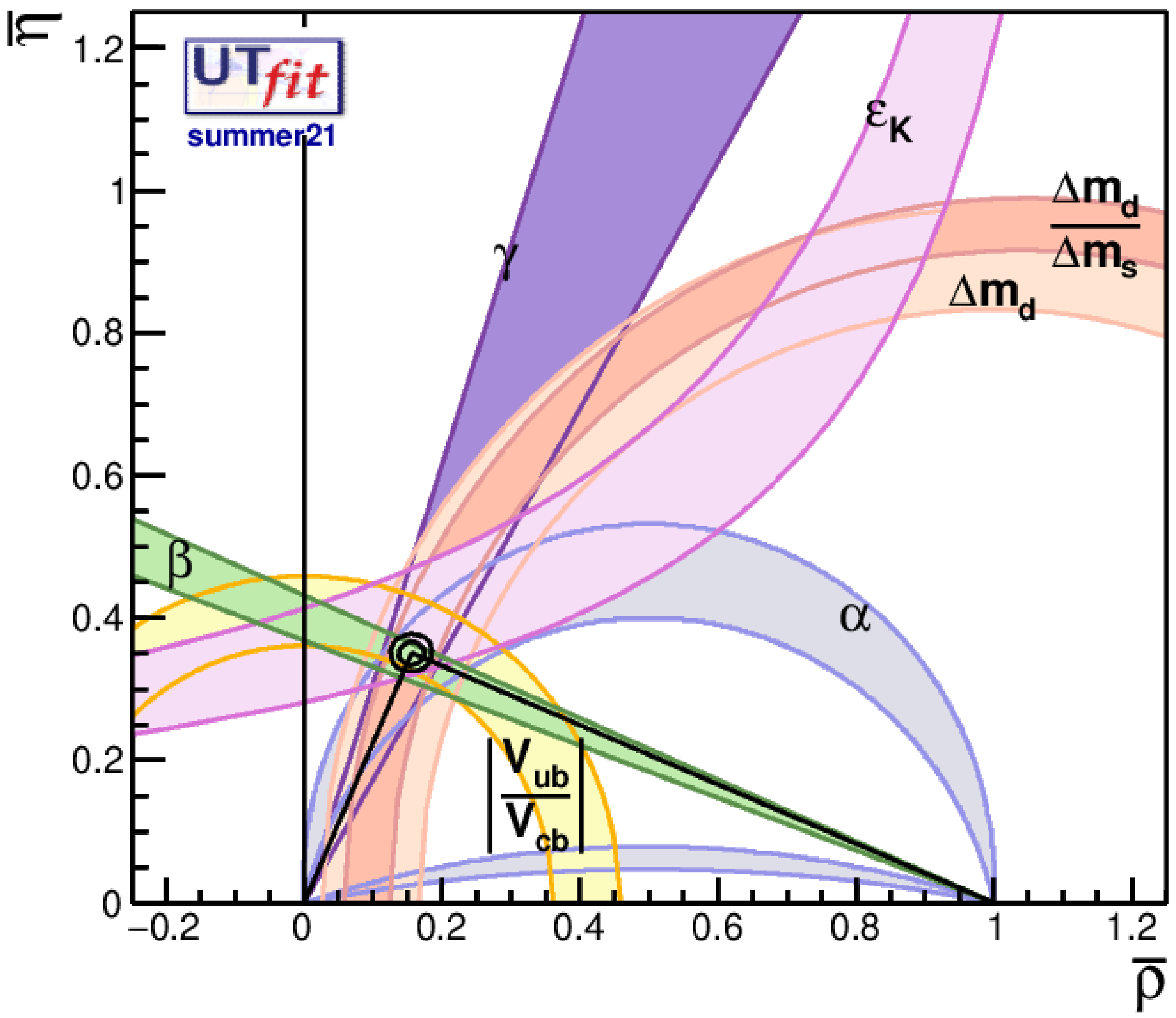
~8%

$\bar{\rho}$ = 0.157 ± 0.012
$\bar{\eta}$ = 0.350 ± 0.010

~3%

Unitarity Triangle analysis in the SM:

zoomed in..



levels @
95% Prob

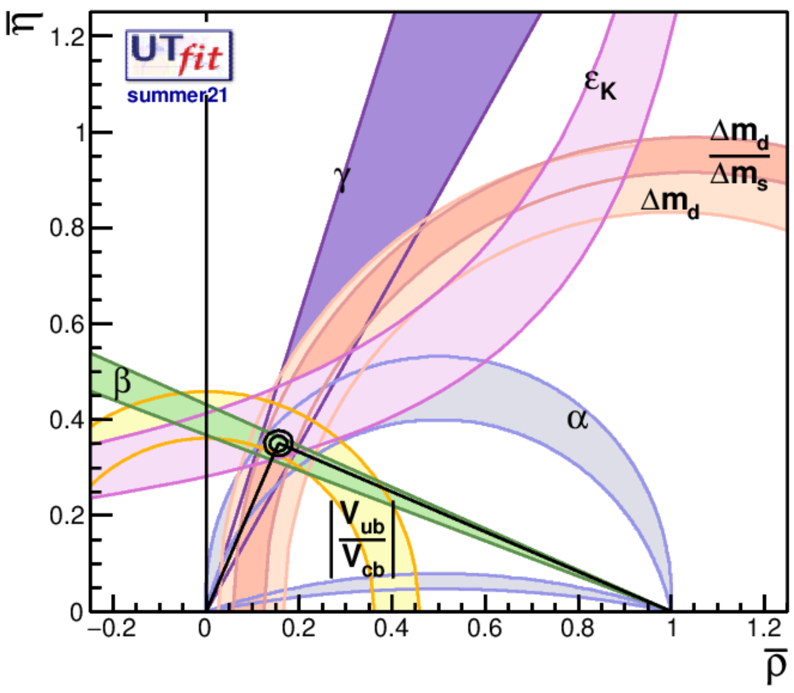
~8%

$\bar{\rho}$ = 0.157 ± 0.012
$\bar{\eta}$ = 0.350 ± 0.010

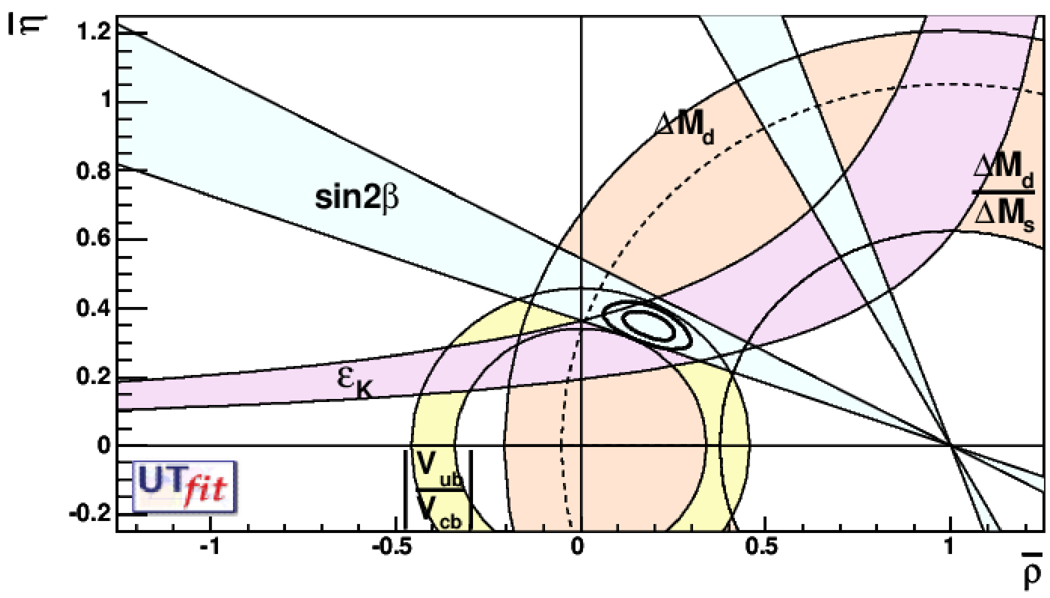
~3%

Unitarity Triangle analysis in the SM:

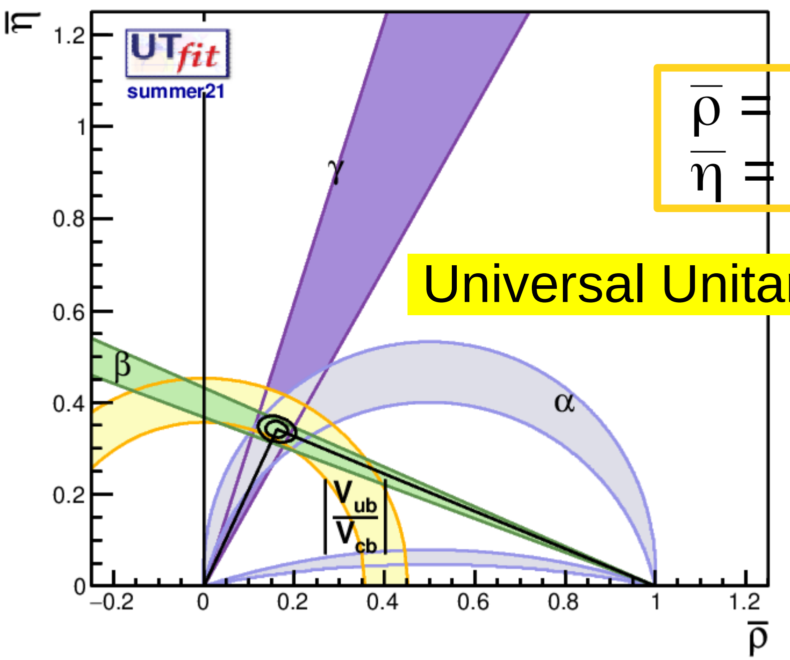
2021



2004



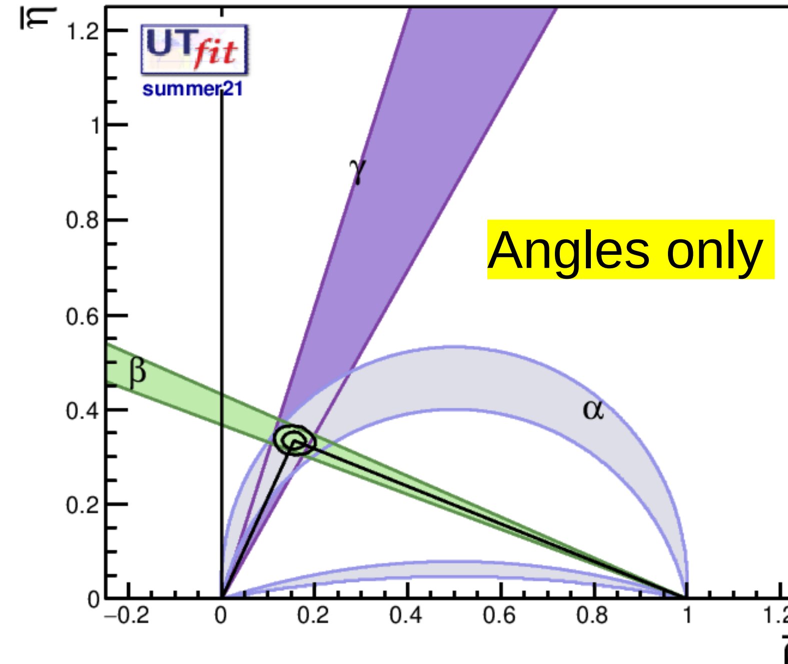
Some interesting configurations



Universal Unitary Triangle

$\bar{\rho} = 0.162 \pm 0.017$ (~10%)

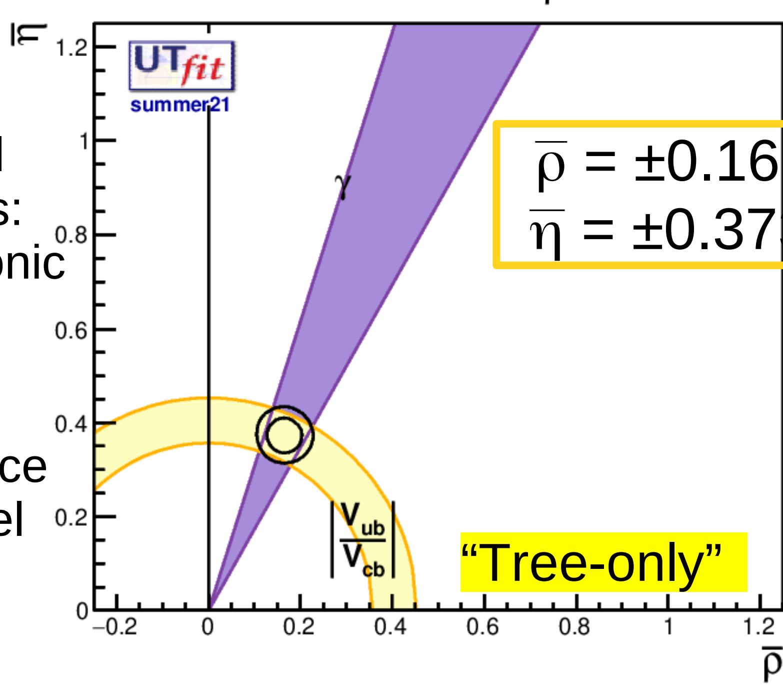
$\bar{\eta} = 0.341 \pm 0.011$ (~3%)



Angles only

Tree-level processes: Semileptonic and DK B decays

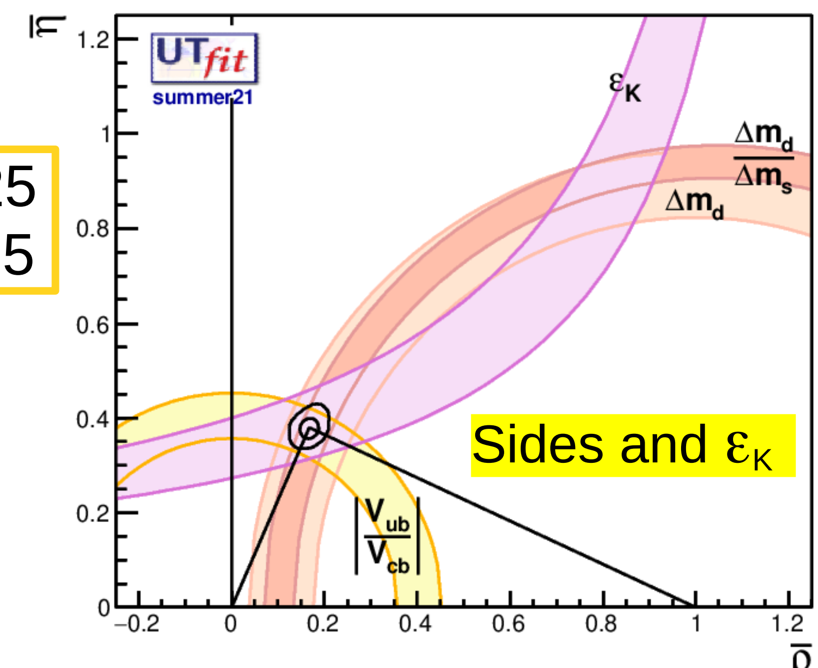
→ reference for model building



"Tree-only"

$\bar{\rho} = \pm 0.165 \pm 0.025$ (~15%)

$\bar{\eta} = \pm 0.373 \pm 0.025$ (~7%)



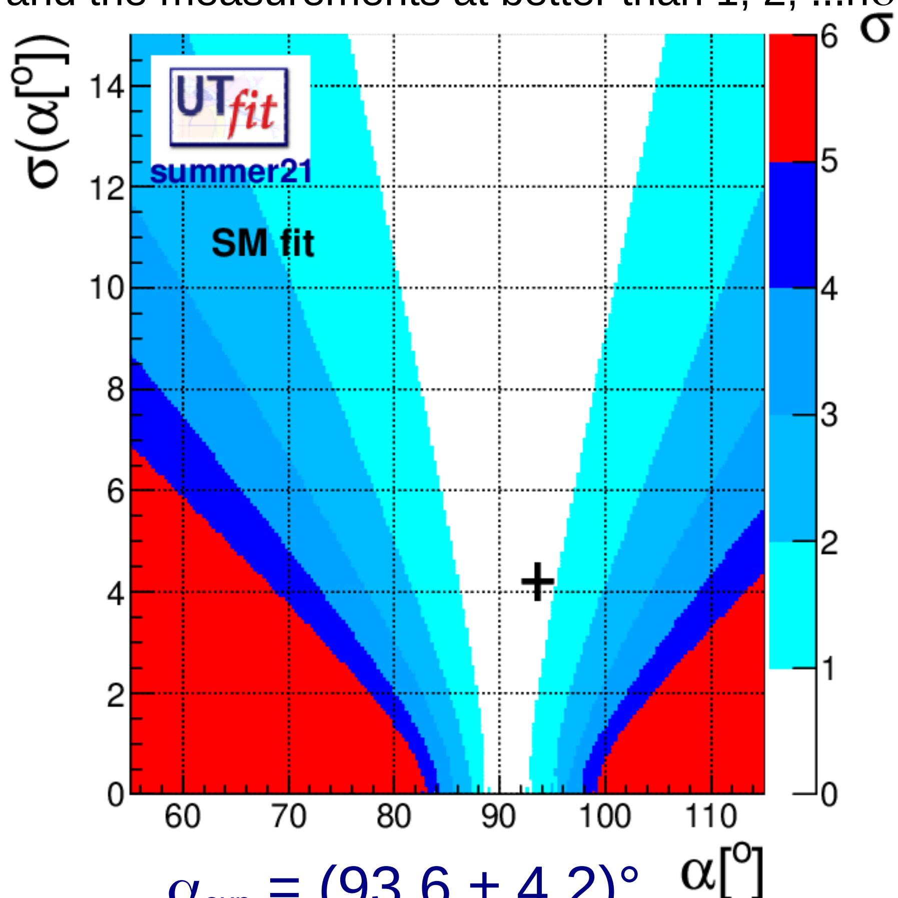
Sides and ε_K

compatibility plots

A way to “measure” the agreement of a single measurement with the indirect determination from the fit using all the other inputs: test for the SM description of the flavour physics

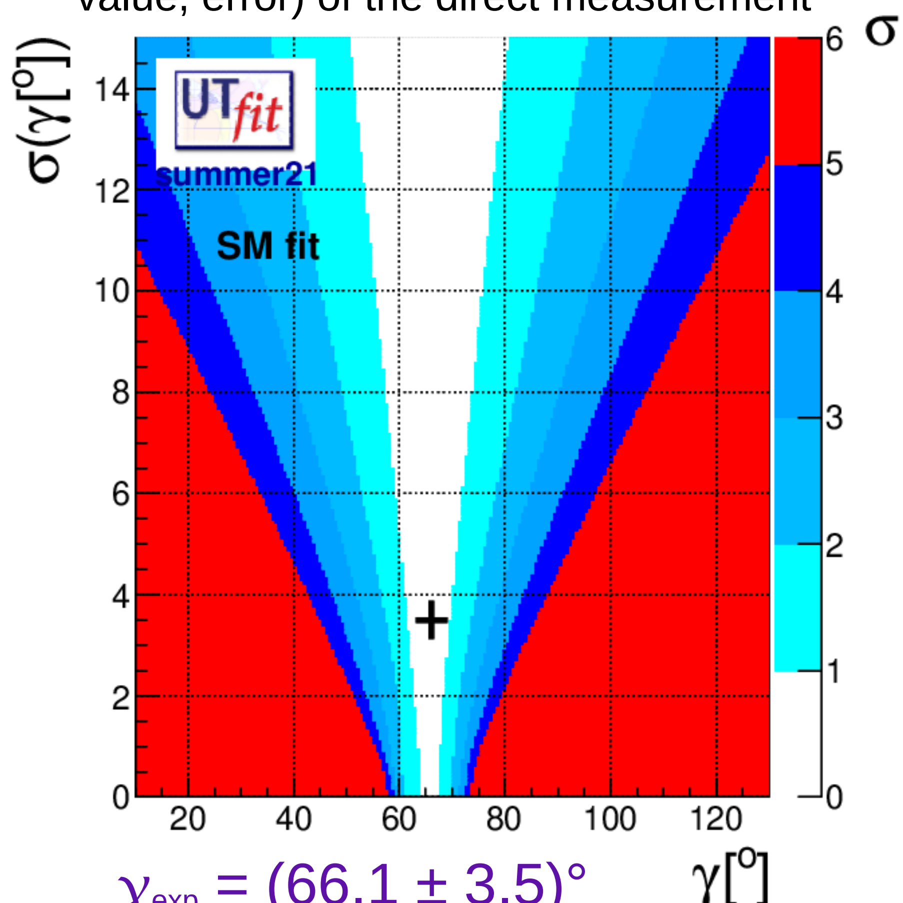
Color code: agreement between the predicted values and the measurements at better than 1, 2, ...nσ

The cross has the coordinates (x,y)=(central value, error) of the direct measurement



$\alpha_{exp} = (93.6 \pm 4.2)°$

$\alpha_{UTfit} = (90.5 \pm 2.1)°$



$\gamma_{exp} = (66.1 \pm 3.5)°$

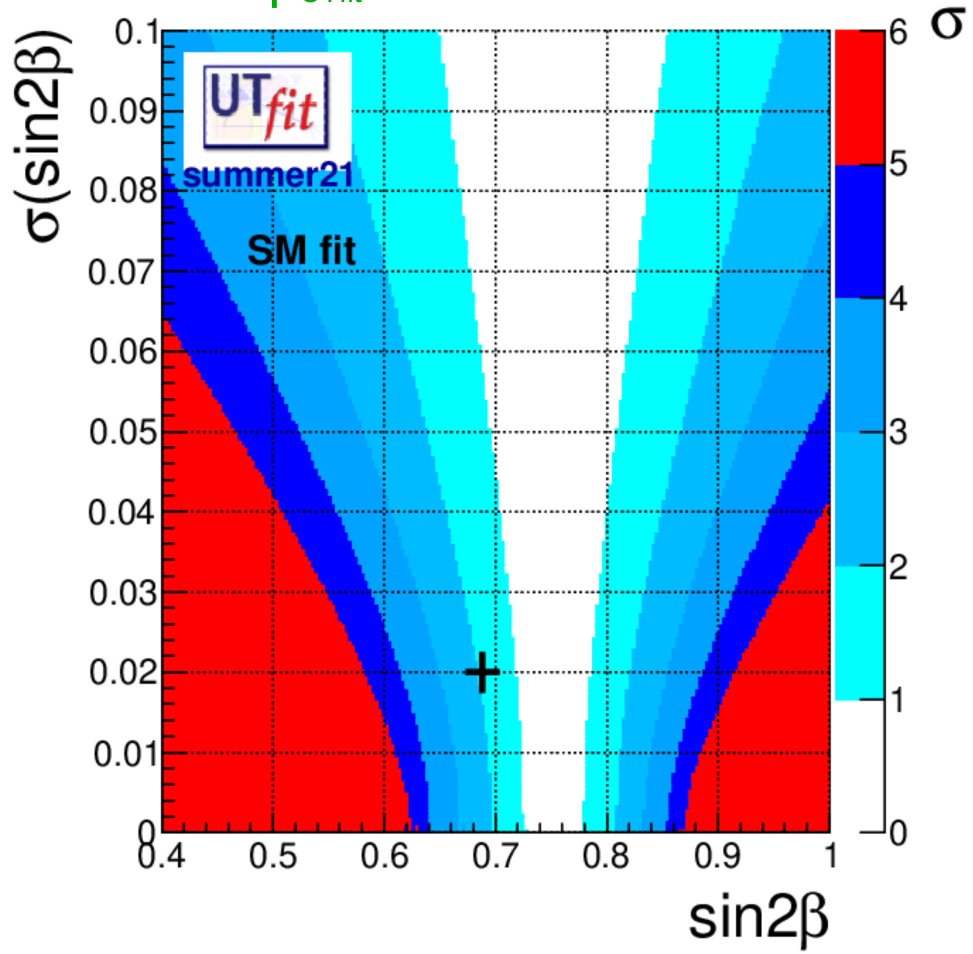
$\gamma_{UTfit} = (66.1 \pm 2.1)°$

Checking the usual *tensions*..

~1.4σ

$\sin 2\beta_{exp} = 0.688 \pm 0.020$

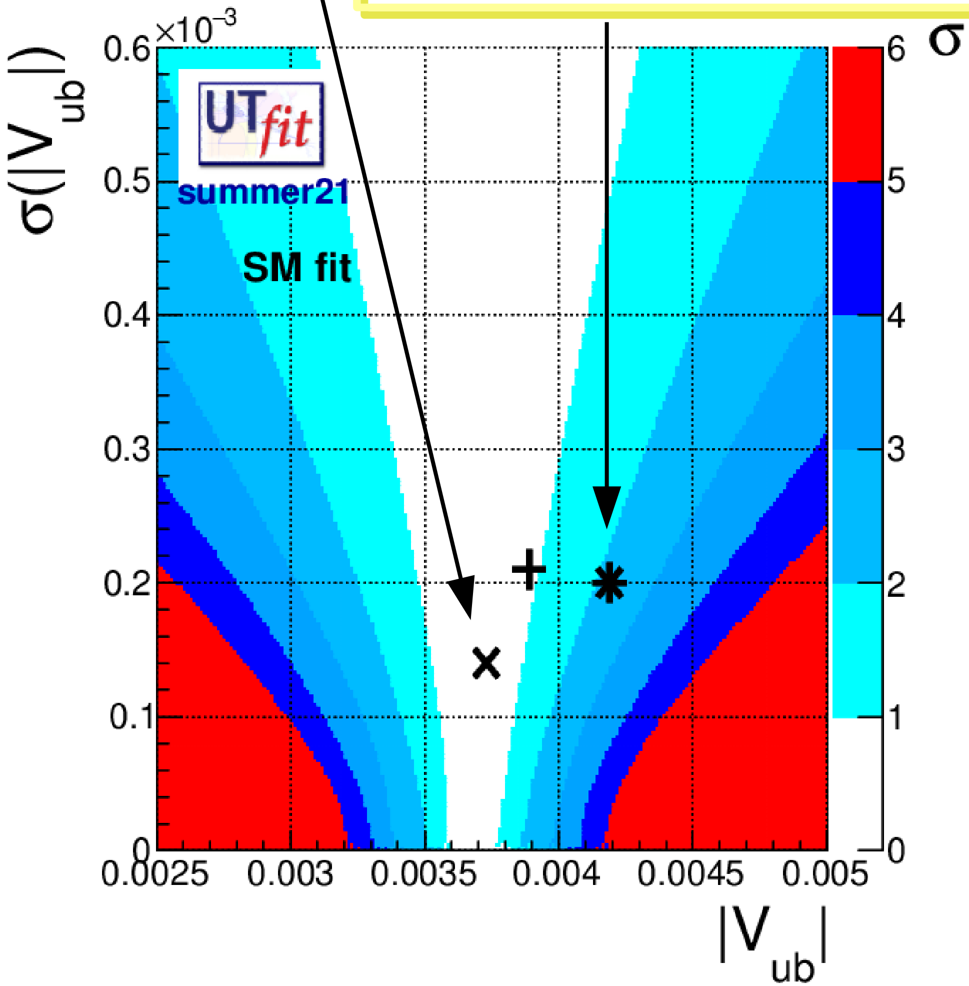
$\sin 2\beta_{UTfit} = 0.751 \pm 0.027$



Checking the usual *tensions*..

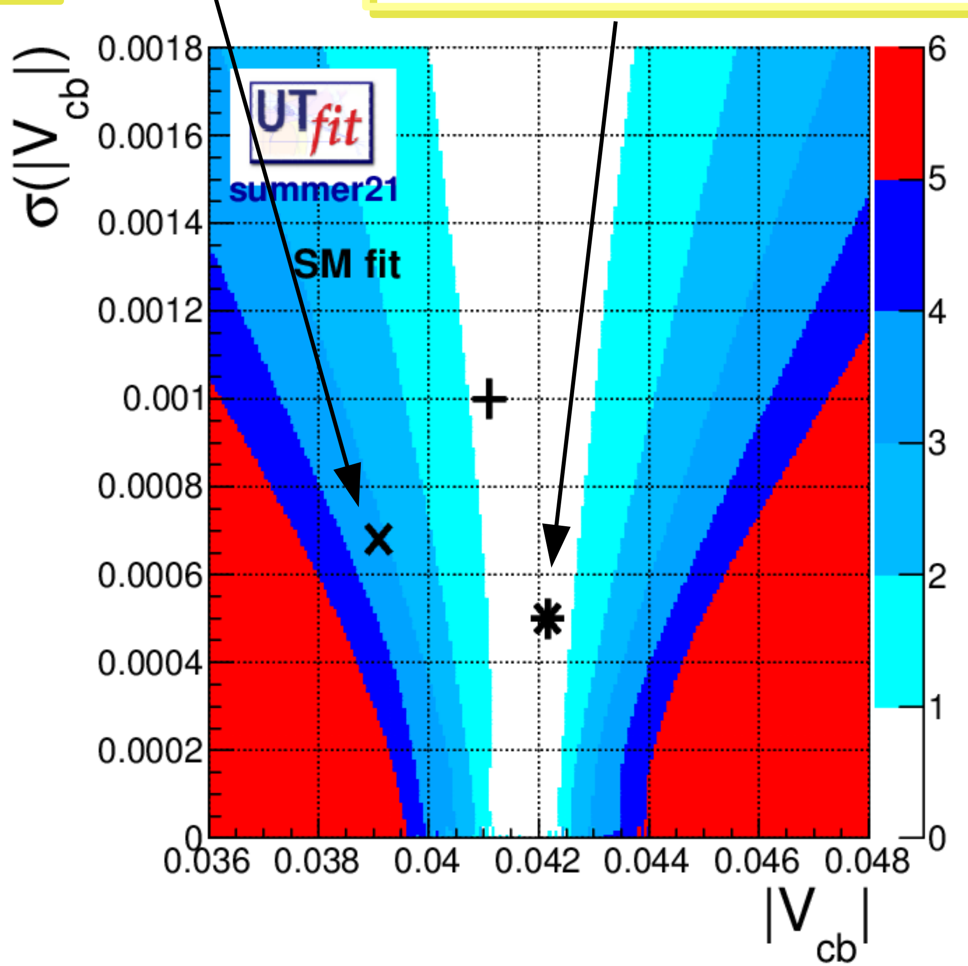
$|V_{ub}|$ *(excl)* = (3.73 ± 0.14) 10^{-3}

$|V_{ub}|$ *(incl)* = (4.19 ± 0.20) 10^{-3}



$|V_{cb}|$ *(excl)* = (39.09 ± 0.68) 10^{-3}

$|V_{cb}|$ *(incl)* = (42.16 ± 0.50) 10^{-3}



Vub_{exp} = (3.89 ± 0.21) · 10^{-3}

Vub_{UTfit} = (3.68 ± 0.10) · 10^{-3}

Vcb_{exp} = (41.1 ± 1.0) · 10^{-3}

Vcb_{UTfit} = (41.9 ± 0.5) · 10^{-3}

Unitarity Triangle analysis in the SM:

obtained excluding the given constraint from the fit

Observables	Measurement	Prediction	Pull (#σ)
$\sin 2\beta$	0.688 ± 0.020	0.751 ± 0.027	~ 1.4
γ	66.1 ± 3.5	66.1 ± 2.1	< 1
α	93.6 ± 4.2	90.5 ± 2.1	< 1
$\varepsilon_K \cdot 10^3$	2.228 ± 0.001	2.05 ± 0.13	~ 1.4
$\|V_{cb}\| \cdot 10^3$	40.4 ± 1.3	41.9 ± 0.5	< 1
$\|V_{cb}\| \cdot 10^3$ (incl)	42.16 0.50		< 1
$\|V_{cb}\| \cdot 10^3$ (excl)	39.09 0.68		~ 2.4
$\|V_{ub}\| \cdot 10^3$	3.89 ± 0.21	3.68 ± 0.10	< 1
$\|V_{ub}\| \cdot 10^3$ (incl)	4.19 ± 0.20	-	~ 1.7
$\|V_{ub}\| \cdot 10^3$ (excl)	3.73 ± 0.14	-	< 1
$BR(B \to \tau\nu)[10^{-4}]$	1.09 ± 0.24	0.87 ± 0.05	< 1
$A_{SL}^d \cdot 10^3$	-2.1 ± 1.7	-0.32 ± 0.03	< 1
$A_{SL}^s \cdot 10^3$	-0.6 ± 2.8	0.014 ± 0.001	< 1

UT analysis including new physics

Consider for example B_s mixing process.
Given the SM amplitude, we can define

$$C_{B_s} e^{-2i\phi_{B_s}} = \frac{\langle \overline{B}_s | H_{eff}^{SM} + H_{eff}^{NP} | B_s \rangle}{\langle \overline{B}_s | H_{eff}^{SM} | B_s \rangle} = 1 + \frac{A_{NP} e^{-2i\phi_{NP}}}{A_{SM} e^{-2i\beta_s}}$$

All NP effects can be parameterized in terms of one complex parameter for each meson mixing, to be determined in a simultaneous fit with the CKM parameters (now there are enough experimental constraints to do so).

For kaons we use *Re* and *Im*, since the two exp. constraints ε_K and Δm_K are directly related to them (with distinct theoretical issues)

$$C_{\varepsilon_K} = \frac{\mathrm{Im}\langle K^0 | H_{eff}^{full} | \overline{K}^0 \rangle}{\mathrm{Im}\langle K^0 | H_{eff}^{SM} | \overline{K}^0 \rangle}$$

$$C_{\Delta m_K} = \frac{\mathrm{Re}\langle K^0 | H_{eff}^{full} | \overline{K}^0 \rangle}{\mathrm{Re}\langle K^0 | H_{eff}^{SM} | \overline{K}^0 \rangle}$$

UT analysis including new physics

fit simultaneously for the CKM and
the NP parameters (generalized UT fit)

- add most general loop NP to all sectors
- use all available experimental info
- extract posteriors on NP contributions to ΔF=2 transitions

B_d and B_s mixing amplitudes
(2+2 real parameters):

$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})}\right) A_q^{SM} e^{2i\phi_q^{SM}}$$

$$\Delta m_{q/K} = C_{B_q/\Delta m_K} (\Delta m_{q/K})^{SM}$$

$$A_{CP}^{B_d \to J/\psi K_s} = \sin 2(\beta + \phi_{B_d})$$

$$A_{SL}^q = \mathrm{Im}\left(\Gamma_{12}^q / A_q\right)$$

$$\varepsilon_K = C_\varepsilon \varepsilon_K^{SM}$$

$$A_{CP}^{B_s \to J/\psi \phi} \sim \sin 2(-\beta_s + \phi_{B_s})$$

$$\Delta\Gamma^q / \Delta m_q = \mathrm{Re}\left(\Gamma_{12}^q / A_q\right)$$

new-physics-specific constraints

$$A^s_{\rm SL} \equiv \frac{\Gamma(\bar{B}_s \to \ell^+ X) - \Gamma(B_s \to \ell^- X)}{\Gamma(\bar{B}_s \to \ell^+ X) + \Gamma(B_s \to \ell^- X)} = {\rm Im}\left(\frac{\Gamma^s_{12}}{A^{\rm full}_s}\right)$$

semileptonic asymmetries in B^0 and B_s: sensitive to NP effects in both size and phase. Taken from the latest HFLAV.

Cleo, BaBar, Belle, D0 and LHCb

same-side dilepton charge asymmetry: admixture of B_s and B_d so sensitive to NP effects in both.

D0 arXiv:1106.6308

$$A^{\mu\mu}_{\rm SL} \times 10^3 = -7.9 \pm 2.0$$

$$A^{\mu\mu}_{\rm SL} = \frac{f_d \chi_{d0} A^d_{\rm SL} + f_s \chi_{s0} A^s_{\rm SL}}{f_d \chi_{d0} + f_s \chi_{s0}}$$

lifetime τ^{FS} in flavour-specific final states: average lifetime is a function to the width and the width difference

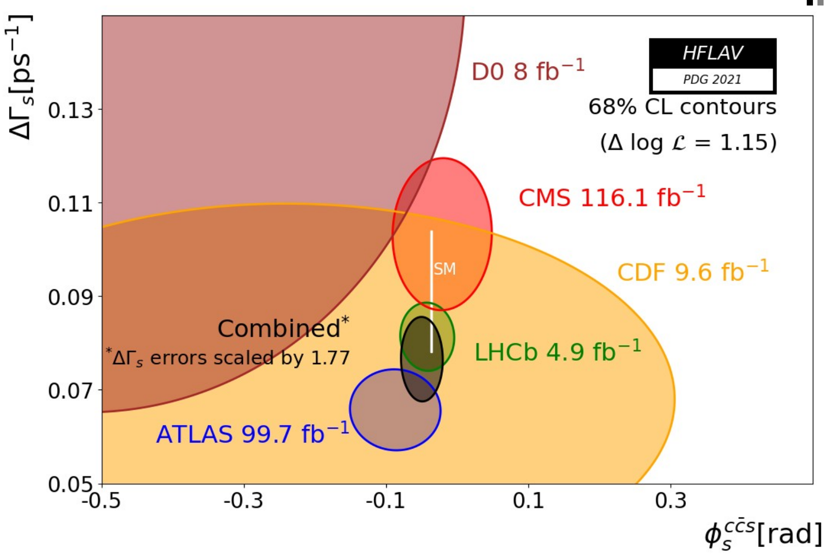
$$\tau^{FS}(B_s) = 1.527 \pm 0.011 \text{ ps}$$

HFLAV

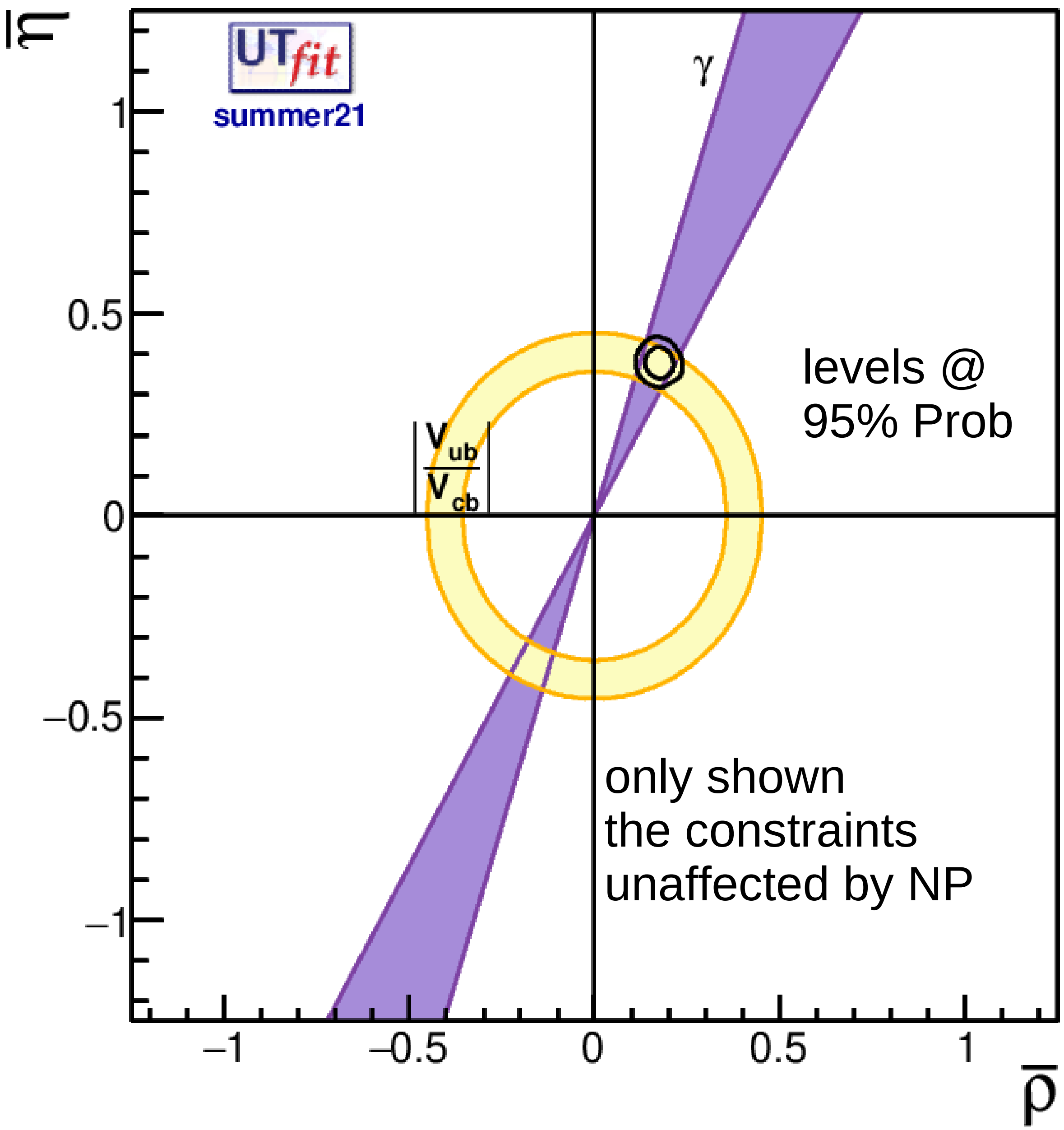
$\phi_s = 2\beta_s$ vs $\Delta\Gamma_s$ from $B_s \to J/\psi\phi$

angular analysis as a function of proper time and b-tagging

$$\phi_s = -0.050 \pm 0.019 \text{ rad}$$



NP analysis results



$\bar{\rho} = 0.175 \pm 0.027$
$\bar{\eta} = 0.380 \pm 0.026$

SM is

$\bar{\rho} = 0.157 \pm 0.012$
$\bar{\eta} = 0.350 \pm 0.010$

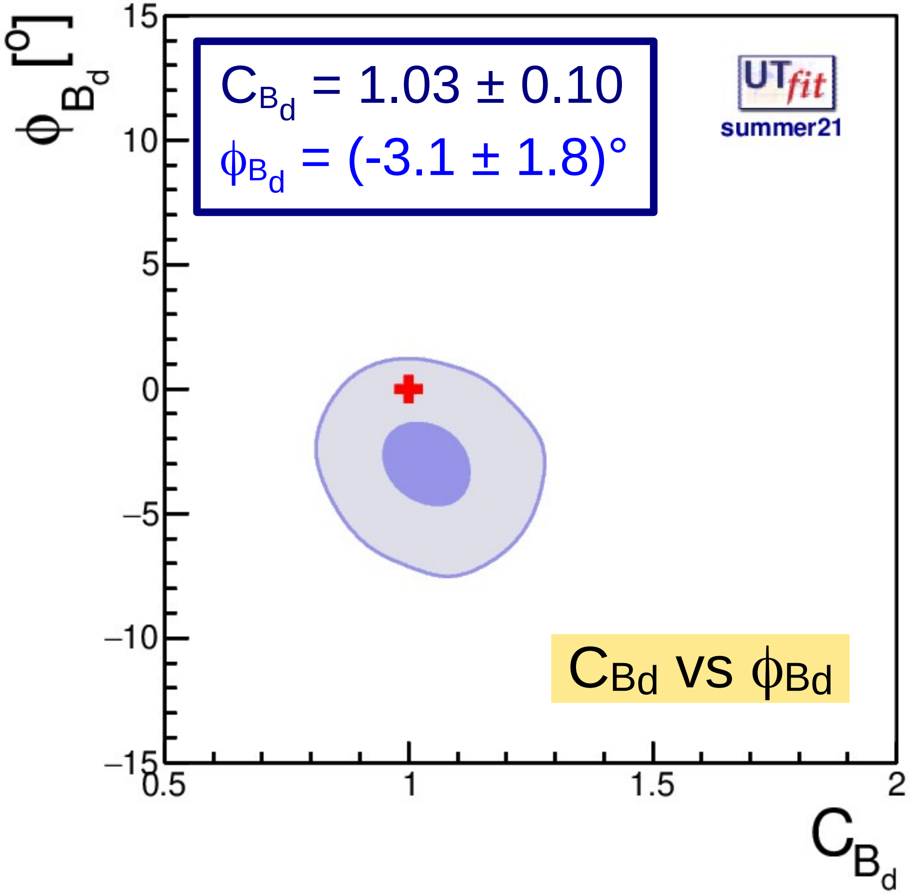
NP parameter results

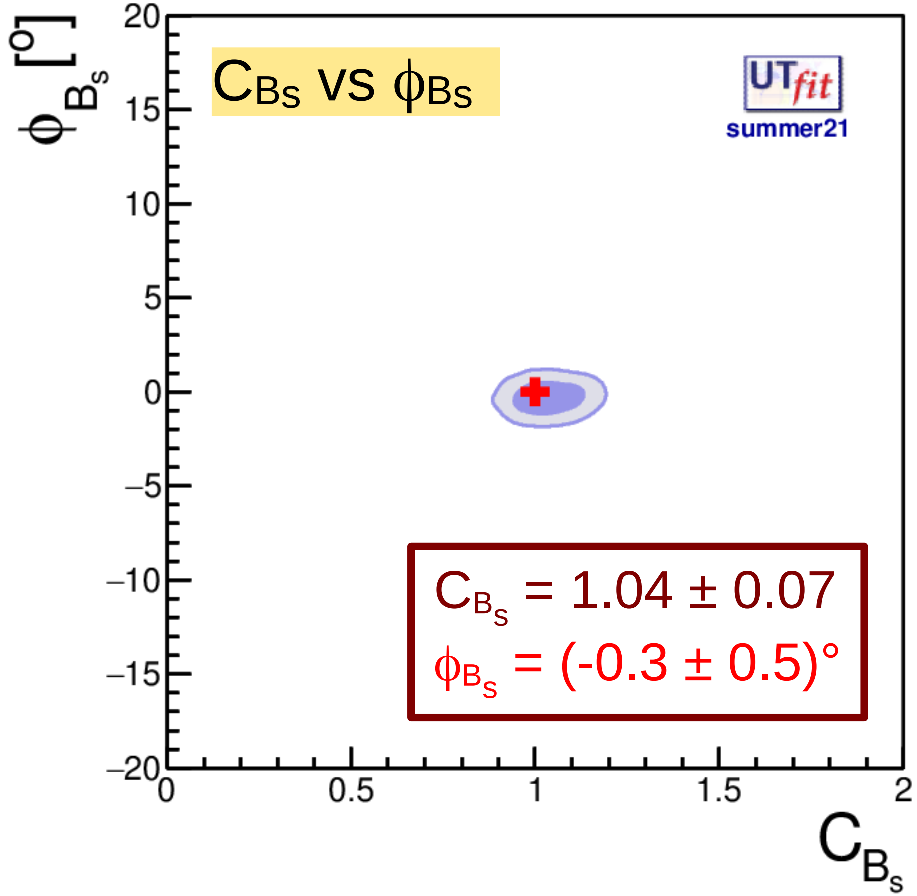
$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}}$$

dark: 68%
light: 95%
SM: red cross

K system

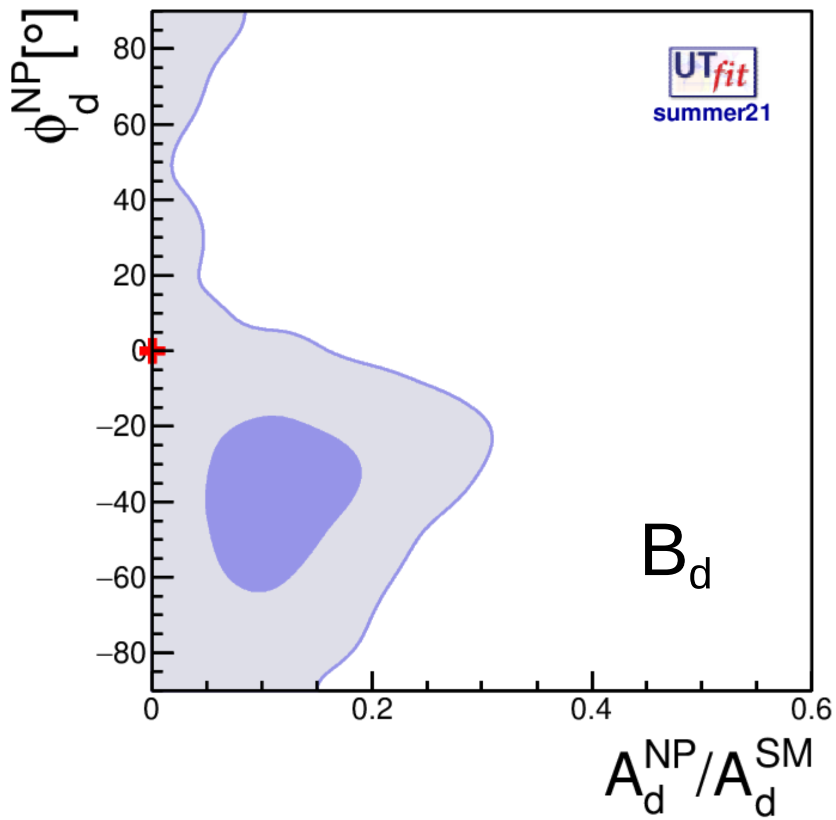
$C_{\epsilon_K} = 1.05 \pm 0.10$

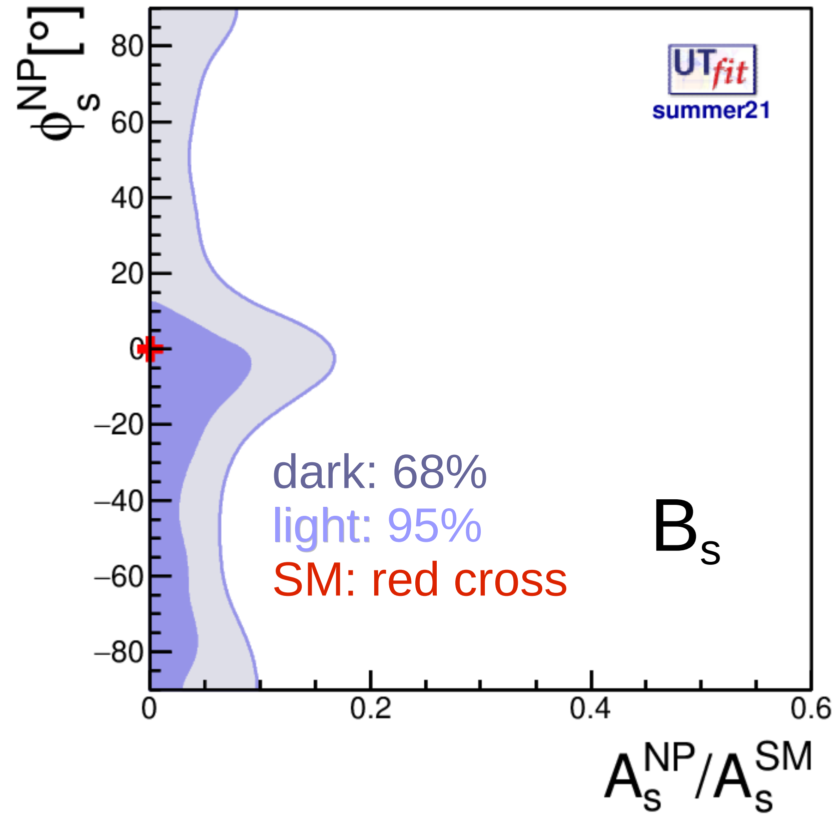




NP parameter results

$$A_q = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})}\right) A_q^{SM} e^{2i\phi_q^{SM}}$$





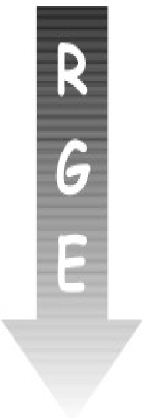
The ratio of NP/SM amplitudes is:

< 18% @68% prob. (30% @95%) in B_d mixing

< 10% @68% prob. (18% @95%) in B_s mixing

testing the new-physics scale

M. Bona *et al.* (UTfit)
JHEP 0803:049,2008
arXiv:0707.0636



At the high scale
new physics enters according to its specific features

At the low scale
use OPE to write the most general effective Hamiltonian.
the operators have different chiralities than the SM
NP effects are in the Wilson Coefficients C

$$\mathcal{H}_{\rm eff}^{\Delta B=2} = \sum_{i=1}^{5} C_i Q_i^{bq} + \sum_{i=1}^{3} \tilde{C}_i \tilde{Q}_i^{bq}$$

$$Q_1^{q_i q_j} = \bar{q}_{jL}^{\alpha} \gamma_\mu q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} \gamma^\mu q_{iL}^{\beta} ,$$

$$Q_2^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jR}^{\beta} q_{iL}^{\beta} ,$$

$$Q_3^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jR}^{\beta} q_{iL}^{\alpha} ,$$

$$Q_4^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} q_{iR}^{\beta} ,$$

$$Q_5^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jL}^{\beta} q_{iR}^{\alpha} .$$

$$C_i(\Lambda) = F_i \frac{L_i}{\Lambda^2}$$

F_i: function of the NP flavour couplings

L_i: loop factor (in NP models with no tree-level FCNC)

Λ: NP scale (typical mass of new particles mediating ΔF=2 processes)

testing the TeV scale

$$C_i(\Lambda) = F_i \frac{L_i}{\Lambda^2}$$

The dependence of C on Λ changes depending on the flavour structure.
We can consider different flavour scenarios:

- Generic: $C(\Lambda) = \alpha/\Lambda^2$ — $F_i \sim 1$, arbitrary phase
- NMFV: $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$ — $F_i \sim |F_{SM}|$, arbitrary phase
- MFV: $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$ — $F_1 \sim |F_{SM}|$, $F_{i\neq1} \sim 0$, SM phase

α (L_i) is the coupling among NP and SM

- $\alpha \sim 1$ for strongly coupled NP
- $\alpha \sim \alpha_W$ (α_S) in case of loop coupling through weak (strong) interactions

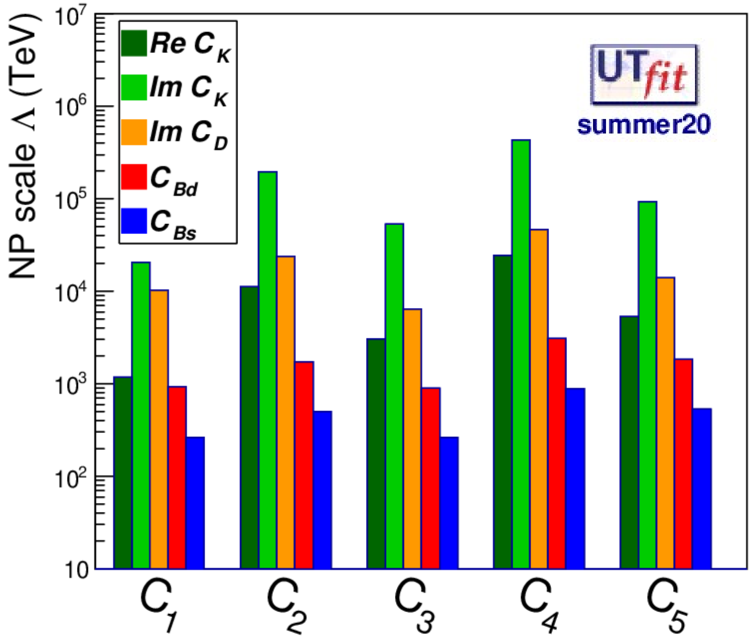
If no NP effect is seen lower bound on NP scale Λ

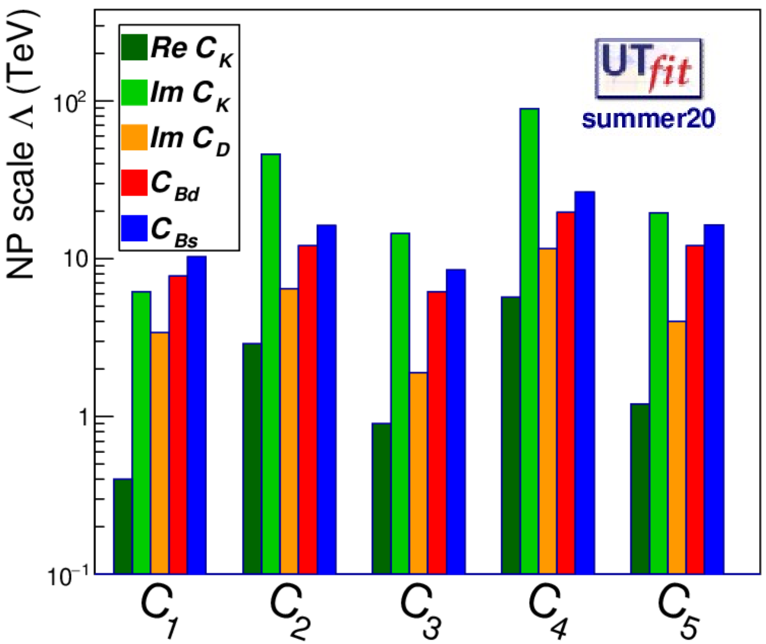
F is the flavour coupling and so
F_{SM} is the combination of CKM factors for the considered process

results from the Wilson coefficients

Generic: $C(\Lambda) = \alpha/\Lambda^2$,
$F_i \sim 1$, arbitrary phase
$\alpha \sim 1$ for strongly coupled NP

NMFV: $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$,
$F_i \sim |F_{SM}|$, arbitrary phase





$\Lambda > 4.3\ 10^5$ TeV

Lower bounds on NP scale
(at 95% prob.)

$\Lambda > 89$ TeV

$\alpha \sim \alpha_W$ in case of loop coupling through weak interactions

$\Lambda > 1.3\ 10^4$ TeV

$\alpha \sim \alpha_W$ in case of loop coupling through weak interactions

$\Lambda > 2.7$ TeV

for lower bound for loop-mediated contributions, simply multiply by α_s (~ 0.1) or by α_W (~ 0.03).

conclusions

- SM analysis displays very good (improved) overall consistency
- Still open discussion on semileptonic inclusive vs exclusive: exclusive fit shows tension, V_{cb} now showing the biggest discrepancy..
- UTA provides determination of NP contributions to $\Delta F=2$ amplitudes. It currently leaves space for NP at the level of 20-30%
- So the scale analysis points to high scales for the generic scenario and at the limit of LHC reach for weak coupling. Indirect searches are not only complementary to direct searches, but they might be the main way to glimpse at new physics.

Back up slides

lattice QCD inputs

updated in early 2020

Observables	Measurement
B_K	0.756 ± 0.016
f_{Bs}	0.2301 ± 0.0012
f_{Bs}/f_{Bd}	1.208 ± 0.005
B_{Bs}/B_{Bd}	1.032 ± 0.038
B_{Bs}	1.35 ± 0.06

FLAG 2019 suggests to take the most precise between the Nf=2+1+1 and Nf=2+1 averages.
We quote, instead, the weighted average of the Nf=2+1+1 and Nf=2+1 results with the error rescaled when chi2/dof > 1, as done by FLAG for the Nf=2+1+1 and Nf=2+1 averages separately

Contribution to the mixing amplitutes

analytic expression for the contribution to the mixing amplitudes

Lattice QCD

$$\langle \bar{B}_q | \mathcal{H}_{\rm eff}^{\Delta B=2} | B_q \rangle_i = \sum_{j=1}^{5} \sum_{r=1}^{5} \left(b_j^{(r,i)} + \eta\, c_j^{(r,i)} \right) \eta^{a_j} C_i(\Lambda) \langle \bar{B}_q | Q_r^{bq} | B_q \rangle$$

arXiv:0707.0636: for "magic numbers" a,b and c, $\eta = \alpha_s(\Lambda)/\alpha_s(m_t)$

analogously for the K system

$$\langle \bar{K}^0 | \mathcal{H}_{\rm eff}^{\Delta S=2} | K^0 \rangle_i = \sum_{j=1}^{5} \sum_{r=1}^{5} \left(b_j^{(r,i)} + \eta\, c_j^{(r,i)} \right) \eta^{a_j} C_i(\Lambda)\, R_r \langle \bar{K}^0 | Q_1^{sd} | K^0 \rangle$$

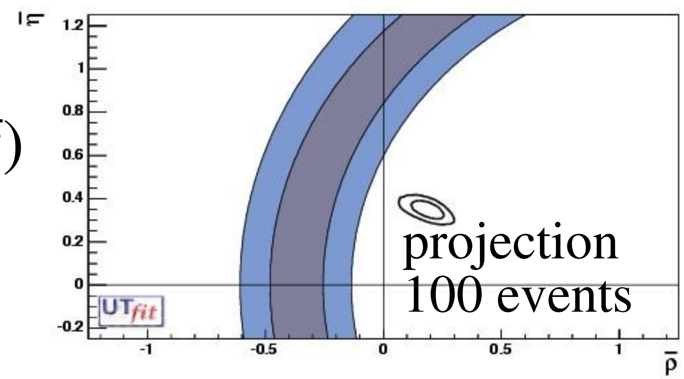
to obtain the p.d.f. for the Wilson coefficients $C_i(\Lambda)$ at the new-physics scale, we switch on one coefficient at a time in each sector and calculate its value from the result of the NP analysis.

some old plots coming back to fashion:

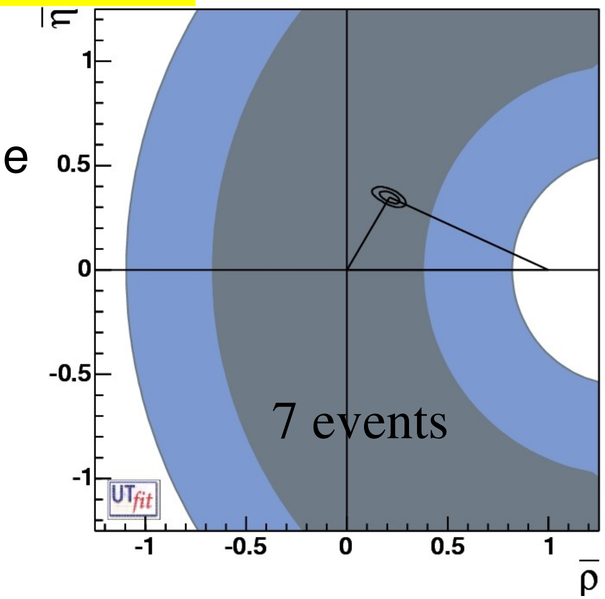
As NA62 and KOTO are analysing data:

$BR(K^+ \rightarrow \pi^+ \nu \bar{\nu})$

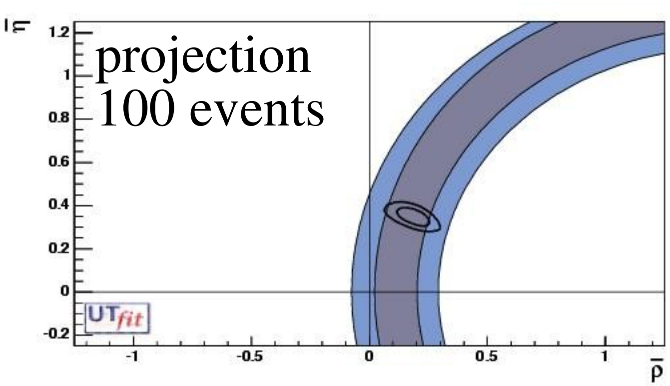
E949 central value

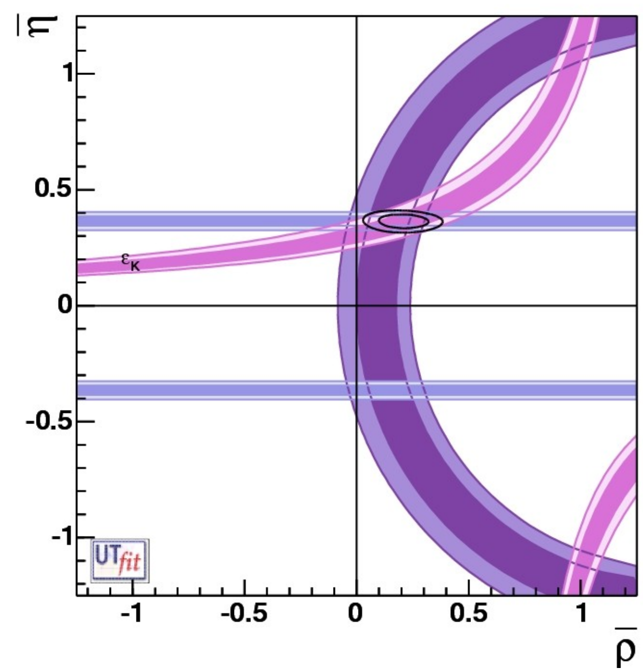


2007 global fit area



SM central value

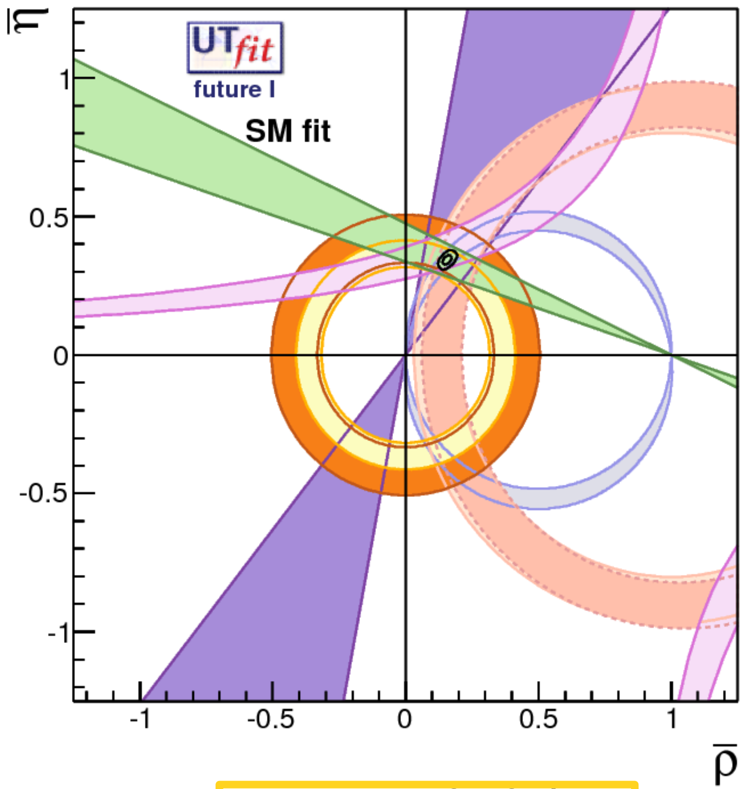




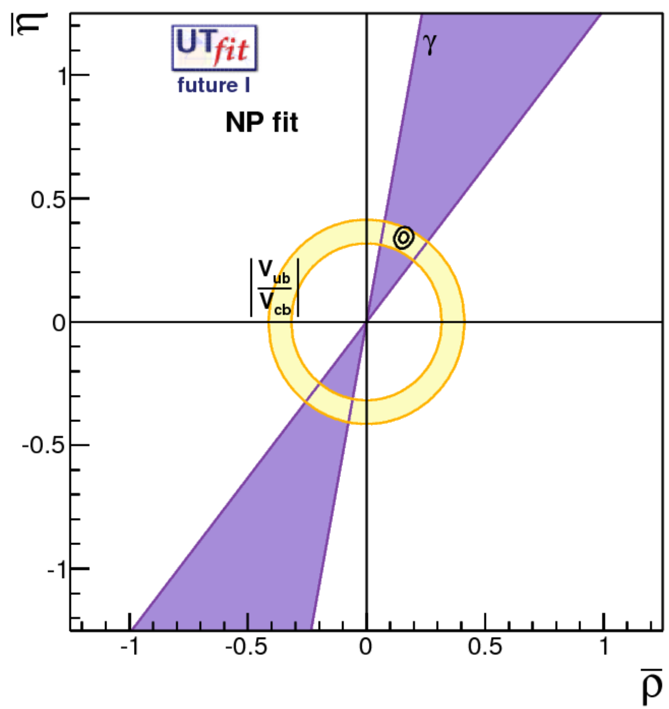
including
$BR(K^0 \rightarrow \pi^0 \nu \bar{\nu})$
SM central value

Look at the near future

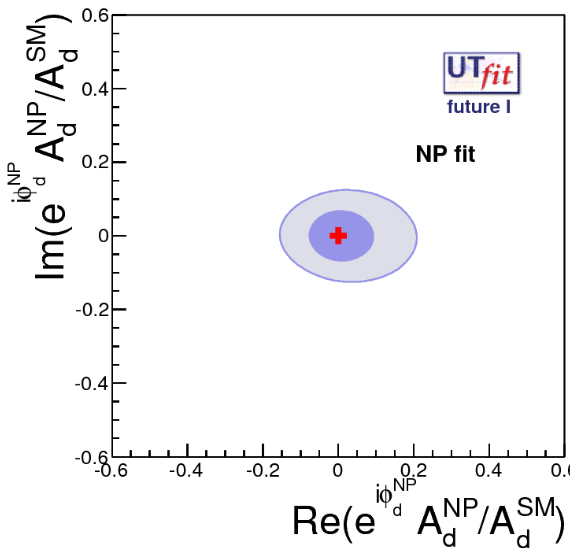
future I scenario:
errors from
Belle II at 5/ab
+ LHCb at 10/fb

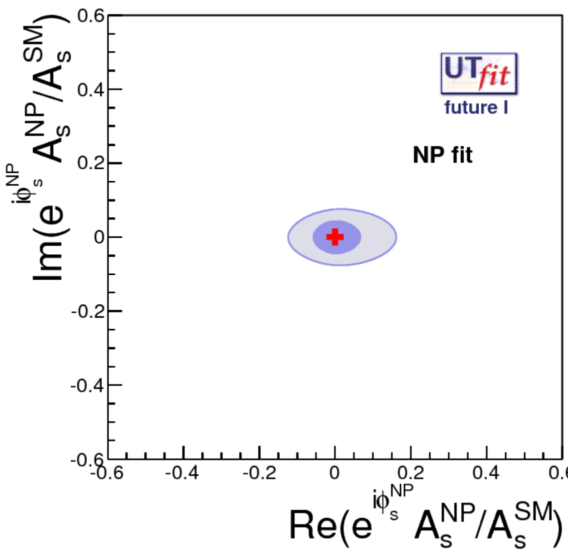


$\rho = \pm\, 0.015$
$\eta = \pm\, 0.015$



$\rho = \pm\, 0.016$
$\eta = \pm\, 0.019$





$\bar{\rho} = 0.154 \pm 0.015$
$\bar{\eta} = 0.344 \pm 0.013$

current sensitivity

$\bar{\rho} = 0.150 \pm 0.027$
$\bar{\eta} = 0.363 \pm 0.025$

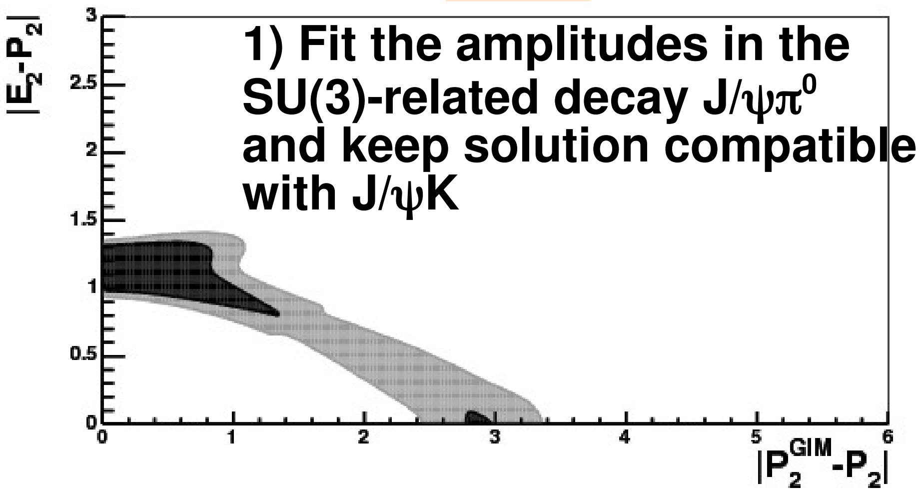
Theory error on sin2β:

A.Buras, L.Silvestrini
Nucl.Phys.B569:3-52(2000)

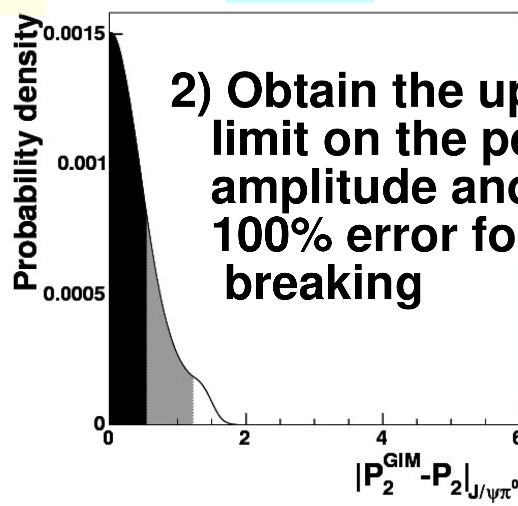
Channel	Cl.	E_1	E_2	EA_2	A_2	P_1	P_2	P_3	P_1^{GIM}	P_2^{GIM}	P_3^{GIM}	P_4	P_4^{GIM}
			$\frac{1}{N}$	$\frac{1}{N^2}$	$\frac{1}{N}$	$\frac{1}{N}$	$\frac{1}{N^2}$	1	$\frac{1}{N}$	$\frac{1}{N^2}$		$\frac{1}{N^3}$	$\frac{1}{N^3}$
$B_d \to J/\psi K^0$	C	–	λ^2	–	–	–	λ^2	–	–	λ^4	–	–	–
$B_d \to \pi^0 J/\psi$	D	–	λ^3	λ^3	–	–	λ^3	–	–	λ^3	–	$[\lambda^3]$	$[\lambda^3]$

$V^*_{cb}V_{cs}$ $V^*_{tb}V_{ts}$ $V^*_{ub}V_{us}$

$V^*_{cb}V_{cd}$ $V^*_{tb}V_{td}$ $V^*_{ub}V_{ud}$

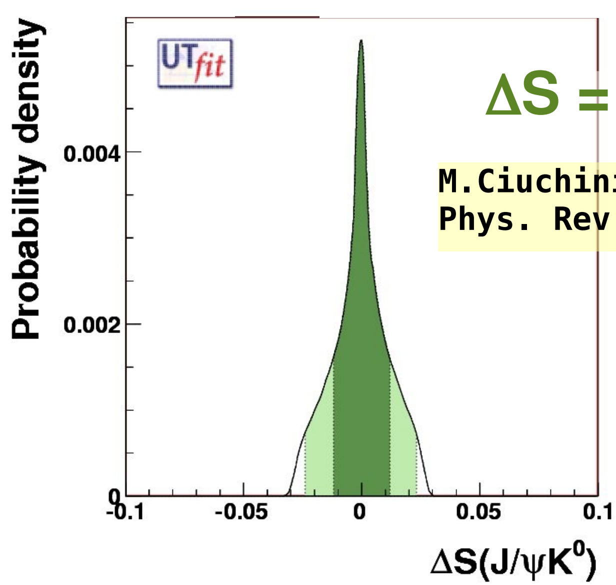


1) Fit the amplitudes in the SU(3)-related decay $J/\psi\pi^0$ and keep solution compatible with $J/\psi K$



2) Obtain the upper limit on the penguin amplitude and add 100% error for SU(3) breaking

3) Fit the amplitudes in $J/\psi K^0$ imposing the upper bound on the CKM suppressed amplitude and extract the error on sin2β



$\Delta S = 0.000 \pm 0.012$

M.Ciuchini, M.Pierini, L.Silvestrini
Phys. Rev. Lett. 95, 221804 (2005)